

6D Beam Envelope Equations: An Ultrafast Computational Approach for Interactive Modeling of Accelerator Structures

By:

M. Dayyani Kelisani

1. Deutsche Elektronen-Synchrotron (DESY), Zeuthen, Germany.
2. Paul Scherrer Institute (PSI), Villigen, CH-5232, Switzerland.
3. Institute for Research in Fundamental Sciences (IPM), Tehran, Iran.
4. European Organization for Nuclear Researches (CERN), Geneva, Switzerland.

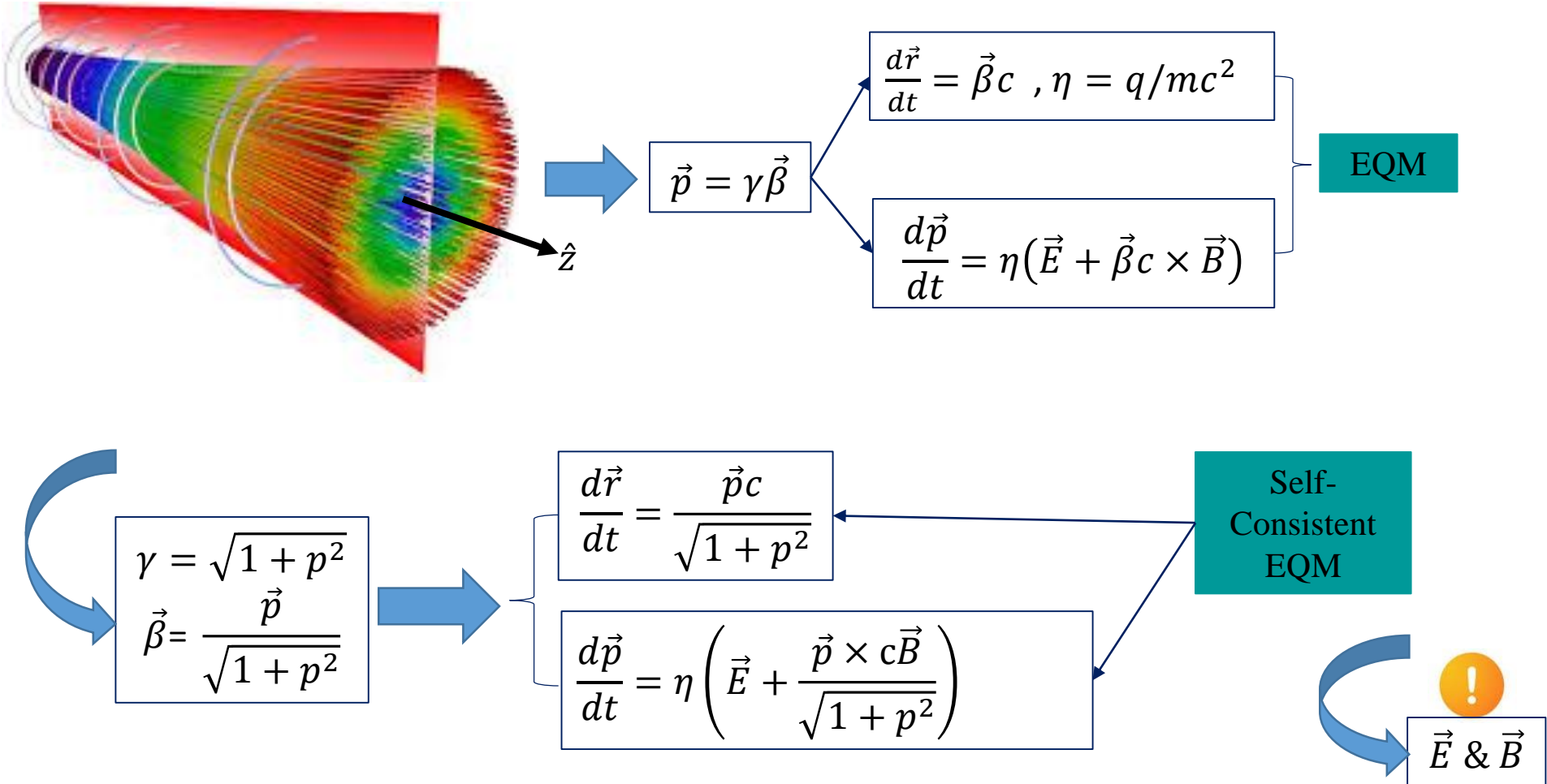
Contents

1. Equations of Motion and Ray Equations

2. 6D Envelope Equations

3. Applications

1.1 Equations of Motion



1.1 Equations of Motion

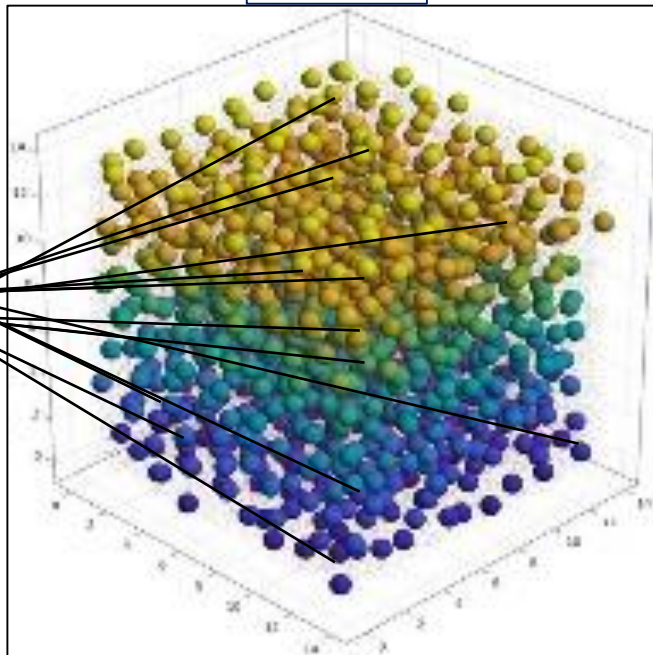
Ex Fields

$$\vec{E} = \vec{E}^e + \vec{E}^s$$

$$\vec{B} = \vec{B}^e + \vec{B}^s$$

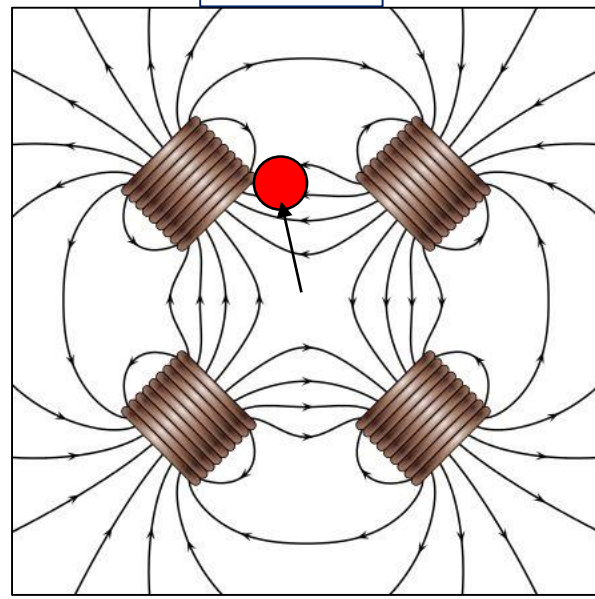
Sp Fields

Sp Fields



Costing-Time
 $\propto N_p^2$

Ex Fields



Costing-Time
 $\propto N_p$

A bunch of 10k Macro-Particles at each time step requires for



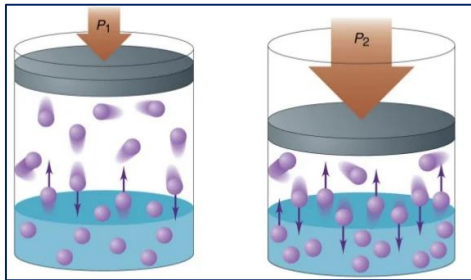
Calculation of 200M EM-fields

+

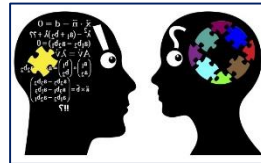
30k Coupled Differential Equation

1.1 Equations of Motion

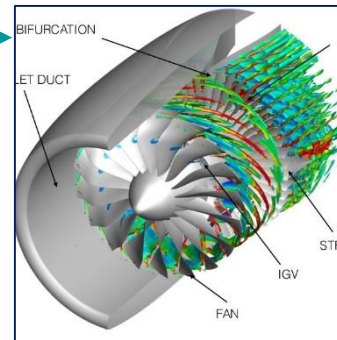
Gas



• A Many-Particle System



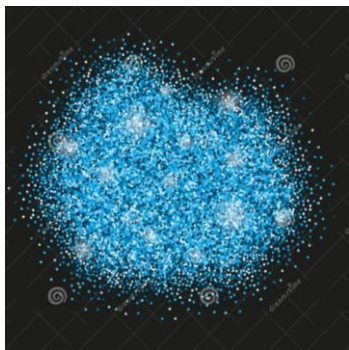
Hydro-Dynamics Consideration



$$pV = Nk_B T$$

Relates all general parameters of the gas totally independent of the gas detailed structure

Bunch

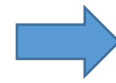
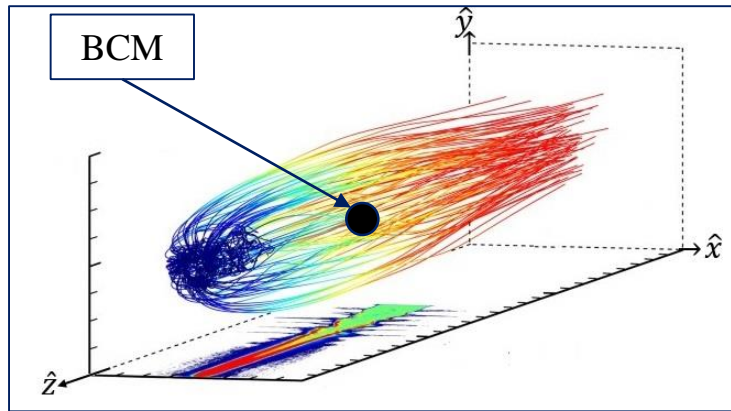


• A Many-Particle System

Electro-Dynamics



1.2 Ray Equations



BCM is not affected by the space-charge forces

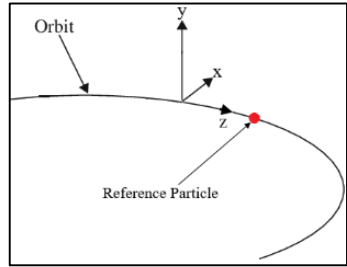
Definition of an artificial reference particle located always at the bunch center of mass $\vec{r}_0 = z_0 \hat{z}$

