

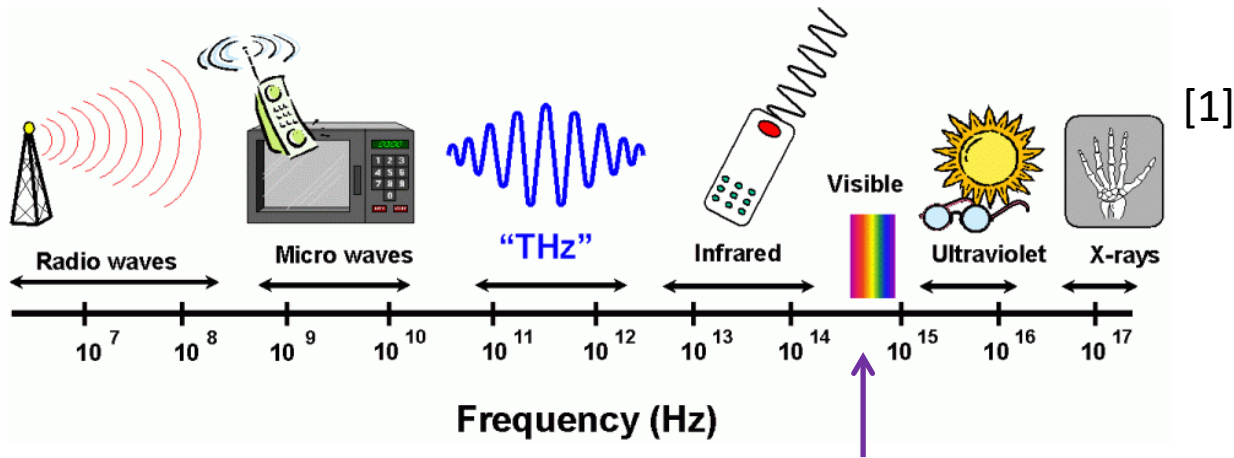
Generation and Detection of coherently emitted THz radiation at DELTA

Carsten Mai

Center for Synchrotron Radiation
TU Dortmund



The Terahertz domain



$$1 / 1 \text{ THz} = 1 \text{ ps}$$

$$c_0 / 1 \text{ THz} = 0.3 \text{ mm}$$

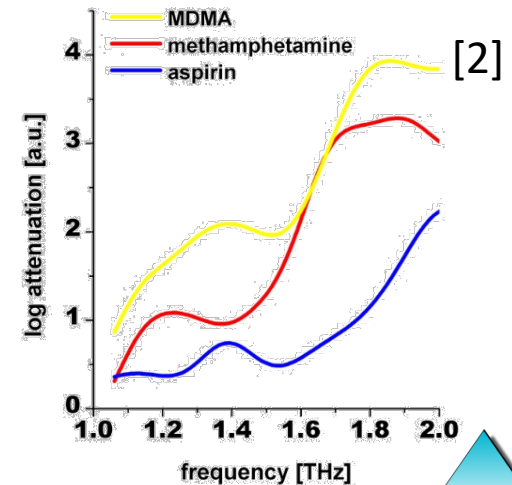
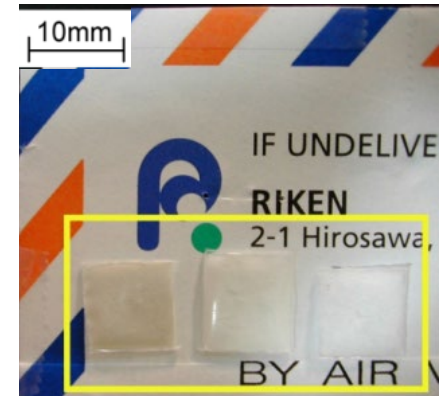
$$800 \text{ nm} \approx c_0 / 375 \text{ THz}$$

- THz gap
- limited access to sources and detectors
- non-destructive material science
- THz radiation is part of synchrotron radiation spectrum

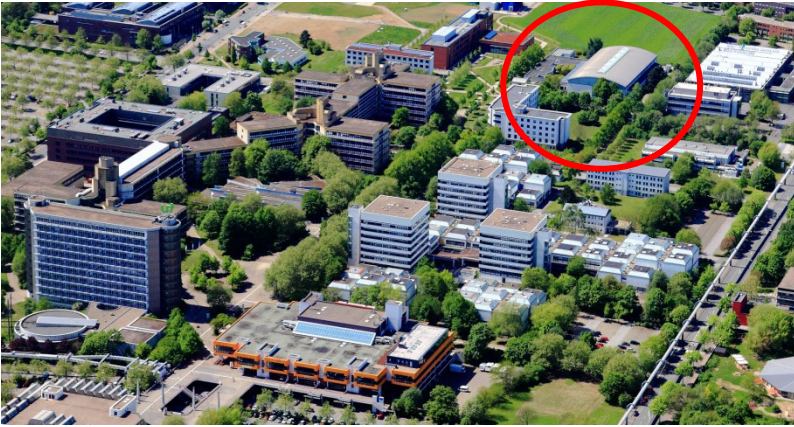
[1] <http://www.terasense.com>

[2] K. Kawase et. al, Optics Express 11, 2549 (2003)

[2]



DELTA – the light source at TU Dortmund University

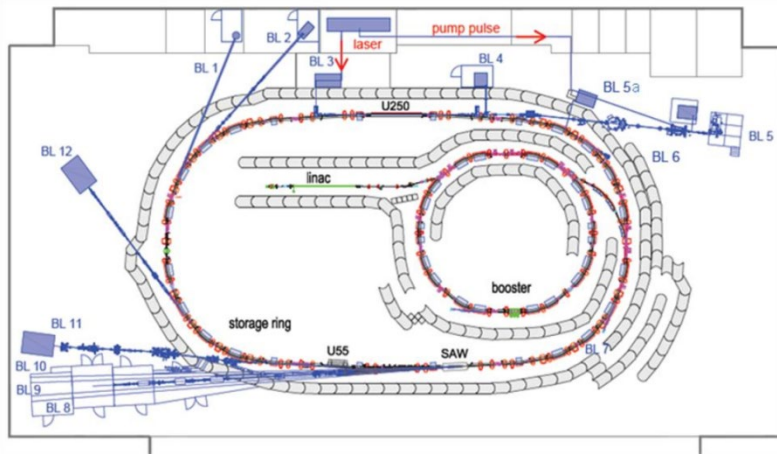


Parameters

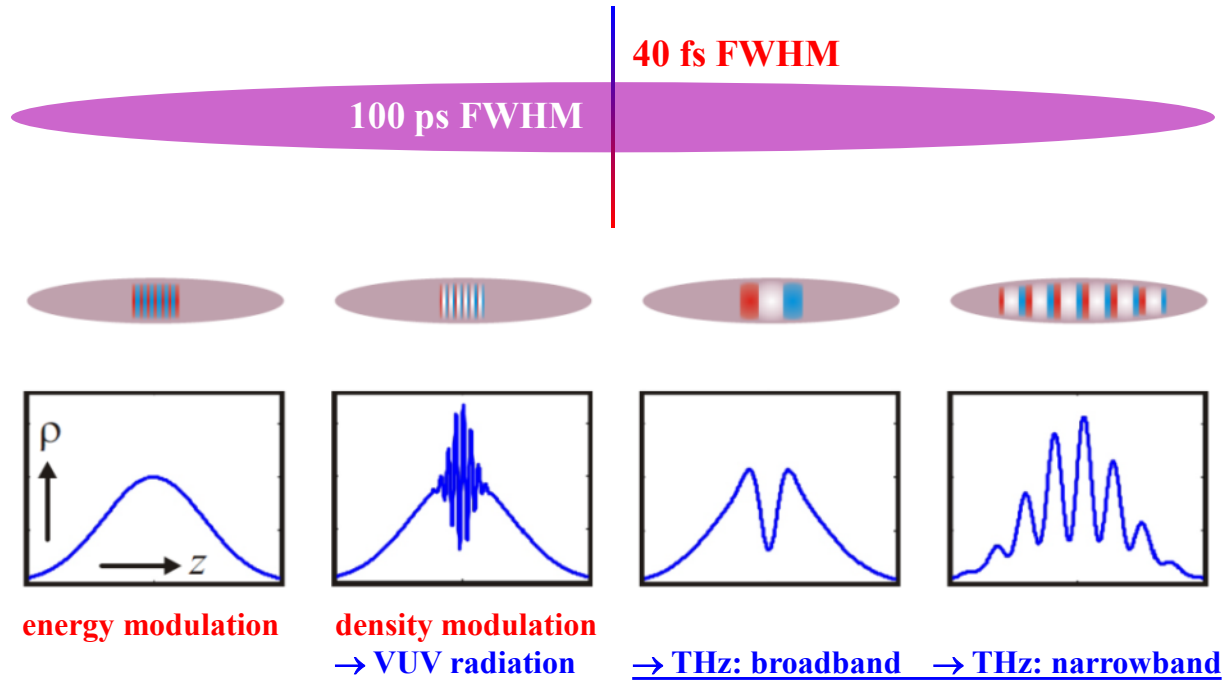
circumference:	115.2 m
beam energy:	1.5 GeV
beam current:	130 mA multi-bunch 20 mA single bunch
beam lifetime:	~15 h @ 100 mA
emittance:	~16 nm rad (horiz.)
bunch length:	100 ps (FWHM)

