Normal-conducting RF Photo Injectors for Free Electron Lasers

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Content:

- Introduction: FELs and their electron sources
- Different NC photo injectors
 - at SLAC for LCLS (low average current)
 - at Berkeley for NGLS (high average current)
 - at PITZ for FLASH and European XFEL (medium average current)
- Outlook and Summary





European XFEL - a next generation light source

The XFEL will deliver:

- > ultra-short pulses (≤ 100 fs)
 → ultra-fast dynamics, "molecular movies"
- ultra-high peak brilliance
 investigations of matter under extreme conditions (Xe²¹⁺)
- transverse spatial coherence
 imaging of single nanoscale objects, possibly down to individual macromolecules (no crystallisation needed !!)

Why brilliance is ~10E+8 higher ? Synchrotrons: $P \sim N \cdot e^2$ FELs (coherence):

$$P \sim (N \cdot e)^2 = N^2 \cdot e^2$$
$$N \sim 10E+8$$





SASE FEL: How does it work?

Coherent motion is all we need !!







SASE FEL: How does it work?



One XFEL key component: → the high brightness electron source

Why electron injector is so important ???

Any linac based short wavelength, high brilliance light source (e.g. SASE-FELs) contains the following main components:

- electron source





property of linacs: beam quality will DEGRADE during acceleration in linac

➔ electron source has to produce lowest possible emittance !!



What is Emittance ?

long.: $\mathcal{E}_{z} \sim (e^{-} \text{ bunch length}) \bullet (energy spread of e^{-} \text{ bunch})$ trans.: $\mathcal{E}_{x,v} \sim (e^{-} \text{ beam size}) \bullet (e^{-} \text{ beam angular divergence})$



E = 6 dimensional phase space volume occupied by given number of particles



effect of <u>acceleration</u> on transverse emittance (adiabatic damping):

Why electron injector is so important ...

• Why emittance must be small ...

FLASH



- XFEL goal: 0.9 mm mrad@injector = 1.4 mm mrad@undulator
- if even smaller emittance \Rightarrow new horizons:

shorter wavelength, higher repetition rate





XFEL

Situation in 1999

ICFA workshop on high brightness beams at UCLA in autumn 1999:

Summary talk of P. O'Shea (U Maryland, USA) on electron source developments:



"Goal for community in next years:

Get transverse normalized emittance of 1 mm mrad @ bunch charge of 1 nC !!!"





Different NC photo injectors

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NC RF Gun Design for LCLS

pulsed / CW	pulsed		
single bunch charge	1.0 / 0.2 nC		
single bunch rep rate	120 Hz		
average current	120 / 24 nA		
norm. trans. emittance (rms, slice)	1.0 / 0.8 mm mrad @ 135 MeV		
rf frequency	2856 MHz		

modified UCLA/BNL/SLAC 1.6 cell S-band gun:

- larger mode separation (3.5 \rightarrow 15 MHz)
- larger iris radius, reduced iris surface field
- dual rf feed, z coupling, racetrack shape
- field probes in both cells
- increased cooling channels
- klystron pulse shaping \rightarrow reduced dissipated power

→improved emittance and stability



From Paul Emma, "First Lasing"-talk at FEL2009 in Liverpool: **Injector Transverse Projected Emittance ~0.5 µm**

Exceptional beam quality from S-band Cu-cath. RF gun...

focusina solenoid

cathode flange

Time-sliced x-emittance: $0.4 \,\mu\text{m}$



From Paul Emma, "First Lasing"-talk at FEL2009 in Liverpool: Measurements and Simulations for 20-pC Bunch at 14 GeV MEASURED SLICE EMITTANCE SIMULATED FEL PULSES





Simulation at 1.5 Å based on measured injector & linac beam & *Elegant* tracking, with CSR, at 20 pC.



linac beam & Elegant tracking, with CSR & 20 pC.

Some more experimental results from the LCLS gun:

> Phy.Rev. ST AB **11**, 030703 (2008):

1 nC, 100A bunch, 135 MeV: $\varepsilon_{n, x}$ (95%) = 1.14 mm mrad, $\mathcal{E}_{n,v}(95\%) = 1.06 \text{ mm mrad}$

P. Emma, "Beam Brightness Measurements in the LCLS Injector", Mini-WS on compact XFELs using HBB, LBNL, Berkeley, USA, 2010.