Reference Particle SEQM

$$\frac{d\gamma_0}{dz_0} = \eta E_z(z_0 \hat{z})$$

$$\frac{dt_0}{dz_0} = \frac{\gamma_0}{c\sqrt{\gamma_0^2 - 1}}$$

Relativistic Factor: γ_0



$$\frac{d}{dt_0} = \frac{dz_0}{dt_0} \frac{d}{dz_0} = \beta_0 c \frac{d}{dz_0}$$

Normalized Velocity: $\beta_0 \hat{z}$

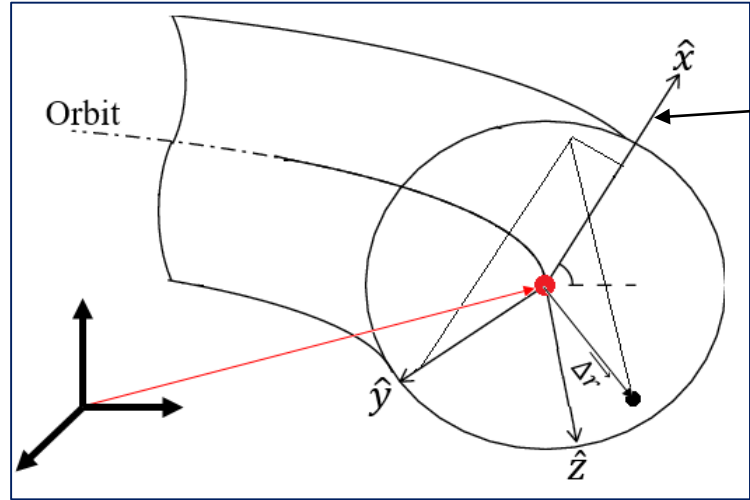
$$t_0 \rightarrow \vec{r}_0 = z_0 \hat{z}$$

Normalized Momentum: $\vec{p}_0 = \gamma_0 \beta_0 \hat{z}$

$$\frac{d(mc^2 \gamma_0)}{dt_0} = q \left(\vec{E}(\vec{r}_0) + \vec{\beta}_0 c \times \vec{B}(\vec{r}_0) \right) \cdot \vec{\beta}_0 c$$

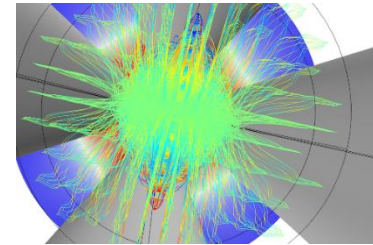
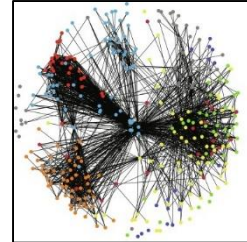
$$\rightarrow \frac{d\gamma_0}{dt_0} = \eta c E_z(\vec{r}_0) \beta_0$$

1.2 Ray Equations



Moving Frame

We report all the particle coordinates with respect to the moving frame by a 6D vector $[\Delta x, \Delta y, \Delta z, p_x, p_y, p_z]$,



Ray Equations

EQM

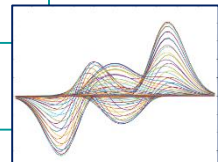
+

$$\frac{d}{dt_0} = \beta_0 c \frac{d}{dz_0}$$

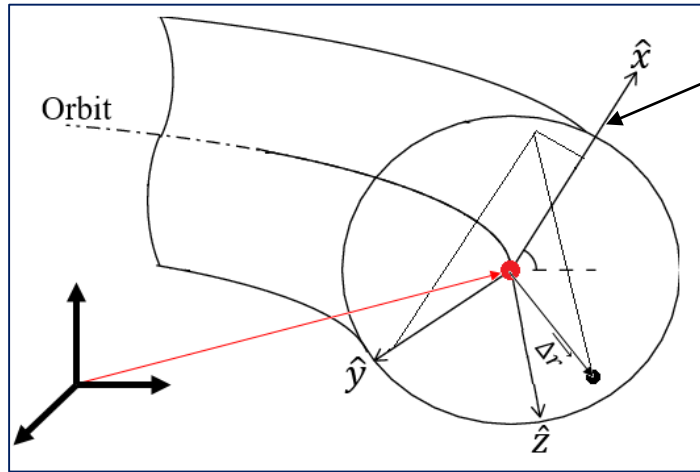


$$\gamma' = \frac{\eta \vec{E} \cdot \vec{\beta}}{\beta_0}, \quad \beta_u = \beta_0 (\hat{u} \cdot \hat{z} + \Delta u')$$

$$\Delta u'' + \frac{\gamma_0'}{p_0^3} \beta_u = \frac{\eta}{\gamma \beta_0^2} [(\vec{E} + \vec{\beta} \times c\vec{B})_u - (\vec{E} \cdot \vec{\beta}) \beta_u]$$

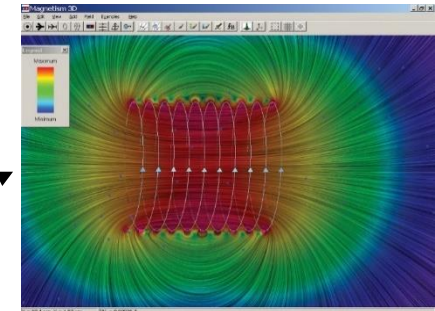


1.2 Ray Equations



Moving Frame

Solenoidal Magnet



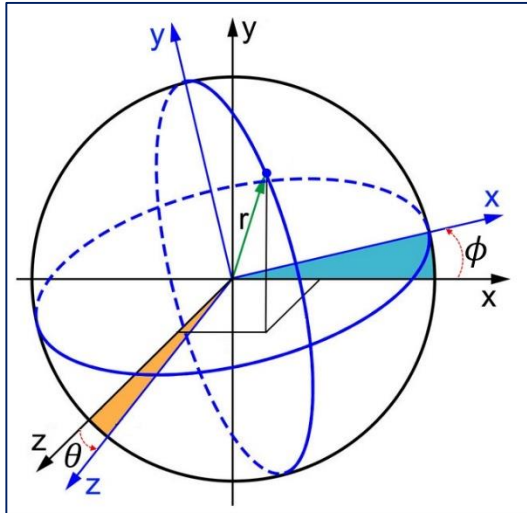
Coupling Due to Solenoidal Fields

$$\vec{B}^{sm} = \left(\hat{z} - \hat{x} \frac{\Delta x}{2} \frac{d}{dz} - \hat{y} \frac{\Delta y}{2} \frac{d}{dz} \right) \mu^{sm}(z)$$

$$\mu^{sm} = (\vec{B}^{sm} \cdot \hat{z}) |_{\Delta \vec{r}=0}$$

$$\begin{aligned} & (\Delta x + i\Delta y)'' \\ & + \left(\frac{\gamma_0 \gamma_0'}{\gamma_0^2 p_0^2} + \frac{\eta(\vec{E} \cdot \vec{\beta})}{\gamma \beta_0} + \frac{i\eta c}{\gamma \beta_0} \mu^{sm} \right) (\Delta x + i\Delta y)' \\ & + \frac{i\eta c \beta_z}{2\gamma \beta_0^2} \frac{d\mu^{sm}}{dz} (\Delta x + i\Delta y) \\ & = \frac{\eta}{\gamma \beta_0^2} (E_x + iE_y) \\ & + \frac{i\eta c \beta_z}{\gamma \beta_0^2} \left((B_x - B_x^{sm}) + i(B_y - B_y^{sm}) \right) \end{aligned}$$

1.2 Ray Equations



$$\theta = 0 \rightarrow [\Delta z]_{prev} = [\Delta z]_{new}$$

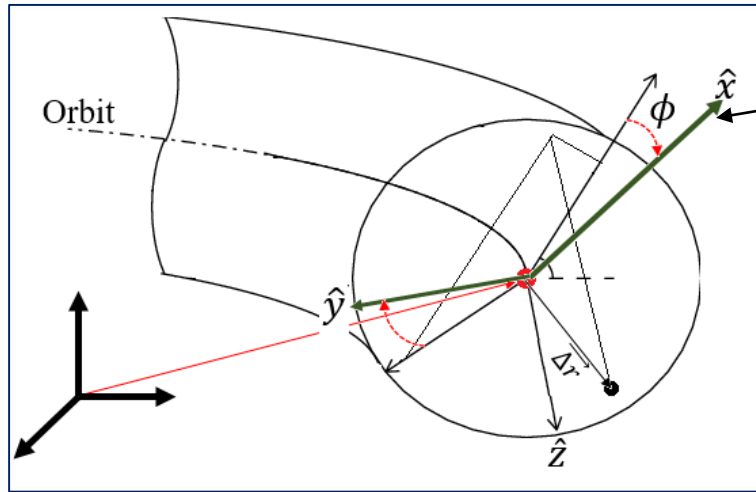
$$[\Delta x + i\Delta y]_{prev} = [\Delta x + i\Delta y]_{new} e^{i\phi}$$

$$\phi' = -\frac{\eta c}{2\gamma\beta_0} \mu^{sm}$$

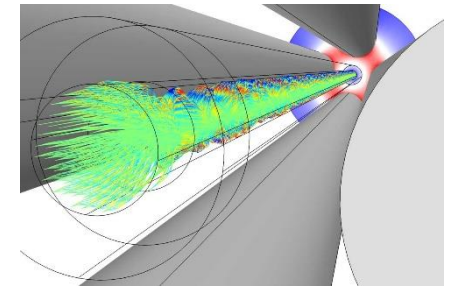
$$\begin{aligned} (\Delta x + i\Delta y)'' + \left(\frac{\gamma_0 \gamma_0'}{\gamma_0^2 p_0^2} + \frac{\eta(\vec{E} \cdot \vec{\beta})}{\gamma\beta_0} \right) (\Delta x + i\Delta y)' \\ = + \frac{\eta}{\gamma\beta_0^2} (E_x + iE_y) - \left(\frac{\eta c}{2\gamma\beta_0} \mu^{sm} \right)^2 (\Delta x + i\Delta y) \\ + \frac{i\eta c \beta_z}{\gamma\beta_0^2} \left((B_x - B_x^{sm}) + i(B_y - B_y^{sm}) \right) \end{aligned}$$