Operation times

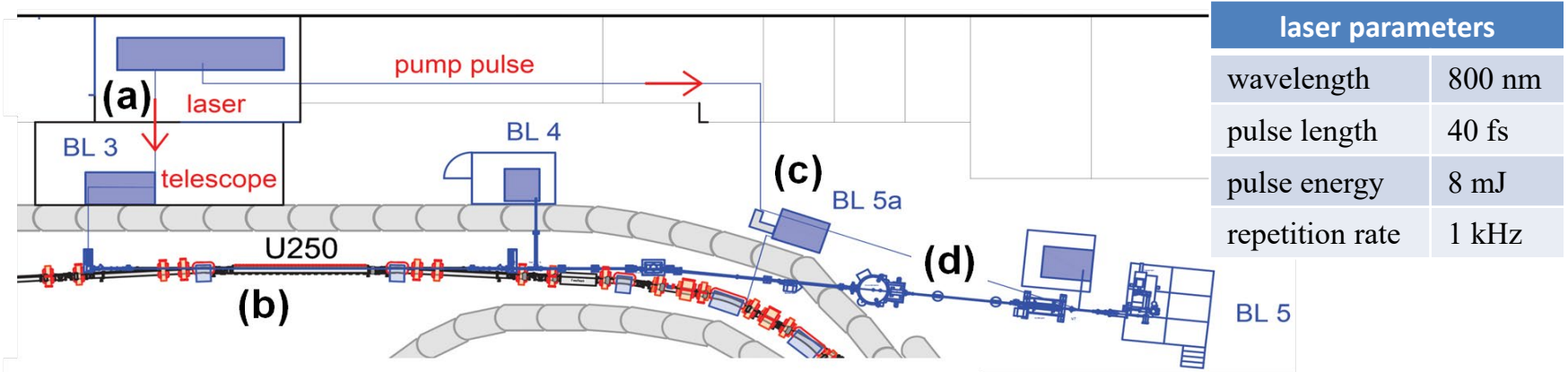
user experiments:	2000 h / year
machine studies:	1000+ h / year



Seeding schemes at DELTA



The DELTA short-pulse facility



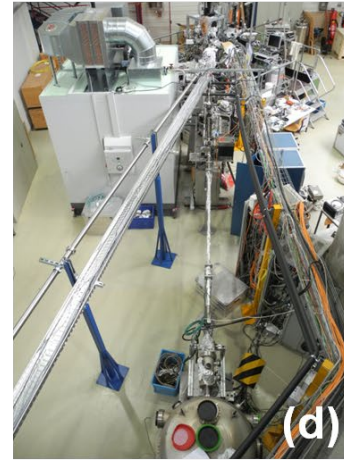
fs laser system



undulator U250

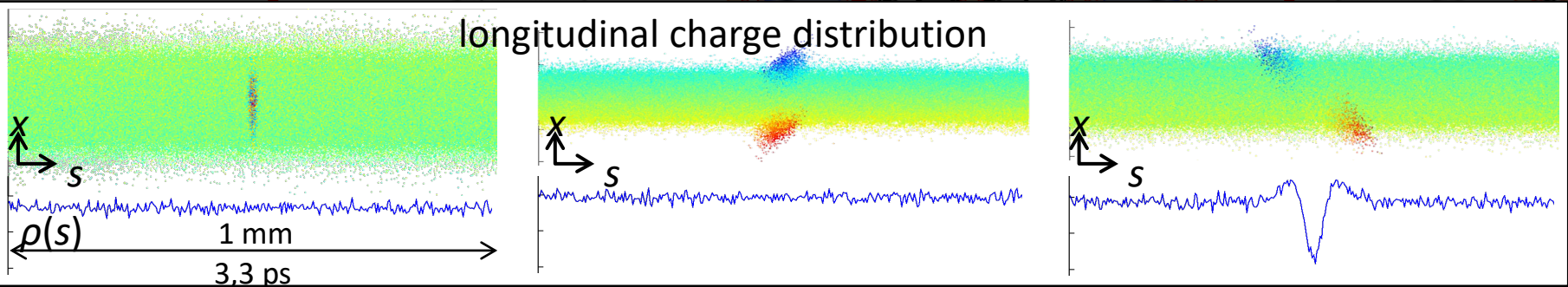
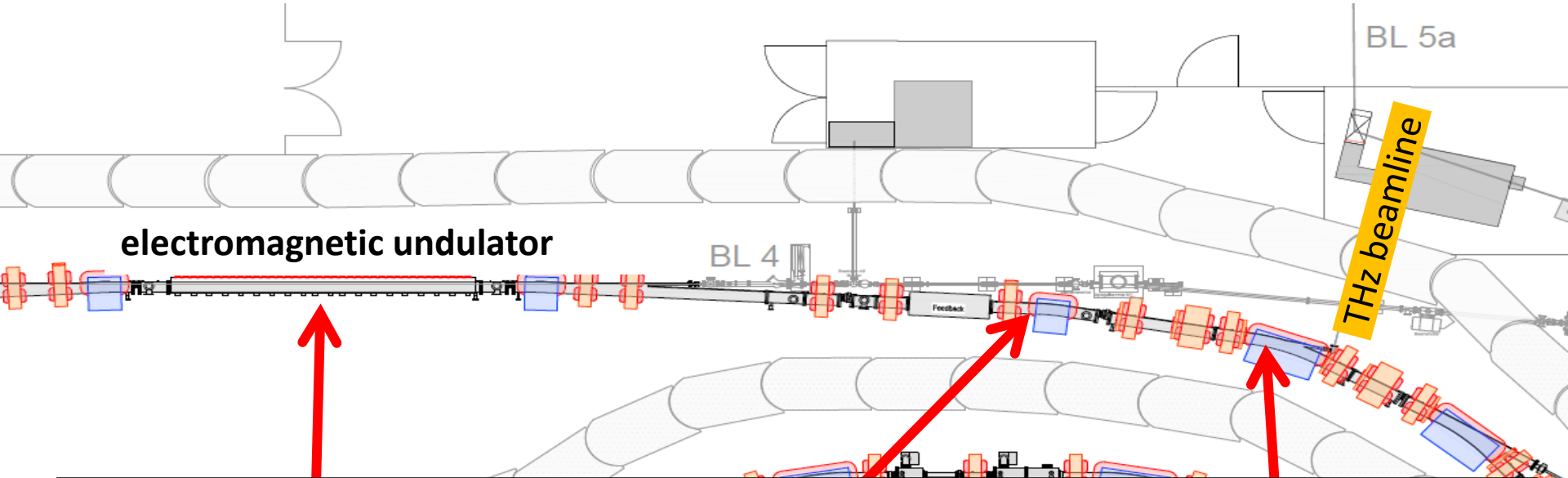