Best Proj. Emittance Measurements at 0.25 nC (35 A) 0.25 nC, 35A, 135 MeV



First day new drive laser was used...

 $\mathcal{E}_{n, x}$ (95%) = 0.32 mm mrad,

 $\epsilon_{n, v}$ (95%) = 0.39 mm mrad





J. Frisch, "Operation and Upgrades of the LCLS", LINAC2010, Tsukuba, Japan: 0.02 nC, 135 MeV: $\varepsilon_{n, x}(95\%) = 0.18 \text{ mm mrad}, \ \varepsilon_{n, y}(95\%) = 0.19 \text{ mm mrad}$

Frank Stephan | NC RF photo injectors for FELs | LA3NET Workshop, CERN

Different NC photo injectors

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Berkeley approach: NC RF gun design for CW operation (1)

Design requirements

\rightarrow high duty cycle, high average beam current, low emittance:

- repetition rate: up to ~1 MHz, CW
- bunch charge: from ~1 pC to ~1 nC,
- normalized beam emittance: from sub 10⁻⁷ (low charge) to 10⁻⁶ m,
- beam energy at the gun exit: ≥ -500 keV (space charge),
- electric field at the cathode: $\geq \sim 10 \text{ MV/m}$ (space charge limit),
- bunch length control at cathode: from ~1ps to 10s of ps (to handle space charge effects, for allowing different modes of operation),

module

slots

lon pump

low RF

frequency

- compatibility with magnetic fields in the cathode and gun regions (mainly for emittance compensation) Pumping
- operational vacuum pressure: 10⁻⁹ 10⁻¹¹ Torr (high QE photo-cathodes),
- "easy" installation and conditioning of different kind of cathodes,
- high reliability compatible with the operation of a user facility.

(from Journal of Modern Optics, Vol. 58, No. 16, 1419-1437)

Berkeley approach: NC RF gun design for CW operation (2)

pulsed / CW	CW			
single bunch charge	0.01 to 1 nC			
single bunch rep rate	up to 1 MHz			
average current	up to 1 mA			
norm. trans. emittance (rms, slice)	< 1.0 mm mrad (design)			
rf frequency	187 MHz			









Berkeley approach: NC RF gun design for CW operation (3)



The following design parameters have already been demonstrated at Berkeley:

- full conditioning to max. RF power (100kW in CW mode, allow for 750keV beam energy) $\sqrt{}$
- max dark current ~8µA @ 19.5MV/m \rightarrow should be OK for high rep rate FEL operation $\sqrt{}$
- first photoelectrons at 1 MHz from temporary molybdenum cathode $~\sqrt{}$
- beam energy: 745 ± 41 keV $\sqrt{}$
- cathode field: >10 MV/m $\sqrt{}$
- pressure: ~5x10⁻¹¹ Torr after baking (RF off, 1 of 20 NEG modules activated)

(factor 3-4 higher with RF on)

(from PRST-AB 15, 103501 (2012))



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The PITZ gun (Photo Injector Test facility at DESY, Zeuthen site)



Capable of high average power \rightarrow long electron bunch trains (SC linac) \rightarrow

Very low normalized transverse emittance



Some parameters of FLASH and European XFEL



	FLASH	European XFEL
final beam energy	1.2 GeV	17.5 GeV
max. repetition rate	10 Hz	10 Hz
max. train length	800 µs	650 µs
bunch spacing	1 – 20 µs	0.2 – 1 µs
required injector emittance (1 nC)	2 mm mrad	0.9 mm mrad
SASE output wavelength	4.1 – 45 nm	0.05 – 6.4 nm









Improvement of the RF gun phase stability (gun has no field probe)