No coupling at least up to the first order

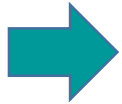
1.2 Ray Equations



Larmor Moving Frame



Ray Equations in Larmor Moving Frame



$$\gamma' = \frac{\eta \vec{E} \cdot \vec{\beta}}{\beta_0}, \quad \beta_u = \beta_0 (\hat{u} \cdot \hat{z} + \Delta u')$$

$$\begin{aligned} & \Delta u'' + \frac{\gamma_0'}{p_0^3} \beta_u \\ &= - \frac{\eta (\vec{E} \cdot \vec{\beta}) \beta_u}{\gamma \beta_0^2} - \frac{(1 - \hat{u} \cdot \hat{z}) \eta^2 c^2 (\overline{B^{sm}} \cdot \hat{z})^2 \Delta u}{4 \gamma^2 \beta_0^2} \\ &+ \frac{\eta (\vec{E} + \vec{\beta} c \times \vec{B} - (1 - \hat{u} \cdot \hat{z}) \vec{\beta} c \times \overline{B^{sm}}) \cdot \hat{u}}{\gamma \beta_0^2} \end{aligned}$$

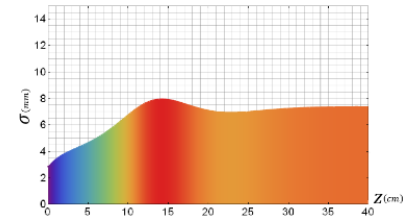
solenoidal magnetic field

2.1 6D Envelope Equations



Envelope Equations

$$\sigma_u'' + \frac{\gamma_0 \gamma_0'}{\gamma_0^2 p_0^2} \sigma_u' = F_u^e + F_u^s + F_u^\varepsilon$$



$$F_u^e = -\frac{\eta}{\beta_0^2 \sigma_u} \left\langle \frac{(\vec{E}^e \cdot \vec{\beta}) \beta_u \Delta u}{\gamma} \right\rangle - \frac{\eta^2 c^2 (1 - \hat{u} \cdot \hat{z})}{4 \beta_0^2 \sigma_u} \left\langle \frac{(\vec{B}^{sm} \cdot \hat{z})^2 \Delta u^2}{\gamma^2} \right\rangle + \frac{\eta}{\beta_0^2 \sigma_u} \left\langle \frac{(\vec{E}^e + \vec{\beta} c \times \vec{B}^e - (1 - \hat{u} \cdot \hat{z}) \vec{\beta} c \times \vec{B}^{sm}) \cdot \Delta u \hat{u}}{\gamma} \right\rangle$$

Ex F-Type Force

$$F_u^s = -\frac{\eta}{\beta_0^2 \sigma_u} \left\langle \frac{(\vec{E}^s \cdot \vec{\beta}) \beta_u \Delta u}{\gamma} \right\rangle + \frac{\eta}{\beta_0^2 \sigma_u} \left\langle \frac{(\vec{E}^s + \vec{\beta} c \times \vec{B}^s) \cdot \Delta u \hat{u}}{\gamma} \right\rangle$$

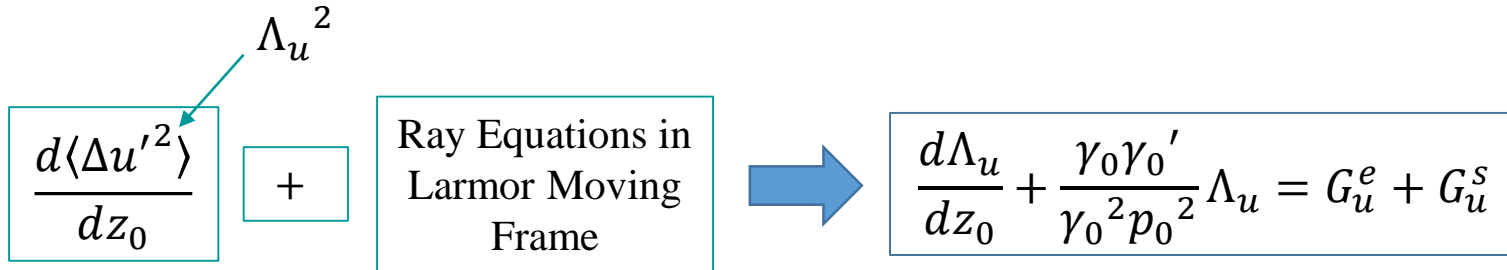
Sp F-Type Force

$$F_u^\varepsilon = \frac{\Lambda_u^2 - \sigma_u'^2}{\sigma_u}, \quad \Lambda_u = \sqrt{\langle \Delta u'^2 \rangle}$$

Emit F-type Force



2.1 6D Envelope Equations



$$G_u^e = -\frac{\eta}{\beta_0^2 \Lambda_u} \left\langle \frac{(\vec{E}^e \cdot \vec{\beta}) \beta_u \Delta u'}{\gamma} \right\rangle - \frac{\eta^2 c^2 (1 - \hat{u} \cdot \hat{z})}{4 \beta_0^2 \Lambda_u} \left\langle \frac{(\overline{B^{sm}} \cdot \hat{z})^2 \Delta u \Delta u'}{\gamma^2} \right\rangle + \frac{\eta}{\beta_0^2 \Lambda_u} \left\langle \frac{(\vec{E}^e + \vec{\beta} c \times \vec{B}^e - (1 - \hat{u} \cdot \hat{z}) \vec{\beta} c \times \overline{B^{sm}}) \cdot \Delta u' \hat{u}}{\gamma} \right\rangle$$

Ex G-type Force

6D Integrals



$$G_u^s = -\frac{\eta}{\beta_0^2 \Lambda_u} \left\langle \frac{(\vec{E}^s \cdot \vec{\beta}) \beta_u \Delta u'}{\gamma} \right\rangle + \frac{\eta}{\beta_0^2 \Lambda_u} \left\langle \frac{(\vec{E}^s + \vec{\beta} c \times \vec{B}^s) \cdot \Delta u' \hat{u}}{\gamma} \right\rangle$$

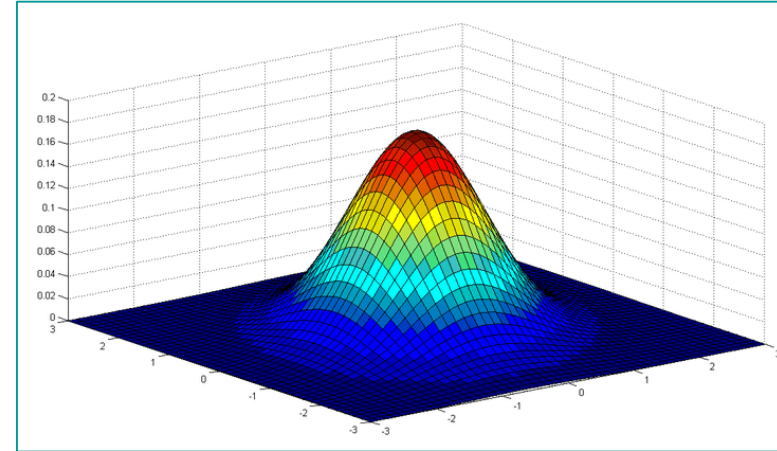
Sp G-type Force

2.3 Bunch Distribution

Dynamics independent of the bunch detailed structure

No matter what is the distribution just it must provide all 6 parameters

$$\mathcal{F}_u(\Delta u, \Delta u') = \frac{p_0}{2\varepsilon_u} e^{-\left[\left(\frac{p_0 \Lambda_u \Delta u}{\sqrt{2}\varepsilon_u} \right)^2 - 2 \frac{p_0^2 \sigma_u \sigma_{u'}}{2\varepsilon_u^2} \Delta u \Delta u' + \left(\frac{p_0 \Lambda_u \Delta u'}{\sqrt{2}\varepsilon_u} \right)^2 \right]}$$



$$\sqrt{\langle \Delta u^2 \rangle} = \sigma_u$$

$$\sqrt{\langle \Delta u'^2 \rangle} = \Lambda_u$$

$$p_0 \sqrt{\langle \Delta u^2 \rangle \langle \Delta u'^2 \rangle - \langle \Delta u \Delta u' \rangle^2} = \varepsilon_u$$

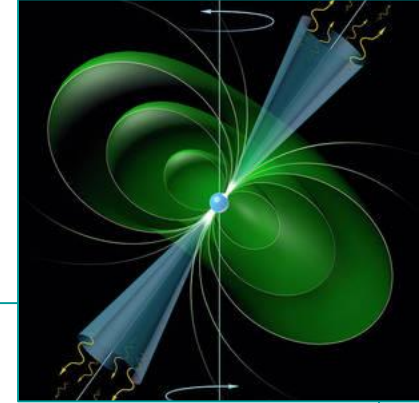
6D Gaussian distribution

$$\mathcal{F}(\Delta x, \Delta y, \Delta z, \Delta x', \Delta y', \Delta z') = \mathcal{F}_x(\Delta x, \Delta x') \mathcal{F}_y(\Delta y, \Delta y') \mathcal{F}_z(\Delta z, \Delta z')$$

3.1 Space-Charge Forces

Space Charge Forces

$$f_b = \frac{\eta q_b}{8\pi\sqrt{\pi}\epsilon_0}$$



$$F_x^s \cong + \frac{f_b}{\beta_0^2 \gamma_0^3} \times \frac{\alpha_x}{\sigma_x \sigma_z}$$

$$- \frac{f_b}{\gamma_0} \left\{ \frac{(\Lambda_x^2 + 2\sigma_x'^2)\alpha_x - \sigma_x'^2 \alpha_{xx}}{2\sigma_x \sigma_z} + \frac{8\Lambda_y^2 \alpha_x - \sigma_y'^2 \alpha_{xy}}{16\sigma_x \sigma_z} + \frac{(1 - p_0^2)(2\Lambda_z^2 \alpha_x - \sigma_z'^2 \alpha_{xz})}{4\sigma_x \sigma_z} \right\}$$

$$- \frac{f_b}{\gamma_0} \left\{ \frac{(\Lambda_x^2 + 2\sigma_x'^2)\alpha_x - \sigma_x'^2 \alpha_{xx}}{\sigma_x \sigma_z} + \frac{\sigma_x' \sigma_y' (8\alpha_y - \alpha_{xy})}{8\sigma_z \sigma_y} + \frac{(1 - p_0^2) \sigma_x' \sigma_z' (2\gamma_0^2 \sigma_z^2 \alpha_z - \sigma_x \sigma_y \alpha_{xz})}{2\sigma_x \sigma_y \gamma_0^2 \sigma_z^2} \right\}$$

$$G_x^s = + \frac{f_b \sigma_x'}{\beta_0^2 \gamma_0^3 \Lambda_x} \frac{\alpha_x}{\sigma_x \sigma_z}$$