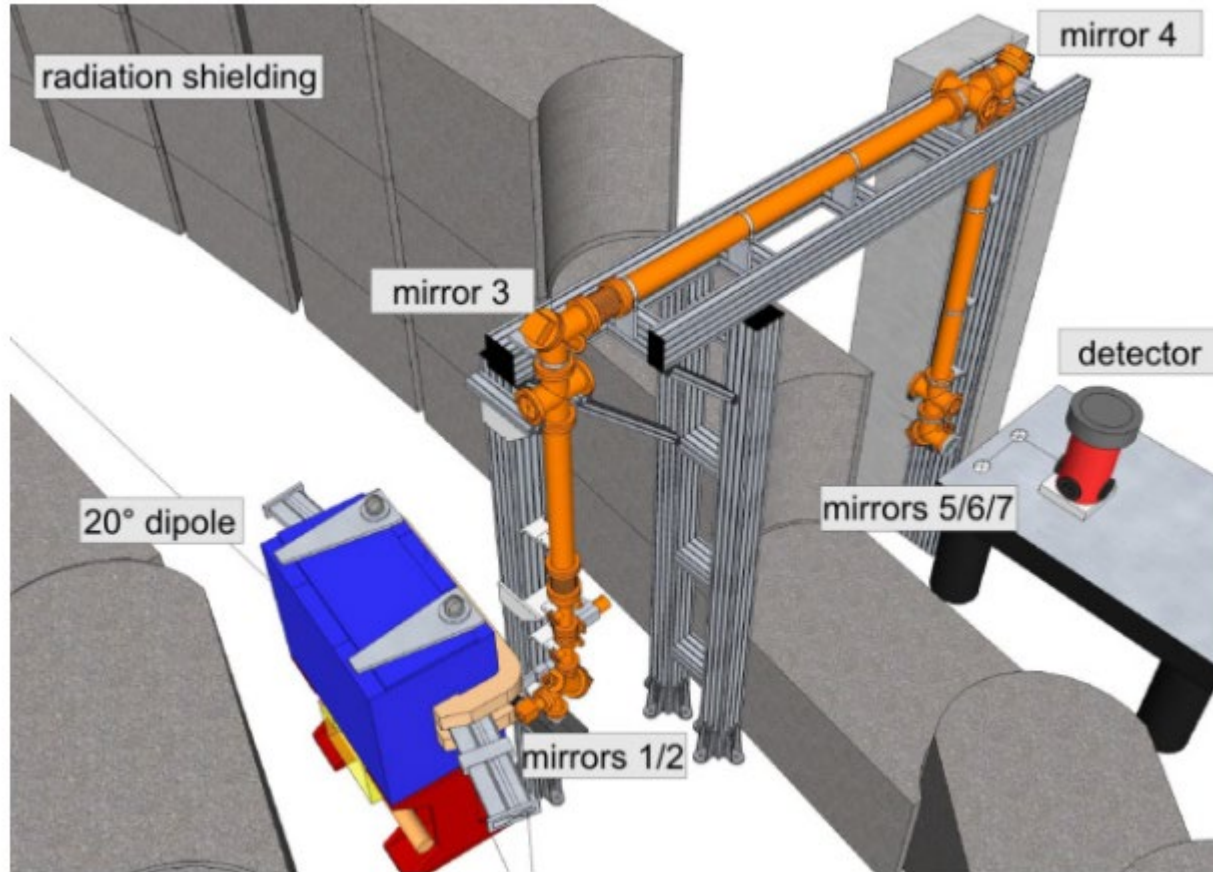
THz beamline

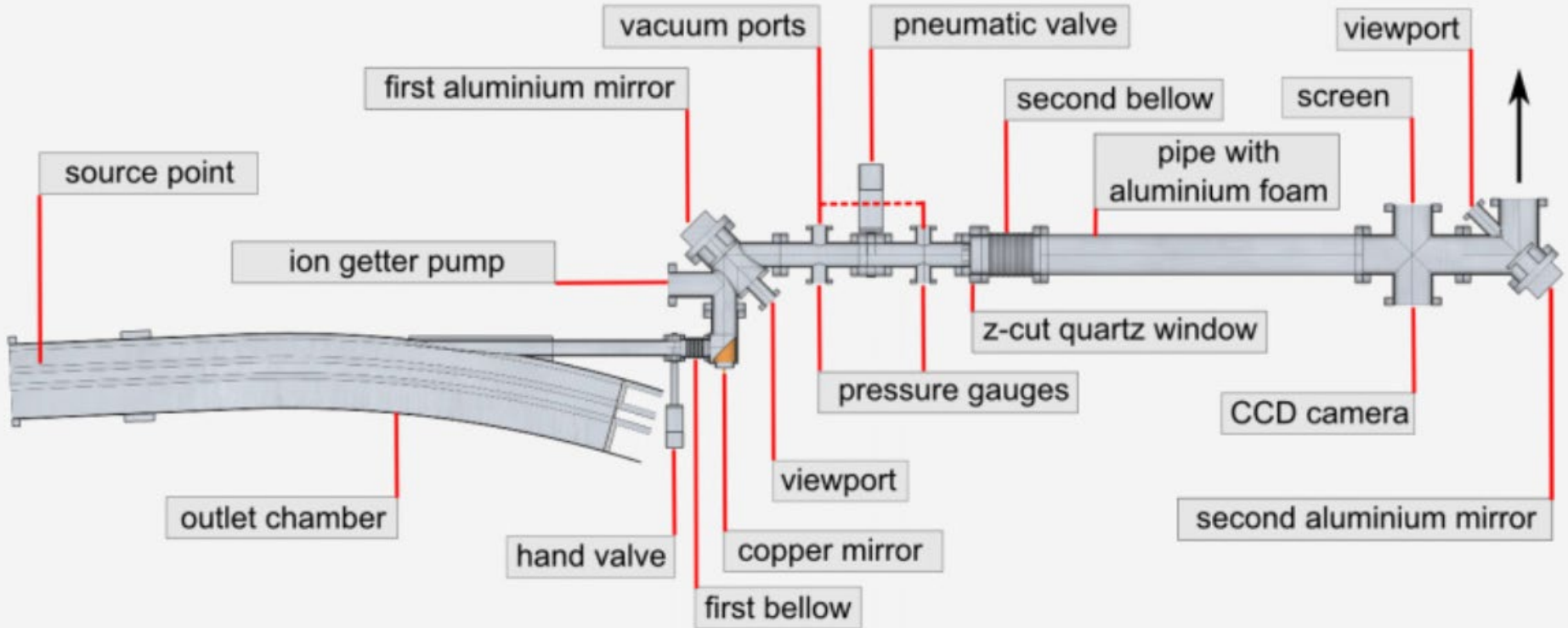


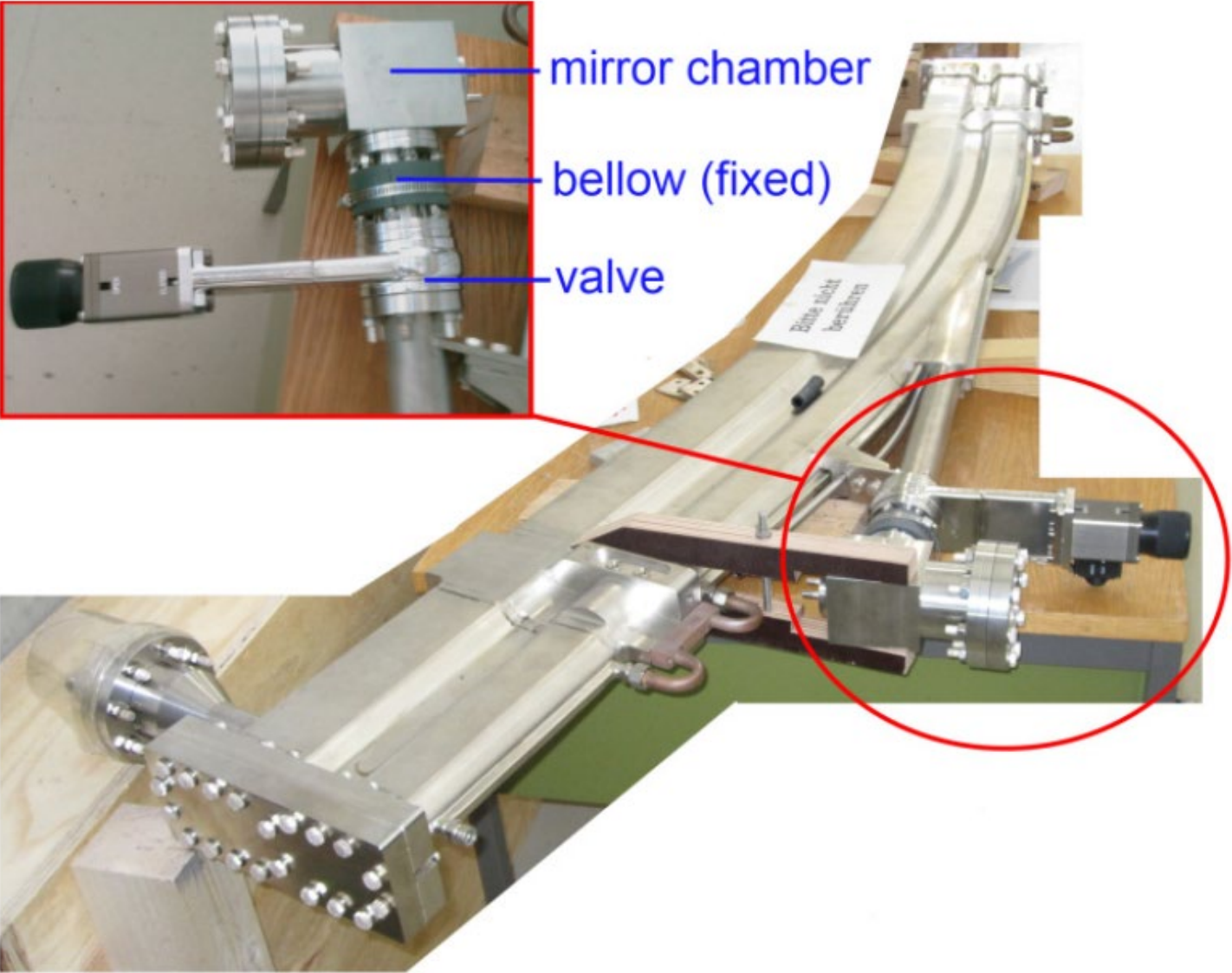
VUV beamline for pump-probe experiments

Coherent emission of THz radiation