DESY

How we measure the transverse projected emittance

Single slit scan technique

- Emittance Measurement SYstem (EMSY) consists of horizontal / vertical actuators with
 - YAG / OTR screens
 - 10 / 50 μm slits
- > Beam size is measured @ slit position using screen

pixel intensity

Beam local divergence is estimated from beamlet sizes @ observation screen (12 bit camera)





value

2D corrected normalized RMS emittance

$$\mathcal{E}_{n} = \underbrace{\sigma_{x}}_{\sqrt{\langle x^{2} \rangle}} \beta \gamma \sqrt{\langle x^{2} \rangle \cdot \langle x'^{2} \rangle} - \langle xx' \rangle^{2}$$

correction factor (>1) introduced to correct for low intensity losses from beamlet measurements

 $\boldsymbol{\sigma_{x}}$ - RMS beam size measured with YAG screen at slit location

SQRT(<x²>) - RMS beam size at slit location estimated from slit positions and beamlet intensities

"100% RMS emittance"

(conservative estimate)

High stability of the measurement



Example, measurement on 5.5.2011: 1nC, Emittance vs. Imain (gun SP phase=6deg)

For all measurements f250 lenses and 2x2 binning were used

398	0.567	0.501	1.104	0.739	0.533	2	23	2	22	16	25	1.034	0.892	15	25	1.068	0.812	1.5	1.3
397	0.359	0.367	0.732	0.583	0.363	1	25	1	24	10	25	0.719	0.650	10	25	0.725	0.616	1.3	1.2
396	0.307	0.320	0.722	0.479	0.313	-	21	-	21	5	25	0.689	0.584	7	24	0.705	0.529	1.5	1.3
395	0.285	0.279	0.770	0.424	0.282	1	18	1	18	3	25	0.779	0.571	5	25	0.774	0.492	1.8	1.6
394	0.269	0.254	0.816	0.408	0.261	1	15	1	16	3	25	0.882	0.646	5	24	0.848	0.513	2.0	1.7
393	0.266	0.234	0.940	0.453	0.249	1	13	1	15	4	23	1.015	0.723	5	25	0.977	0.572	2.1	1.7

 $\epsilon_x = (0.707 \pm 0.032) \text{ mm mrad}$ $\epsilon_y = (0.686 \pm 0.024) \text{ mm mrad}$ $\epsilon_{xy} = (0.697 \pm 0.026) \text{ mm mrad}$



Emittance vs. Laser Spot size for various charges in 2011



Minimum emittance

Charge, nC	PITZ, 100%, mm mrad	LCLS, 95% mm mrad
2	1.25	
1	0.70	1.10
0.7		0.80
0.25	0.33	0.35
0.1	0.21	
0.02	0.12	0.19

 PITZ is setting a new benchmark for minimum emittance at a given charge + is capable of operation at fairly high duty cycle !!

LCLS data:

P. Emma, "Beam Brightness Measurements in the LCLS Injector", Mini-WS on compact XFELs using HBB, LBNL, Berkeley, USA, 2010.

J. Frisch, "Operation and Upgrades of the LCLS", LINAC2010, Tsukuba, Japan.

PITZ data:

• M. Krasilnikov et. al., "Experimentally minimized beam emittance from an L-band photoinjector", PRST-AB 15, 100701 (2012).

Core Emittance for various bunch charges

Idea: Cut low intensity region of MEASURED phase space (i.e. remove



2011 PITZ 1.8: measurements vs. simulations



Reasons of discrepancy for high Q? \rightarrow Emission from the cathode?



- Direct plug-in machine settings into ASTRA does not produce 1nC at the gun operation phase (+6deg), whereas 1nC and even higher charge (~1.2nC) are experimentally detected
- Simulated (ASTRA) phase scans w/o Schottky effects (solid thick lines) have different shapes than the experimentally measured (thin lines with markers)

1.6 measured charge (XYrms=0.3mm, 0deg) 1.4 С simulated charge (XYrms=0.3mm, 0deg) charge@LOW.ICT1, 1.2 1.0 0.8 0.6 0.4 bunch 0.2 0.0 0.0 0.2 0.4 0.6 0.8 1.0 1.2 1.4 1.6 1.8 2.0 ~ laser intensity, nC

Measured and simulated laser energy scan (1nC)

• Laser intensity (LT) scan for the MMMG phase (red curve with markers) shows higher saturation level, whereas the simulated charge even goes slightly down while the laser intensity (Qbunch) increases

Photo emission (bunch charge) needs more detailed modeling in simulations.