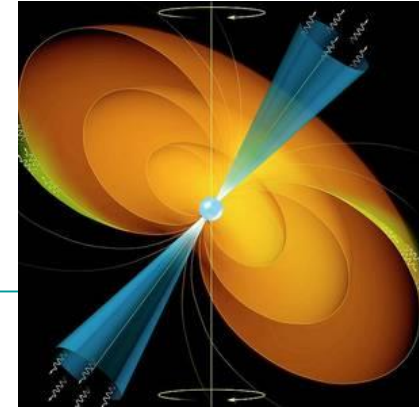
$$- \frac{f_b \sigma_x'}{2\gamma_0 \Lambda_x} \left\{ 3 \frac{3\Lambda_x^2 \alpha_x - \sigma_x'^2 \alpha_{xx}}{\sigma_x \sigma_z} + \frac{8\Lambda_y^2 \alpha_x - \sigma_y'^2 \alpha_{xy}}{8\sigma_x \sigma_z} + (1 - p_0^2) \frac{2\Lambda_z^2 \alpha_x - \sigma_z'^2 \alpha_{xz}}{2\sigma_x \sigma_z} \right\}$$

$$- \frac{f_b \sigma_y'}{8\gamma_0 \Lambda_x} \left\{ \frac{8\Lambda_x^2 \alpha_y - \sigma_x'^2 \alpha_{xy}}{\sigma_y \sigma_z} \right\} - \frac{f_b \sigma_z'}{2\gamma_0 \Lambda_x} \left\{ (1 - p_0^2) \frac{2\gamma_0^2 \sigma_z^2 \Lambda_x^2 \alpha_z - \sigma_x'^2 \sigma_x \sigma_y \alpha_{xz}}{\sigma_x \sigma_y \gamma_0^2 \sigma_z^2} \right\}$$

3.1 Space-Charge Forces

Space Charge Forces

$$f_b = \frac{\eta q_b}{8\pi\sqrt{\pi}\epsilon_0}$$

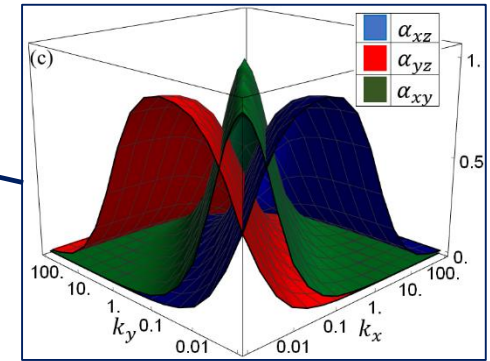
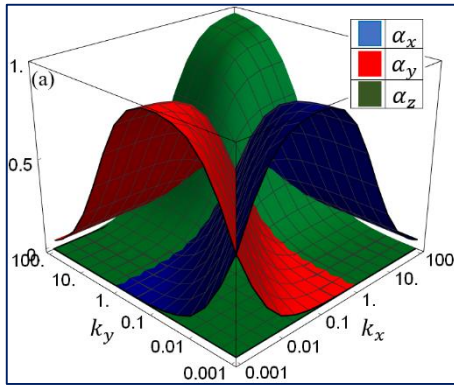


$$F_z^s \cong + \frac{f_b}{\beta_0^2 \gamma_0^3} \times \frac{\alpha_z}{\sigma_x \sigma_y} - \frac{f_b}{\gamma_0} \left\{ \frac{3(1-p_0^2)(2(\Lambda_z^2 + 2\sigma_z'^2)\alpha_z - \sigma_z'^2\alpha_{zz})}{4\sigma_x \sigma_y} + \frac{2\gamma_0^2 \sigma_z^2 \Lambda_x^2 \alpha_z - \sigma_x'^2 \sigma_x \sigma_y \alpha_{xz}}{4\sigma_x \sigma_y \gamma_0^2 \sigma_z^2} \right\}$$

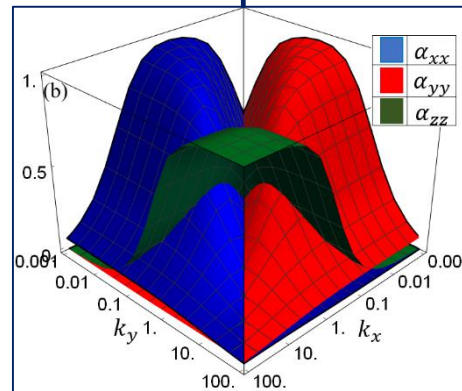
$$G_z^s = + \frac{f_b \sigma_z'}{\beta_0^2 \gamma_0^3 \Lambda_z} \times \frac{\alpha_z}{\sigma_x \sigma_y} - \frac{f_b \sigma_z'}{2\gamma_0 \Lambda_z} \left\{ 3(1-p_0^2) \frac{6\Lambda_z^2 \alpha_z - \sigma_z'^2 \alpha_{zz}}{2\sigma_x \sigma_y} + \frac{2\gamma_0^2 \sigma_z^2 \Lambda_x^2 \alpha_z - \sigma_x'^2 \sigma_x \sigma_y \alpha_{xz}}{2\sigma_x \sigma_y \gamma_0^2 \sigma_z^2} + \frac{2\gamma_0^2 \sigma_z^2 \Lambda_y^2 \alpha_z - \sigma_y'^2 \sigma_x \sigma_y \alpha_{yz}}{2\sigma_x \sigma_y \gamma_0^2 \sigma_z^2} \right\} - \frac{f_b \sigma_x'}{2\gamma_0 \Lambda_z} \left\{ \frac{2\Lambda_z^2 \alpha_x - \sigma_z'^2 \alpha_{xz}}{\sigma_x \sigma_z} \right\} - \frac{f_b \sigma_y'}{2\gamma_0 \Lambda_z} \left\{ \frac{2\Lambda_z^2 \alpha_y - \sigma_z'^2 \alpha_{yz}}{\sigma_y \sigma_z} \right\}$$

3.1 Space-Charge Forces

1 st order	$\{\alpha_x, \alpha_y, \alpha_z\}$
2 nd Order	$\{\alpha_{xx}, \alpha_{yy}, \alpha_{zz}\}$
3 rd Order	$\{\alpha_{xy}, \alpha_{xz}, \alpha_{yz}\}$



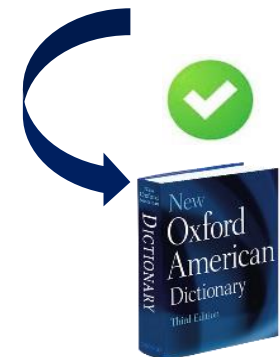
$$k_u = \frac{\sigma_u}{\gamma_0 \sigma_z}$$



$$\alpha_x = \int_0^{+\infty} \frac{k_x^2 s ds}{\sqrt{(s^2 + k_x^2)^3 (s^2 + k_y^2) (s^2 + 1)}}$$

$$\alpha_{xz} = \int_0^{+\infty} \frac{k_x^2 s ds}{\sqrt{(s^2 + k_x^2)^3 (s^2 + k_y^2) (s^2 + 1)^3}}$$

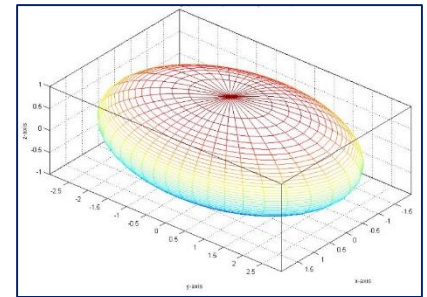
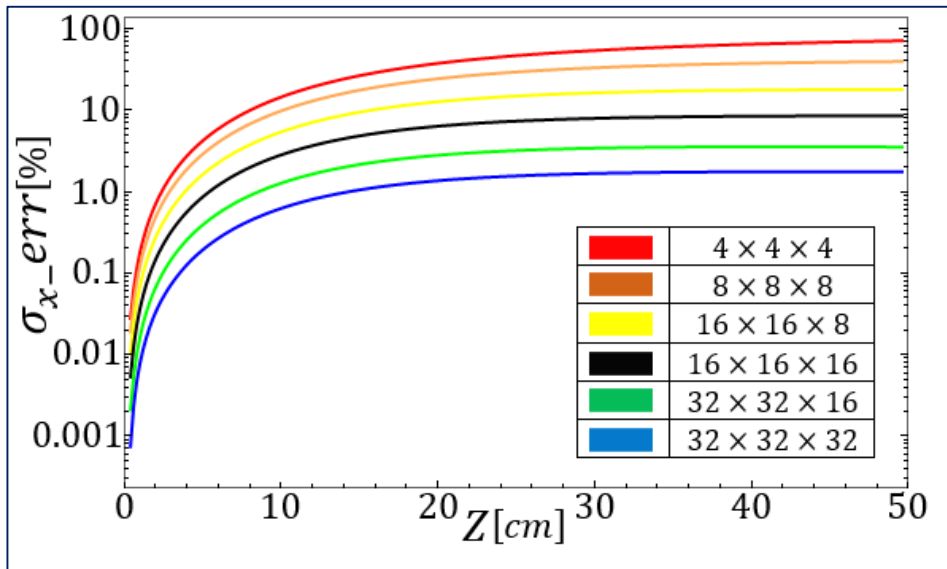
$$\alpha_{xx} = \int_0^{+\infty} \frac{3k_x^4 s ds}{2\sqrt{(s^2 + k_x^2)^5 (s^2 + k_y^2) (s^2 + 1)}}$$



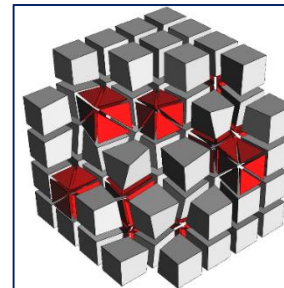
3.1 Space-Charge Forces

Specifications	$E_k [MeV]$	$\sigma_E [%]$	$\sigma_x [mm]$	$\sigma_y [mm]$	$\sigma_z [mm]$	$\epsilon_{nx} [\mu m]$	$\epsilon_{ny} [\mu m]$	$\epsilon_{nz} [\mu m]$
Value	5	1	$\sqrt{2}$	$2\sqrt{2}$	0.030	0.05	0.10	2.93