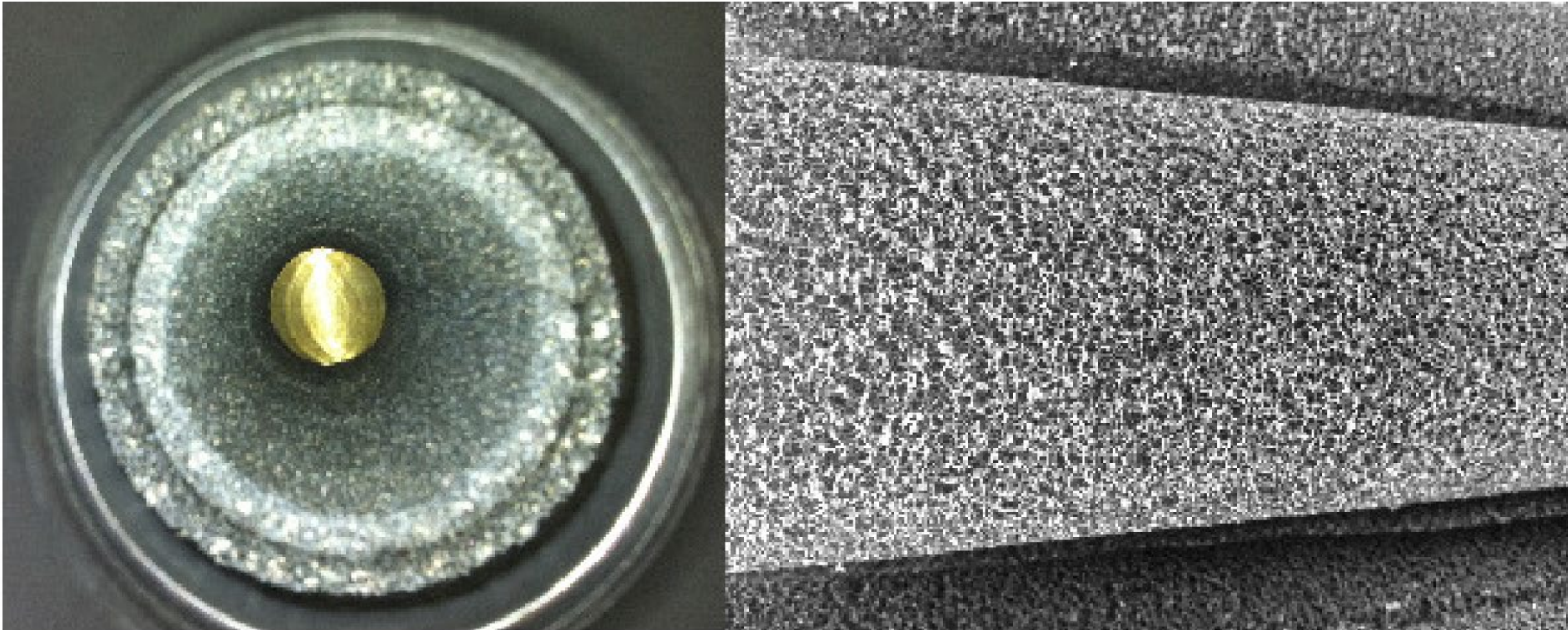






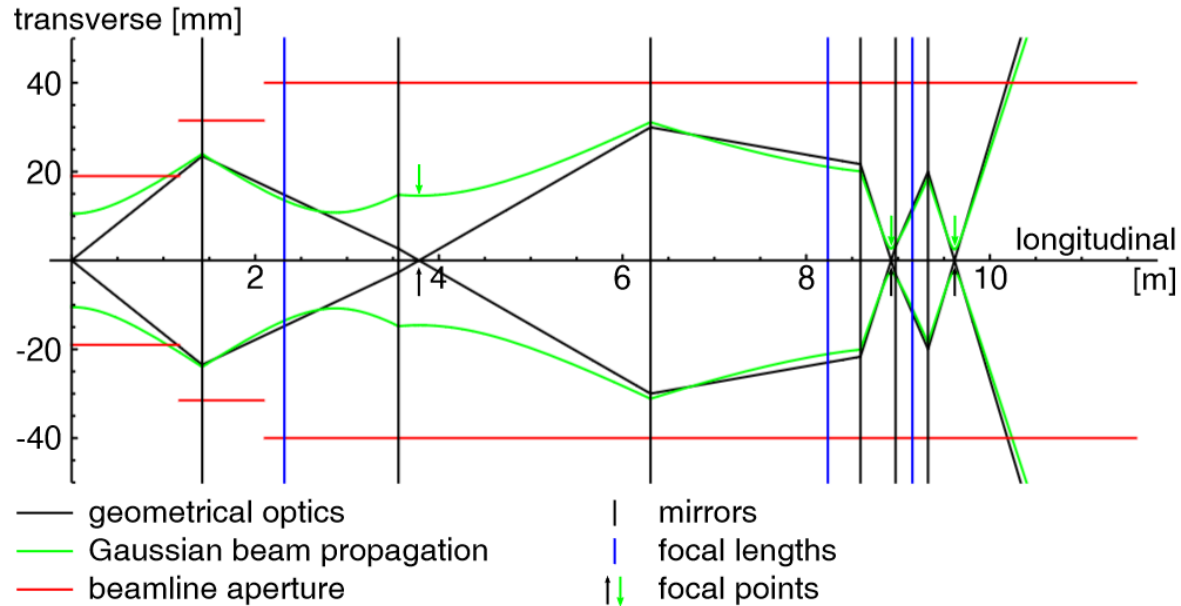


Reflectionless Beam transport

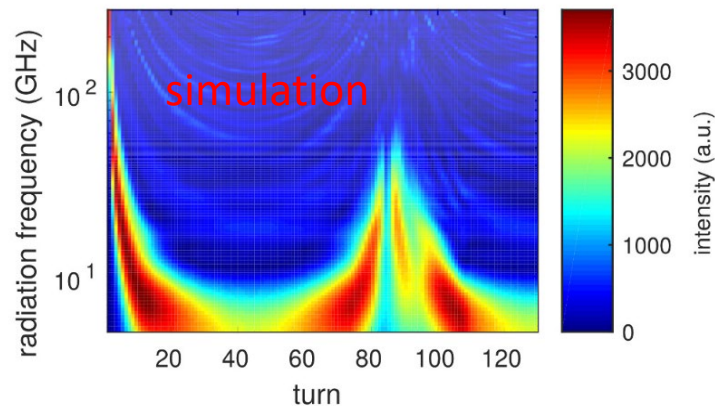
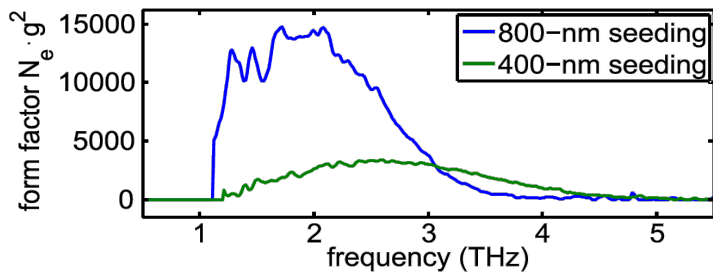