Key components for successful operation of NC RF photoinjectors

> good RF stability

- > good cooling water stability (increasingly difficult for high average RF power)
 - e.g. better than ±0.05 degrees are needed at 50 kW average power for the PITZ case
- robust and high quantum efficiency (QE) photo cathodes (especially important for high average current)
 - e.g. the PITZ photo cathode system was developed by INFN-LASA Milano
 → see talk of Daniele Sertore

> stable photo cathode laser system with high flexibility

- e.g. the PITZ photo cathode laser system was developed by the Max-Born-Institute in Berlin
 - \rightarrow see talk of Ingo Will
- High flexibility of photo cathode laser opens new research possibilities
 → excursion to plasma acceleration
- fancy shaping of the photo cathode laser pulses can further improve the beam quality





Photo cathode laser system installed at PITZ



Photo cathode laser: temporal pulse shaping

(for more MBI see I. Will)

Multicrystal birefringent pulse shaper containing 13 crystals

Gaussian:







Simulated pulse-stacker





birefringent shaper, **13** crystals

→ important for good beam quality → high flexibility → new options !



Particle driven plasma wakefield acceleration (PDPWA)



LAOLA @ PITZ: Studies for Particle Driven Plasma Acceleration

- 1) Self-modulation of electron beam (proof of principle for CERNs AWAKE exp.)
 - use high flexibility of photo cathode laser system:



Example: flat-top e-beam through plasma cell:



2) Study high transformer ratio

- resonantly drive plasma wave with specially shaped electron bunch (5 bunchlets inside the bunch)
- \rightarrow high transformation ratio:



to be sent to bunch compressor





New option for the photo cathode laser \rightarrow 3D ellipsoid



> Benefits from 3D ellipsoidal laser pulses for ALL linac driven light sources:

- 30-50% lower average slice emittance → higher brilliance
- ~pure sinusoidal longitudinal phase space +3rd harm. → simplify/allow required compression
- ~no beam halo → better signal/noise, reduced radiation damage
- less sensitive to machine settings → higher stability

Summary

- Normal conducting RF photo injectors are a mature and reliable technology to produce high brightness beams for FEL applications
- From low to medium average currents (~30µA) a new benchmark in optimizing photo injector performance was demonstrated

TABLE IV. Core *xy* emittance (mm mrad) measured for various bunch charges and gun phases. Only statistical errors are shown (see text).

Bunch charge	Gun phase	0%	Charge cut 5%	10%	20%	
2.0 nC	0 deg	1.558 ± 0.050	1.324 ± 0.045	1.173 ± 0.039	0.936 ± 0.031	from
2.0 nC	6 deg	1.251 ± 0.064	1.064 ± 0.054	0.939 ± 0.048	0.728 ± 0.037	
1.0 nC	0 deg	0.833 ± 0.038	0.711 ± 0.033	0.629 ± 0.029	0.511 ± 0.024	IVI. Kraslinikov et al.,
1.0 nC	6 deg	0.696 ± 0.020	0.596 ± 0.017	0.529 ± 0.015	0.427 ± 0.013	PRST-AB 15 , 100701
0.25 nC	0 deg	0.328 ± 0.010	0.289 ± 0.009	0.260 ± 0.008	0.213 ± 0.006	(2012).
0.10 nC	0 deg	0.212 ± 0.006	0.188 ± 0.006	0.170 ± 0.006	0.141 ± 0.006	
0.02 nC	0 deg	0.121 ± 0.001	0.108 ± 0.001	0.098 ± 0.001	0.082 ± 0.002	

- ➤ For high duty cycle / CW applications the main challenges (cooling, vacuum properties) can be solved → see Berkeley approach
- High brightness photo injectors are also interesting beyond FEL applications, e.g. plasma acceleration, electron diffraction (REGAE), ...
- Sincere Acknowledgements: to all international colleagues for discussions over the years, colleagues at Hamburg and Zeuthen, the PITZ team, ...