Astra Vs Exact Approach



ASTRA

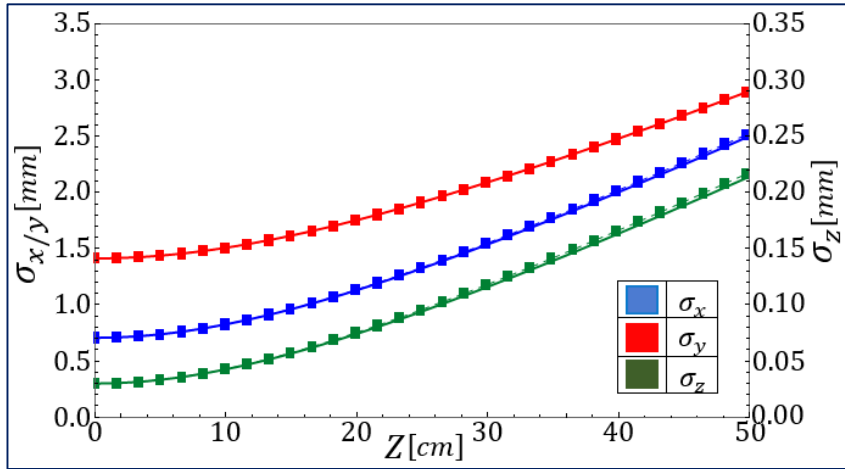


$$2^n \times 2^m \times 2^l$$



$2^5 \times 2^5 \times 2^5$
Simulation-Time :
> 6h

3.1 Space-Charge Forces

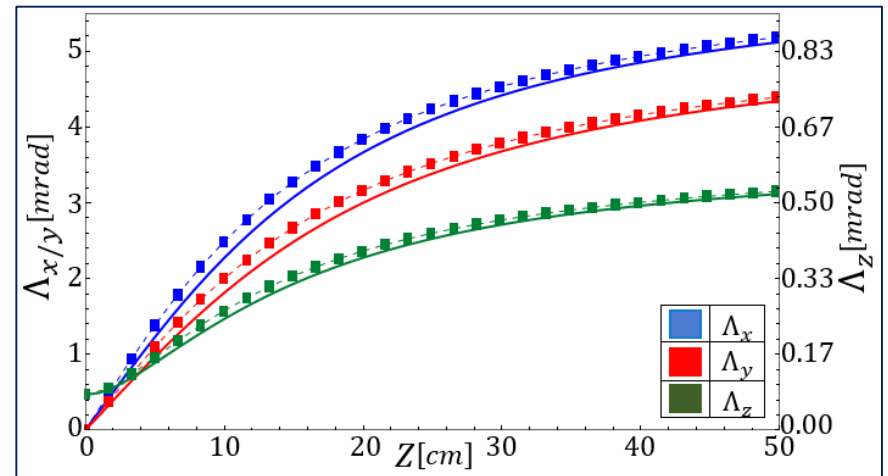


Simulation-Time : > 6h

ASTRA



Simulation-Time :
< 1s



3.2 Solenoidal Magnet Forces

$$F_x^e = -\frac{\eta^2 c^2}{4p_0^2} (\mu^{sm2} + \mu_z^{sm2} \sigma_z^2) \sigma_x$$

$$+ \eta^2 c^2 \mu^{sm} \mu_z^{sm} \sigma_x \sigma_z \sigma_z' + \frac{\eta^2 c^2 \mu^{sm2}}{4} \sigma_x (2\sigma_x'^2 + \Lambda_x^2 + \Lambda_y^2 + \Lambda_z^2)$$

$$+ \frac{\eta^2 c^2 \mu_z^{sm2}}{4} \sigma_x \sigma_z^2 (2\sigma_x'^2 + 2\sigma_z'^2 + \Lambda_x^2 + \Lambda_y^2 + \Lambda_z^2)$$

$$F_z^e = 0$$

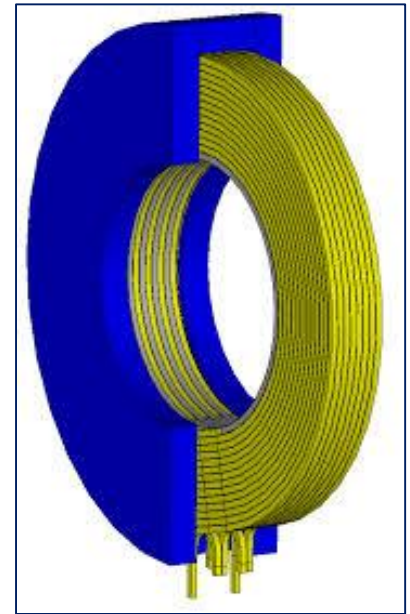
$$G_x^e = -\frac{\eta^2 c^2}{4p_0^2 \Lambda_x} (\mu^{sm2} + \mu_z^{sm2} \sigma_z^2) \sigma_x \sigma_x'$$

$$+ \frac{\eta^2 c^2 \mu^{sm} \mu_z^{sm}}{\Lambda_x} \sigma_x \sigma_x' \sigma_z \sigma_z' + \frac{\eta^2 c^2 \mu^{sm2}}{4\Lambda_x} \sigma_x \sigma_x' (3\Lambda_x^2 - \Lambda_y^2 - \Lambda_z^2)$$

$$+ \frac{\eta^2 c^2 \mu_z^{sm2}}{4\Lambda_x} \sigma_x \sigma_x' \sigma_z^2 (2\sigma_z'^2 + 3\Lambda_x^2 + \Lambda_y^2 + \Lambda_z^2)$$

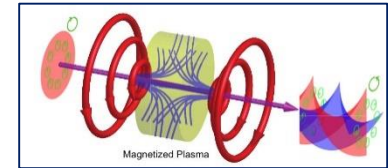
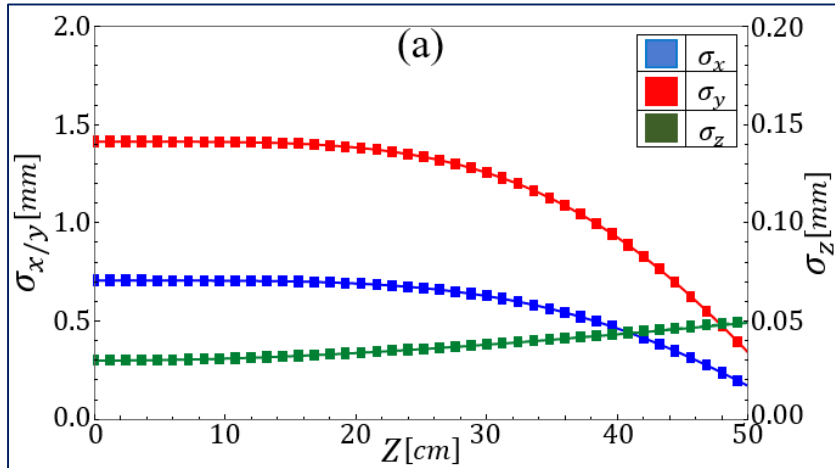
$$G_z^e = 0$$

$$\mu^{sm} = (\vec{B}^{sm} \cdot \hat{z}) |_{\Delta \vec{r}=0}$$



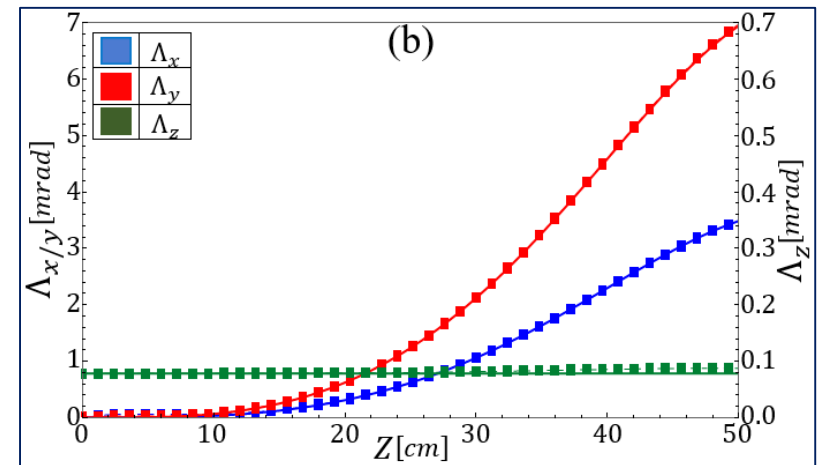
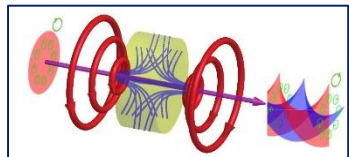
$$\mu_z^{sm} = \left(\frac{\partial}{\partial z} \vec{B}^{sm} \cdot \hat{z} \right) |_{\Delta \vec{r}=0}$$

3.2 Solenoidal Magnet Forces



$$\mu^{sm}(z) = 0.5 z$$

ASTRA



3.3 Quadrupole Magnet Forces

$$F_x^e = -k^{qm}\sigma_x + \frac{k^{qm}p_0^2}{2}\sigma_x\{(2 + \gamma_0^2)\Lambda_z^2 + 2\sigma_x'^2 + \Lambda_x^2 + \Lambda_y^2\}$$

$$F_y^e = +k^{qm}\sigma_y - \frac{k^{qm}p_0^2}{2}\sigma_y\{(2 + \gamma_0^2)\Lambda_z^2 + 2\sigma_y'^2 + \Lambda_y^2 + \Lambda_x^2\}$$

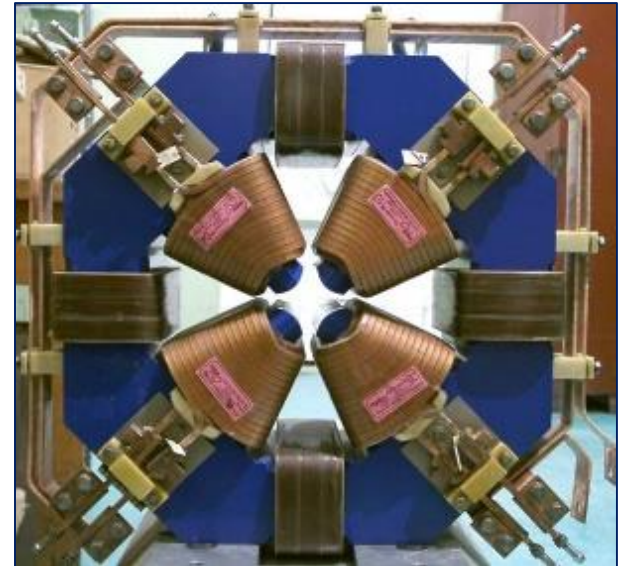
$$F_z^e = -k^{qm}p_0^2\sigma_z'(\sigma_x\sigma_x' - \sigma_y\sigma_y')$$

$$G_x^e = -\frac{k^{qm}}{\Lambda_x}\sigma_x\sigma_x' + \frac{k^{qm}p_0^2}{2\Lambda_x}\sigma_x\sigma_x'\{(2 + \gamma_0^2)\Lambda_z^2 + 3\Lambda_x^2 + \Lambda_y^2\}$$