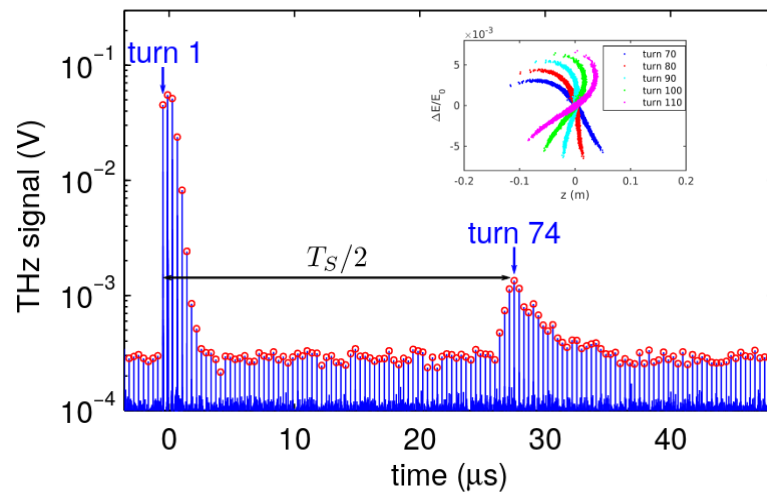
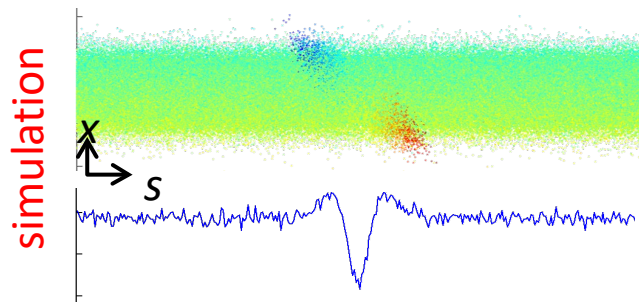




Beam propagation

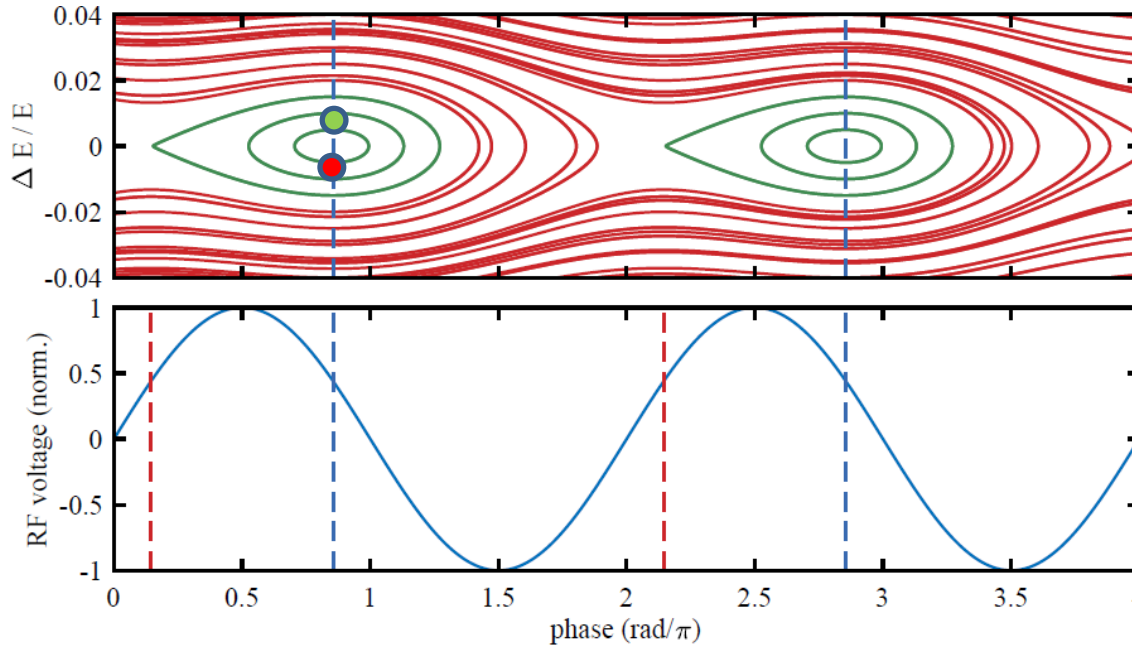
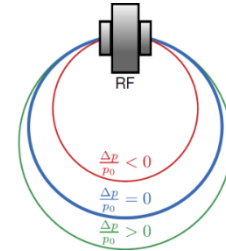
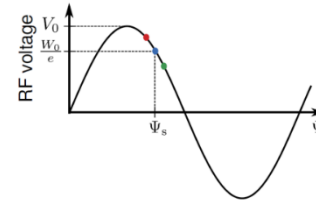


Broadband THz generation: spectrotemporal evolution



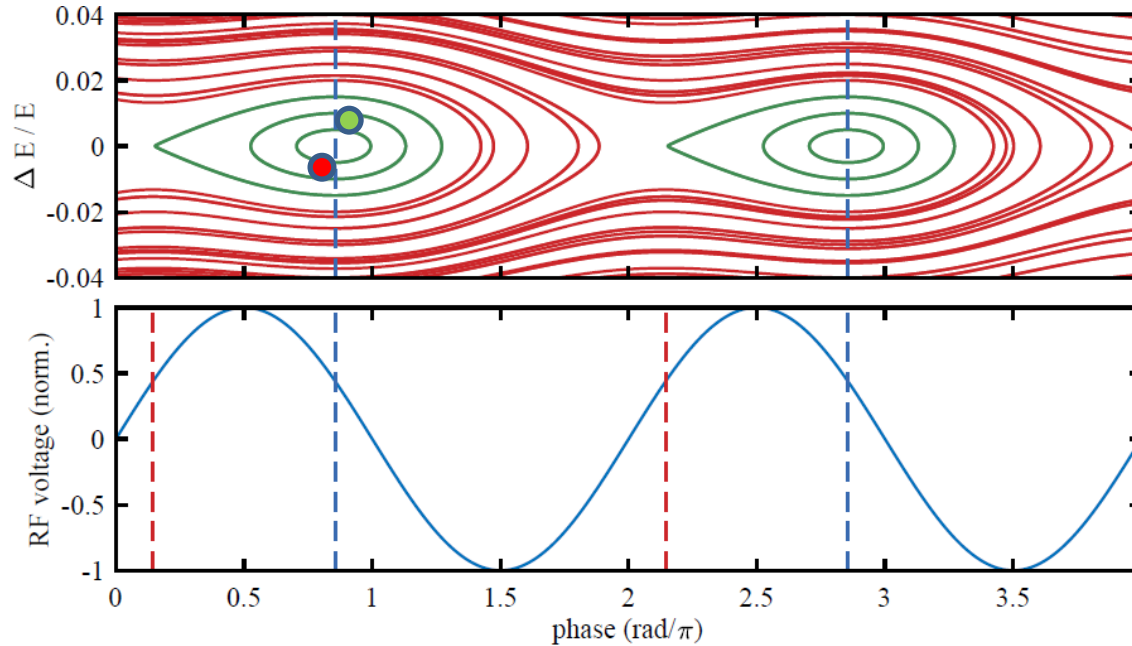
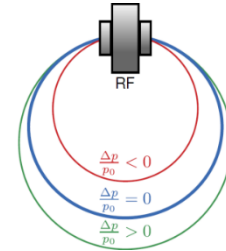
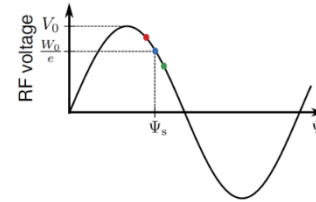
synchrotron frequency:

$$\Omega^2 = \omega_{\text{rev}}^2 \left(\frac{\alpha_c h}{2\pi E_0} \right) e \frac{dV}{d\psi}(\psi_s)$$



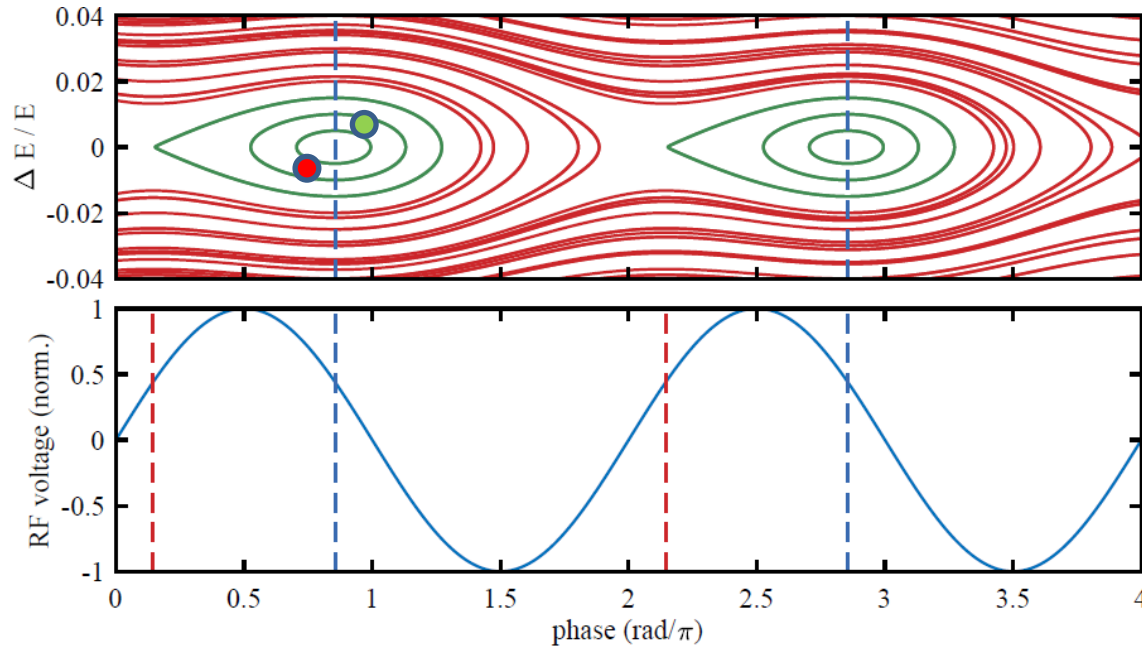
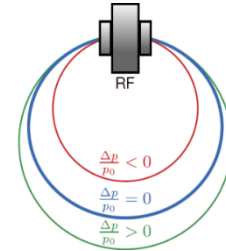
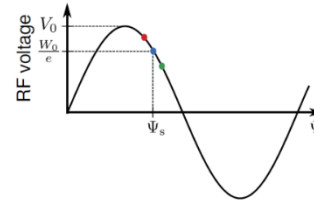
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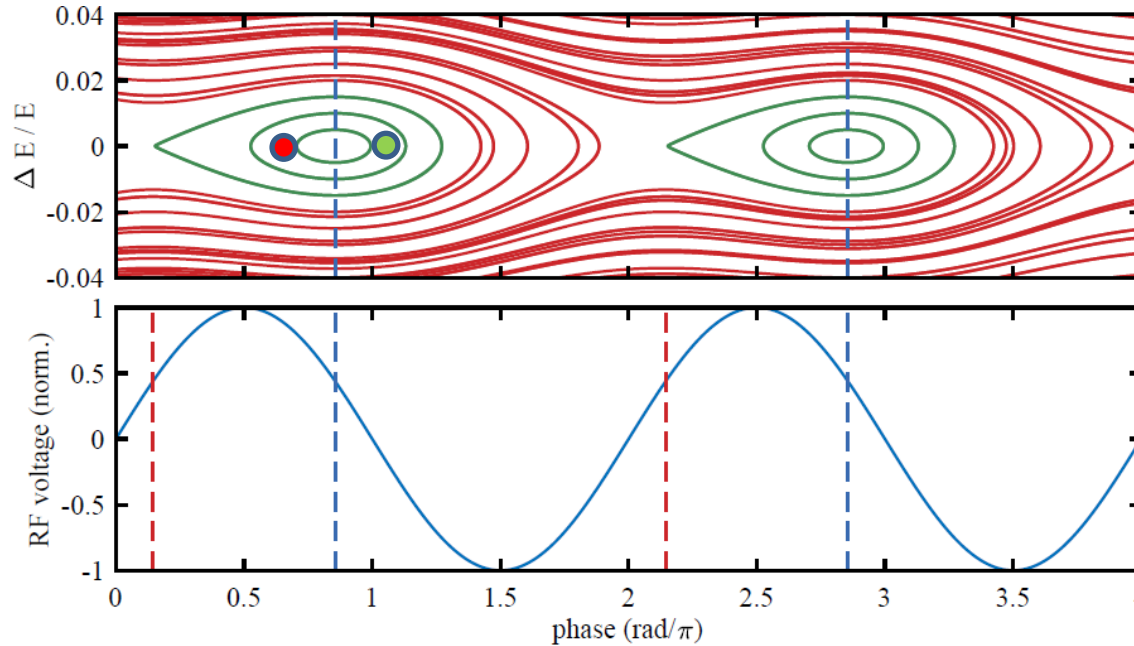
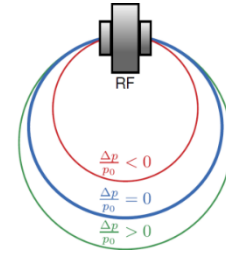
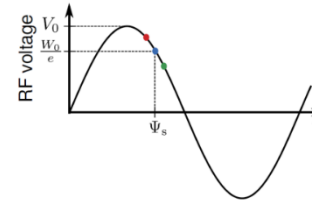
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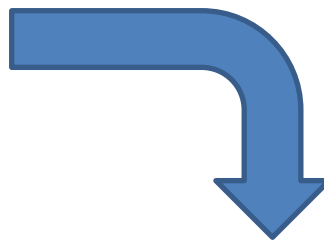
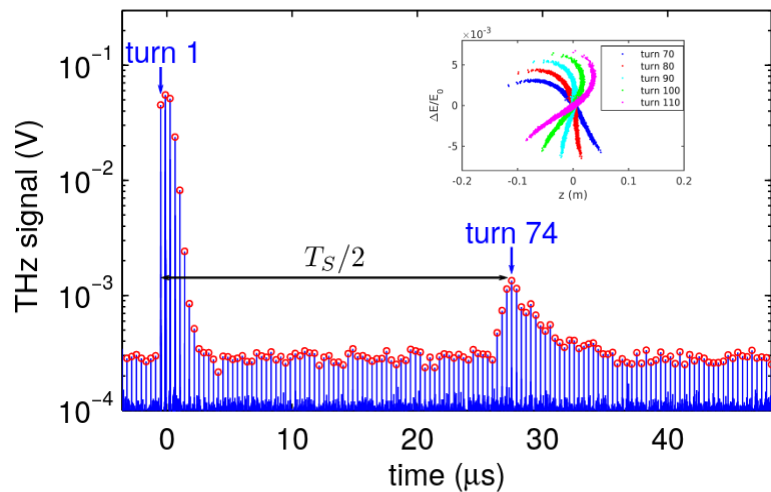