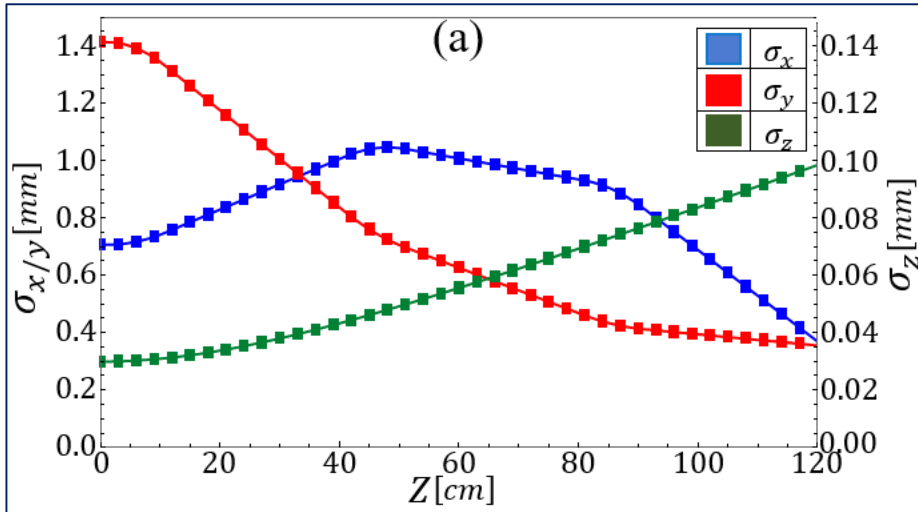
$$G_y^e = +\frac{k^{qm}}{\Lambda_y}\sigma_y\sigma_y' - \frac{k^{qm}p_0^2}{2\Lambda_y}\sigma_y\sigma_y'\{(2 + \gamma_0^2)\Lambda_z^2 + 3\Lambda_y^2 + \Lambda_x^2\}$$

$$G_z^e = -k^{qm}p_0^2\Lambda_z(\sigma_x\sigma_x' - \sigma_y\sigma_y')$$

$$\vec{B}^{qm} = \frac{p_0k^{qm}}{\eta c}(\Delta y\hat{x} + \Delta x\hat{y})$$

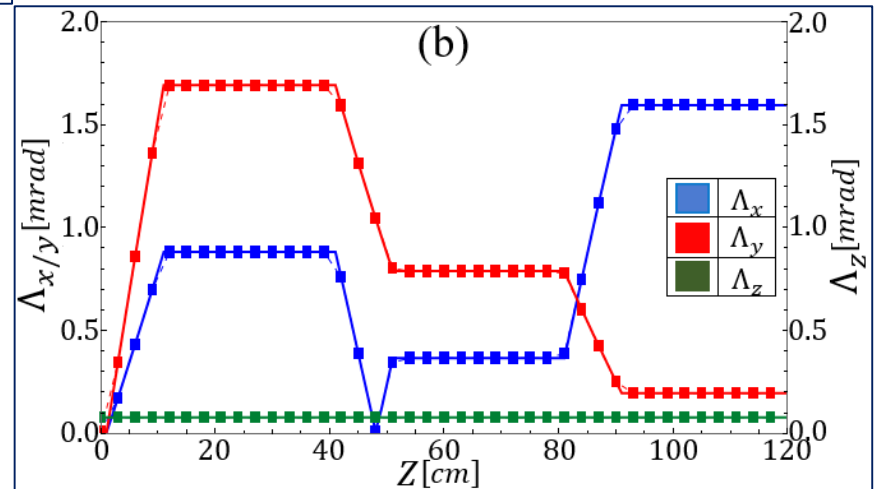
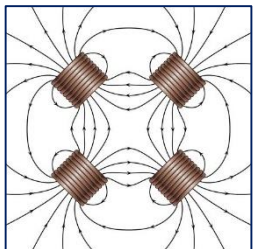


3.3 Quadrupole Magnet Forces



Element	k^{qm} [1/m ²]	L [cm]	z [cm]
Q_1	-12.221	10	6
Q_2	+12.018	10	46
Q_3	+13.776	10	86

ASTRA



3.4 Electrostatic and RF Forces

$$\begin{aligned}
 F_x^e &= -\frac{\eta(\mathcal{E}_z^{rf} + \beta_0 \mathcal{E}_t^{rf})}{2\gamma_0 \beta_0^2} \sigma_x - \frac{\eta \mathcal{E}^{rf}}{\gamma_0} \sigma_x' \\
 &+ \frac{\eta \mathcal{E}_z^{rf}}{2\gamma_0} \sigma_x' \{\sigma_y \sigma_y' - 2(1 - p_0^2) \sigma_z \sigma_z'\} + \frac{\eta \mathcal{E}_z^{rf}}{4\gamma_0} \sigma_x \{\gamma_0^4 \Lambda_z^2 + 2(2 + \gamma_0^2) \sigma_x'^2 + (2 + \gamma_0^2) \Lambda_x^2 + \gamma_0^2 \Lambda_y^2\} \\
 &+ \frac{\eta p_0 \mathcal{E}_t^{rf}}{4} \sigma_x \{(2 + \gamma_0^2) \Lambda_z^2 + 2\sigma_x'^2 + (\Lambda_x^2 + \Lambda_y^2)\} + \frac{\eta p_0^2 \mathcal{E}^{rf}}{2\gamma_0} \sigma_x' \{(2 + \gamma_0^2) \Lambda_z^2 + 3\Lambda_x^2 + \Lambda_y^2\}
 \end{aligned}$$

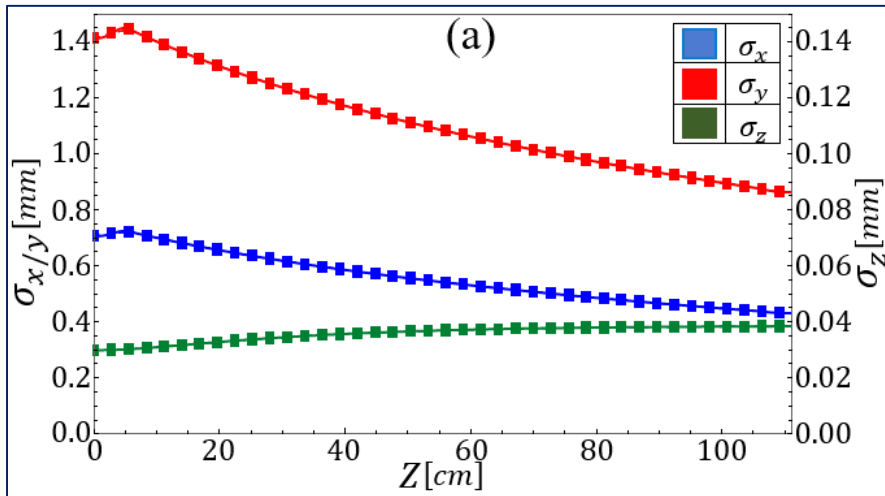
$$\begin{aligned}
 F_z^e &= +\frac{\eta \mathcal{E}_z^{rf}}{\gamma_0 p_0^2} \sigma_z - \frac{3\eta \mathcal{E}^{rf}}{\gamma_0} \sigma_z' \\
 &+ \frac{\eta(1 - p_0^2) \mathcal{E}_z^{rf}}{2\gamma_0} \sigma_z' (\sigma_x \sigma_x' + \sigma_y \sigma_y') - \frac{\eta p_0 \mathcal{E}_t^{rf}}{2} \sigma_z' (\sigma_x \sigma_x' + \sigma_y \sigma_y') \\
 &+ \frac{\eta \mathcal{E}_z^{rf}}{2\gamma_0} \sigma_z \{(6 - 3\gamma_0^2 - 2\gamma_0^4)(2\sigma_z'^2 + \Lambda_z^2) - (\Lambda_x^2 + \Lambda_y^2)\} + \frac{\eta p_0^2 \mathcal{E}^{rf}}{2\gamma_0} \sigma_z' \{3(2 + \gamma_0^2) \Lambda_z^2 + (\Lambda_x^2 + \Lambda_y^2)\}
 \end{aligned}$$

$$\begin{aligned}
 G_x^e &= -\frac{\eta(\mathcal{E}_z^{rf} + \beta_0 \mathcal{E}_t^{rf})}{2\gamma_0 \beta_0^2 \Lambda_x} \sigma_x \sigma_x' - \frac{\eta \mathcal{E}^{rf}}{\gamma_0} \Lambda_x \\
 &+ \frac{\eta \mathcal{E}_z^{rf}}{2\gamma_0} \Lambda_x \{3\sigma_x \sigma_x' + \sigma_y \sigma_y' - 2(1 - p_0^2) \sigma_z \sigma_z'\} + \frac{\eta \gamma_0 \mathcal{E}_z^{rf}}{4\Lambda_x} \sigma_x \sigma_x' (\gamma_0^2 \Lambda_z^2 + 3\Lambda_x^2 + \Lambda_y^2) \\
 &+ \frac{\eta p_0 \mathcal{E}_t^{rf}}{4\Lambda_x} \sigma_x \sigma_x' \{(3 + p_0^2) \Lambda_z^2 + 3\Lambda_x^2 + \Lambda_y^2\}
 \end{aligned}$$

$$\begin{aligned}
 G_z^e &= +\frac{\eta \mathcal{E}_z^{rf}}{\gamma_0 p_0^2 \Lambda_z} \sigma_z \sigma_z' - \frac{3\eta \mathcal{E}^{rf}}{\gamma_0} \Lambda_z \\
 &+ \frac{\eta(1 - p_0^2) \mathcal{E}_z^{rf}}{2\gamma_0} \Lambda_z (\sigma_x \sigma_x' + \sigma_y \sigma_y') - \frac{\eta p_0 \mathcal{E}_t^{rf}}{2} \Lambda_z (\sigma_x \sigma_x' + \sigma_y \sigma_y') \\
 &+ \frac{\eta \mathcal{E}_z^{rf}}{2\gamma_0 \Lambda_z} \sigma_z \sigma_z' \{3(6 - 3\gamma_0^2 - 2\gamma_0^4) \Lambda_z^2 - (\Lambda_x^2 + \Lambda_y^2)\}
 \end{aligned}$$