synchrotron frequency:

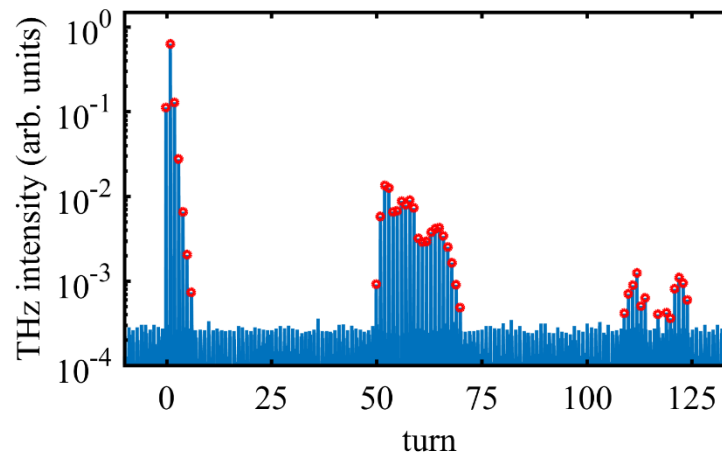
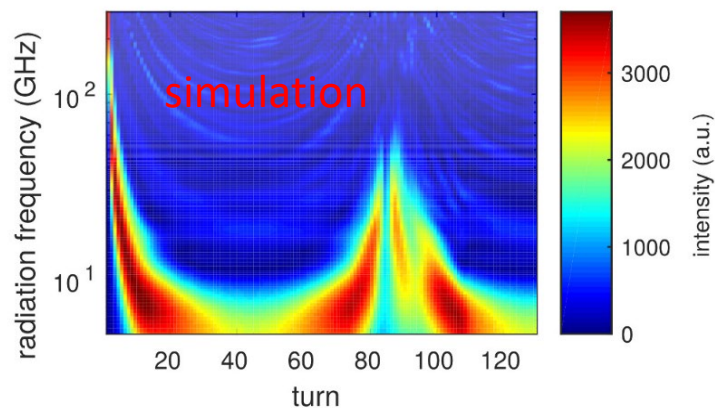
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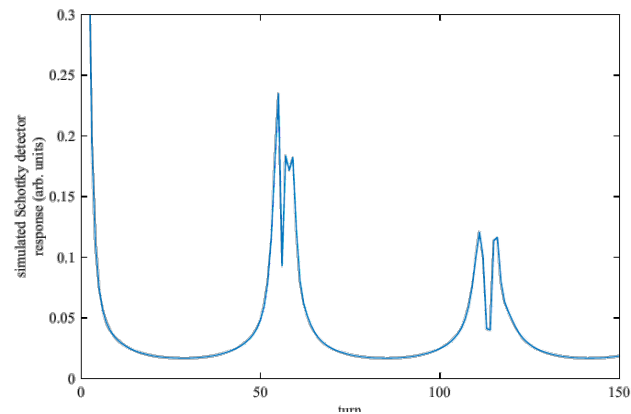
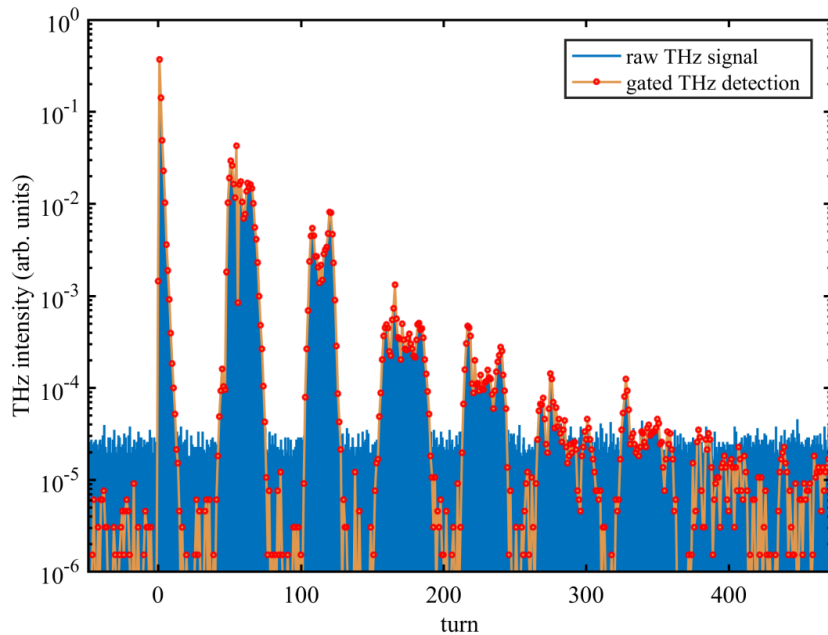
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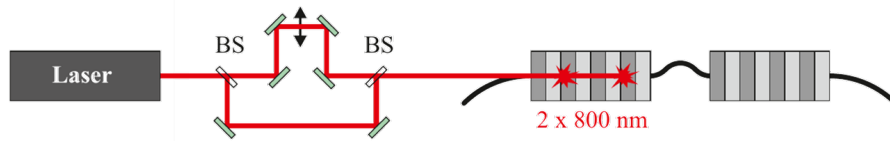
- RF Upgrade (second cavity)
- improved detector readout