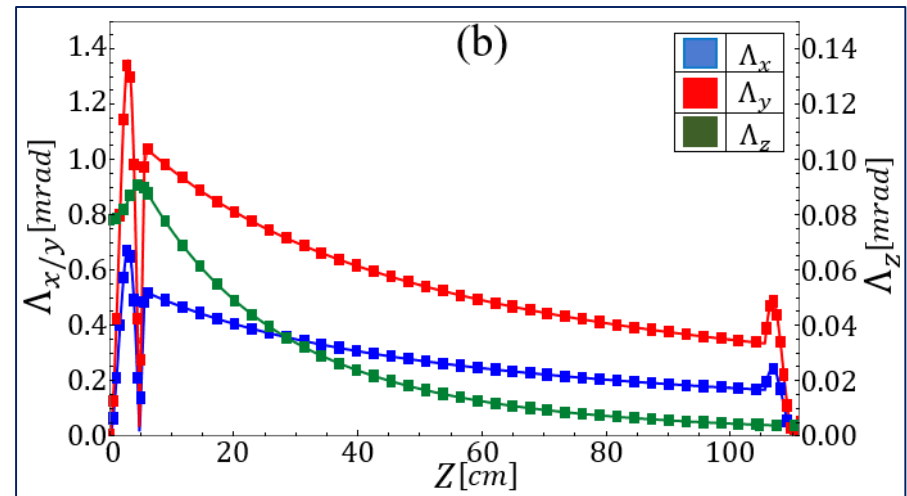
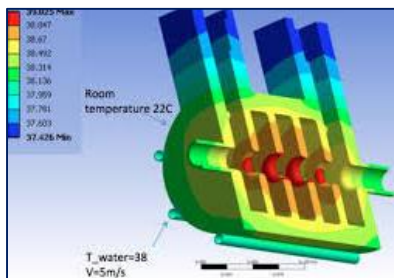


3.4 Electrostatic and RF Forces



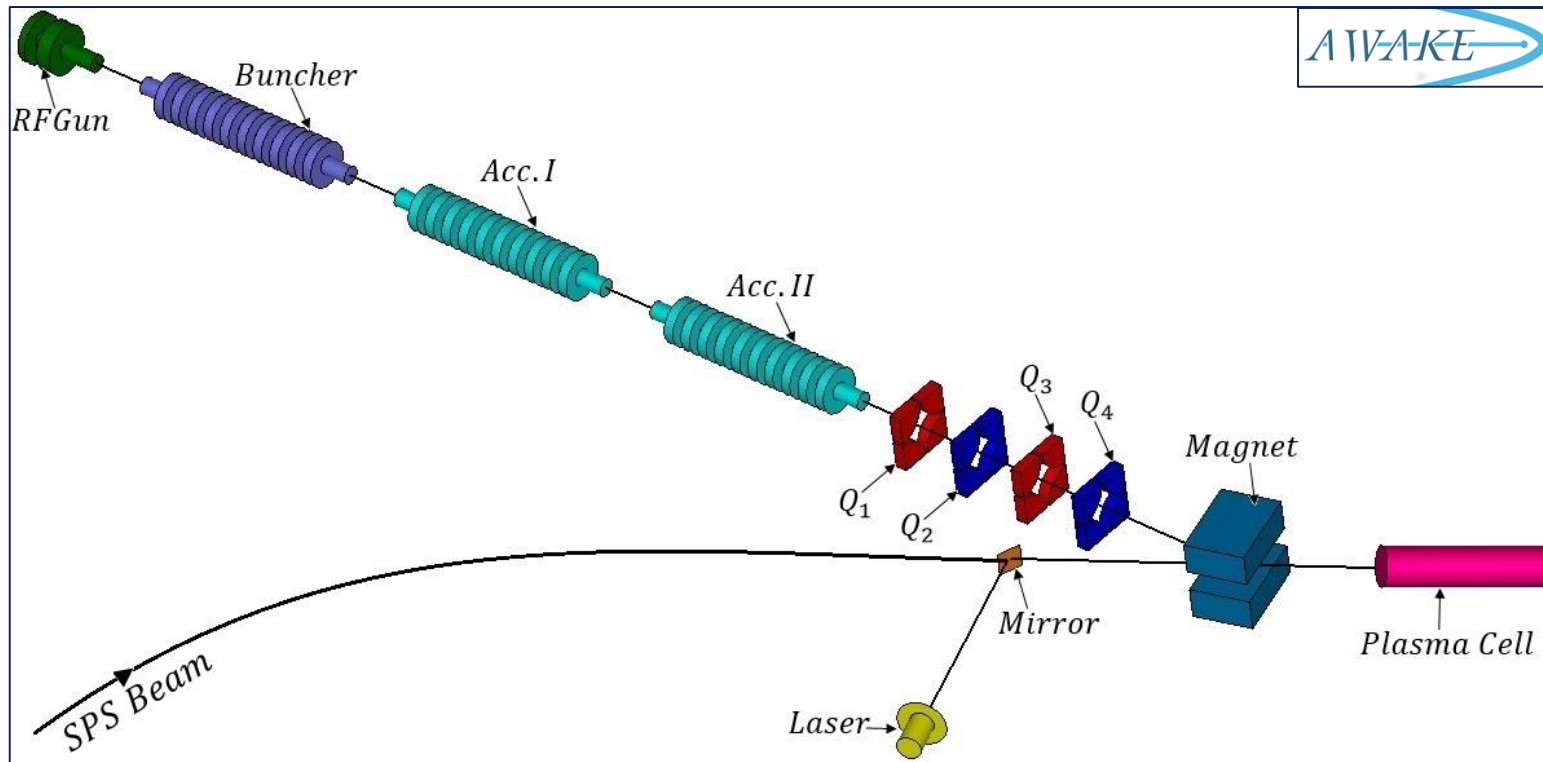
Parameter	Cell	f[GHz]	ϕ [rad]	E_0 [Mv/m]
Value	31	3	$2\pi/3$	15

ASTRA



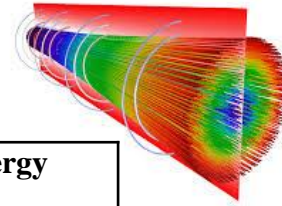
3.5 AWAKE Experiment

Injection of a **compact** and **high-quality** electron bunch at a **right phase** allows for a propagation over long distances with preserving emittance.



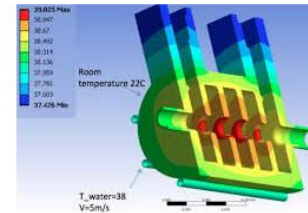
3.5 AWAKE Experiment

Beam Characteristics



Type	Bunch Charge	Bunch Length	Energy Spread	Emittance	Length	Energy
2nd	$\cong 200 \text{ pC}$	$< 200 \text{ fs}$	$< 1\%$	$< 2 \text{ } \mu\text{m}$	$< 5 \text{ m}$	$\cong 200 \text{ MeV}$

Parameter	RF Gun	Buncher	Acc. I	Acc. II
Frequency	3.0	12.0	12.0	12.0
Max Gradient	120MV/m	50MV/m	80MV/m	80MV/m
N. Cell	1.5	30	120	120



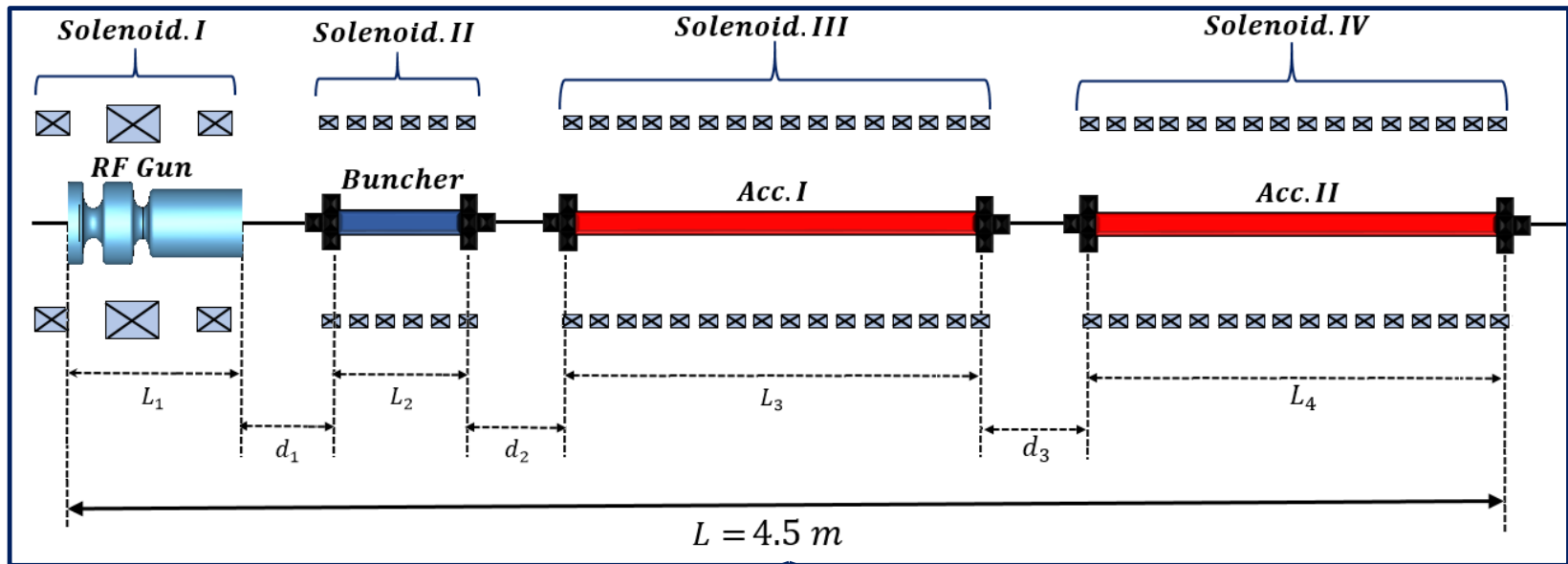
RF Characteristics

Laser Characteristics

$\lambda[\text{nm}]$	$w[\text{ev}]$	$r[\text{mm}]$	$t[\text{ps}]$	$q[\text{pc}]$
262	4.31	1.0-2.0	1.0	100-600



3.5 AWAKE Experiment



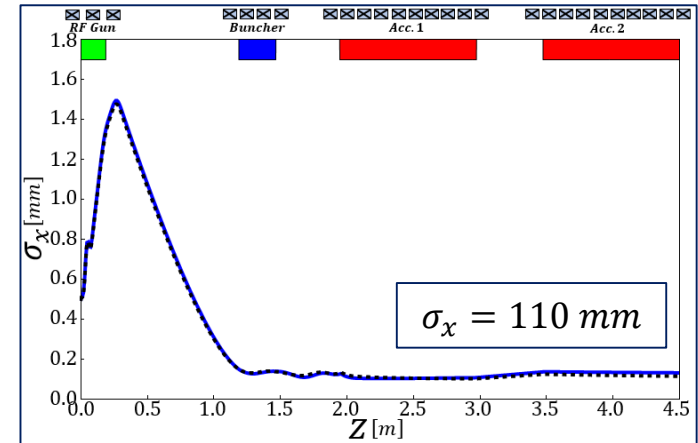
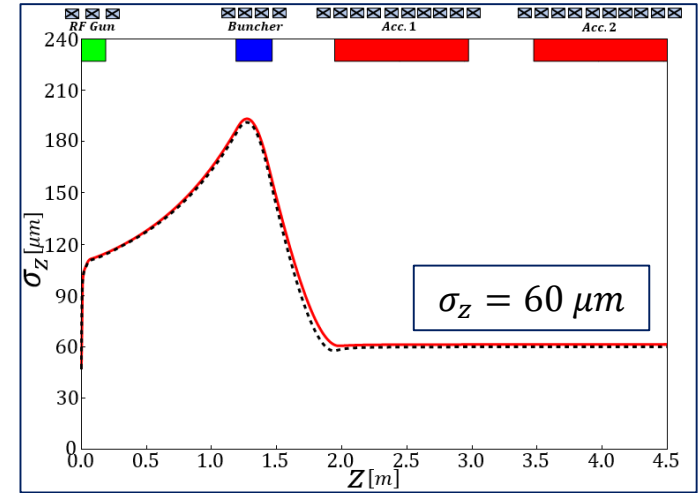
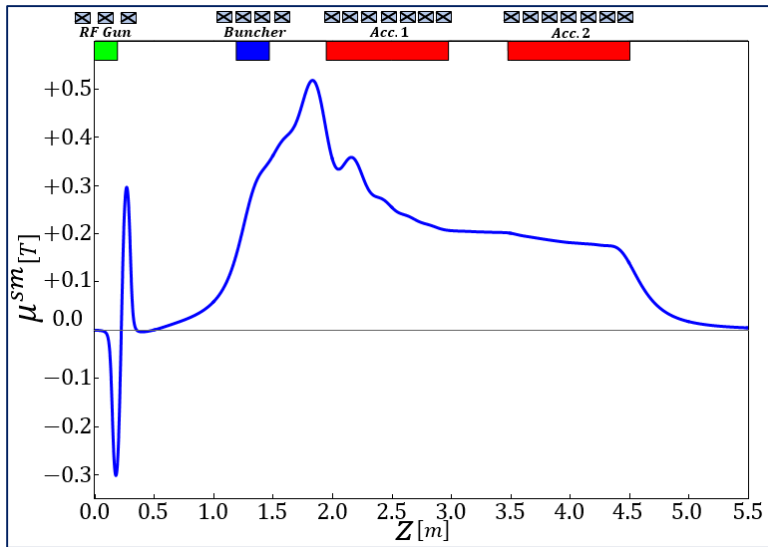
S Band RF Gun Very High Quality EBeam Generation

X Band Buncher for Very Strong Bunching

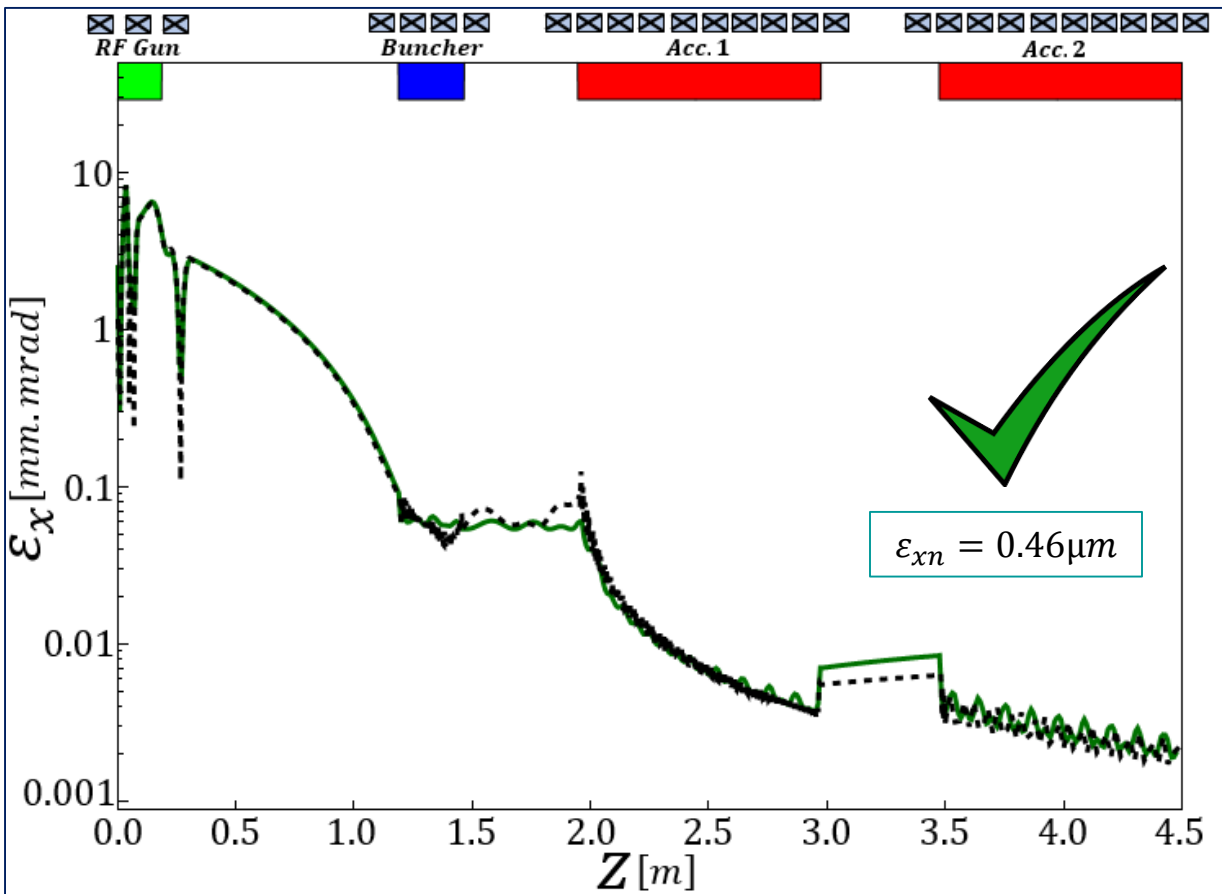
X Band Acc for Very Compact Acceleration

3.5 AWAKE Experiment

Parameter	$d_1[cm]$	$d_2[cm]$	$d_3[cm]$	$E_2[mV/m]$	$E_3[mV/m]$
Gradient	100	48	50	32	80



3.5 AWAKE Experiment



Six-Dimensional Beam-Envelope Equations: An Ultrafast Computational Approach for Interactive Modeling of Accelerator Structures

M.D. Kelisani^{1,2,*}, S. Barzegar², P. Craievich³, and S. Doebert¹

¹BE-RF Department, European Organization for Nuclear Research (CERN), Geneva, CH-1211 Switzerland

²Institute for Research in Fundamental Sciences (IPM), P.O. Box 19395-5531, Tehran, Iran

³Paul Scherrer Institute (PSI), Villigen, CH-5232 Switzerland

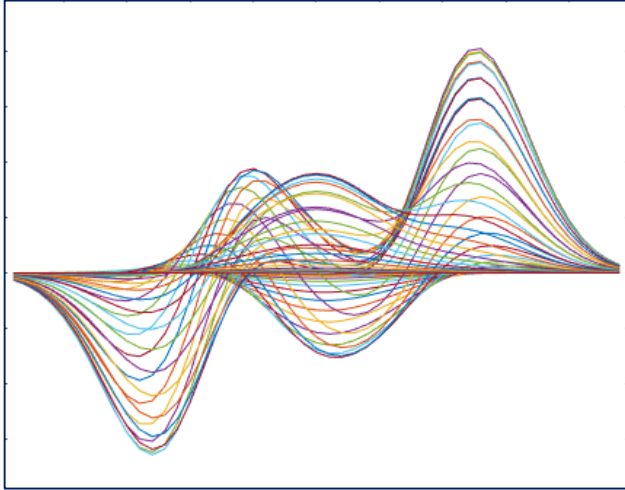
(Received 3 October 2022; revised 1 March 2023; accepted 3 April 2023; published 3 May 2023)

The design and implementation of accelerators capable of providing high-quality bunches require precise and efficient online modeling tools. Current comprehensive beam dynamics studies are prohibitively

Thanks for Attention



2.2 Perturbation



Perturbational Methods

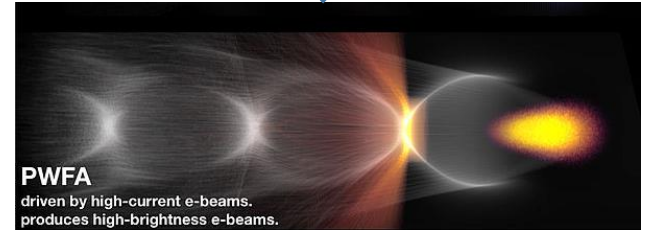
$$\beta_u = \beta_0(\hat{u} \cdot \hat{z} + \Delta u')$$

$$\frac{1}{\gamma} \cong \frac{1}{\gamma_0} \left\{ 1 - p_0^2 \Delta z' - \frac{p_0^2 \gamma_0^2}{2} \Delta z'^2 - \frac{p_0^4 \gamma_0^2}{2} \Delta z'^3 - \frac{p_0^2}{2} (1 + p_0^2 \Delta z') (\Delta x'^2 + \Delta y'^2) \right\}$$

$$\frac{1 - \beta_z \beta_0}{\gamma} \cong \frac{1}{\gamma_0^3} \left(1 - 2p_0^2 \Delta z' - \frac{p_0^2 (1 - p_0^2)}{2} \Delta z'^2 - \frac{p_0^2}{2} (\Delta x'^2 + \Delta y'^2) \right)$$

$$\frac{\beta_0 - \beta_z}{\gamma} \cong \frac{\beta_0}{\gamma_0} \left(-\Delta z' + p_0^2 \Delta z'^2 + \frac{p_0^2 \gamma_0^2}{2} \Delta z'^3 + \frac{p_0^2}{2} \Delta z' (\Delta x'^2 + \Delta y'^2) \right)$$

3.5 AWAKE Experiment



Run 2c: Demonstrate Electron Acceleration and Emittance Preservation