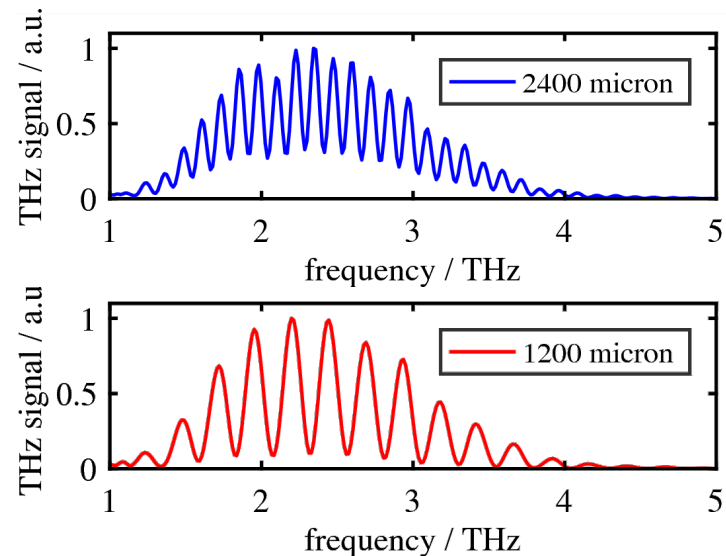
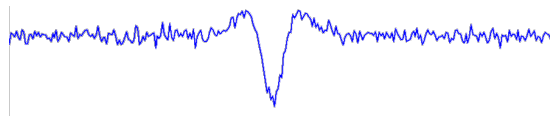
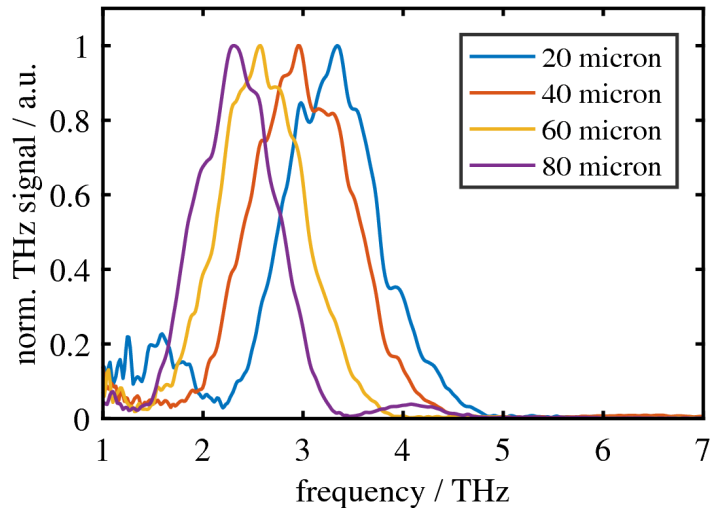
More accurate measurement: gated detection



Increasing spectral control of THz generation

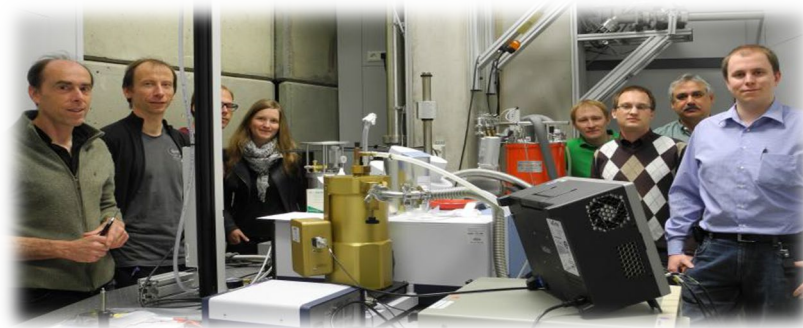


Step 1: replicate textbook knowledge...

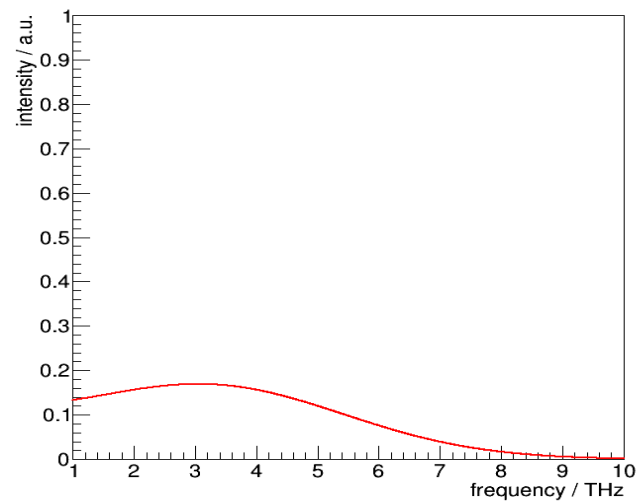
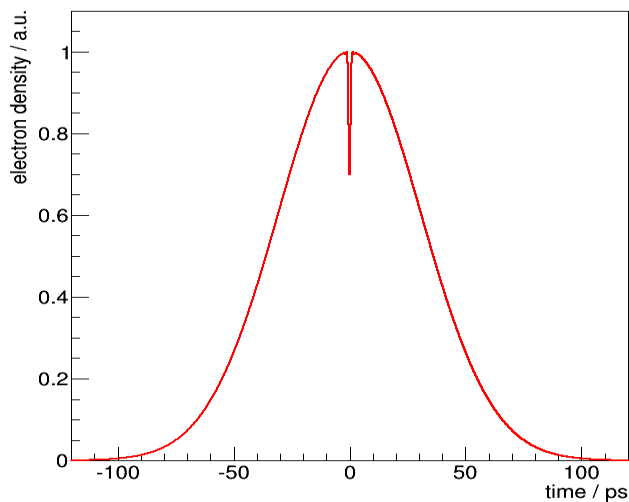


Towards tunable narrowband radiation

cooperation with PhLAM, Lille

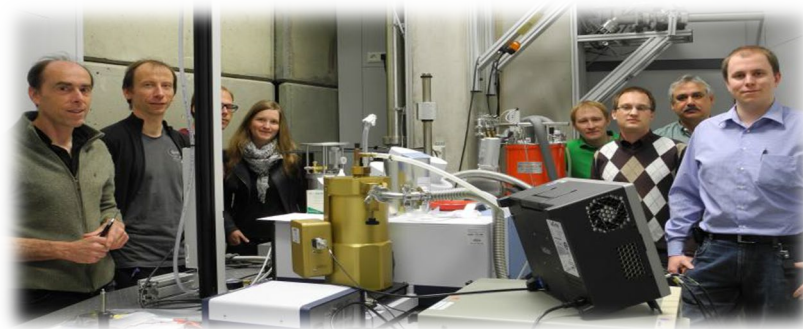


simple model:



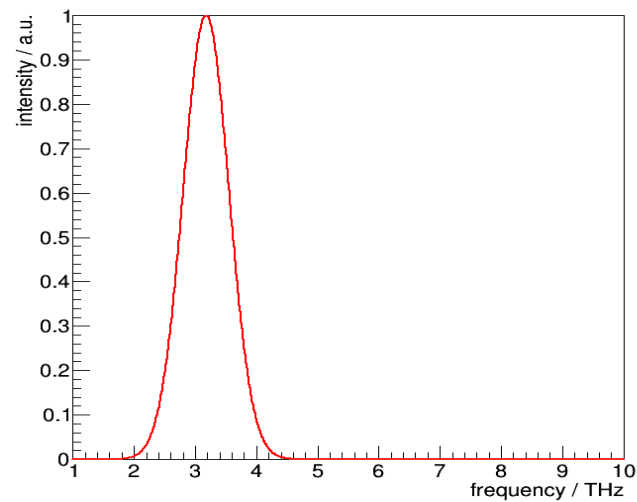
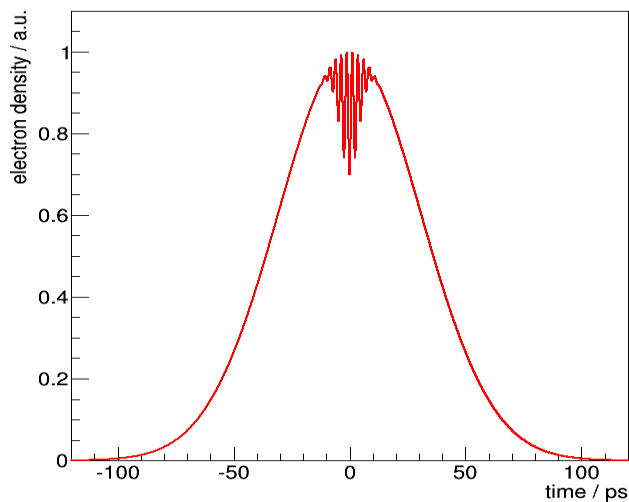
Towards tunable narrowband radiation

cooperation with PhLAM, Lille



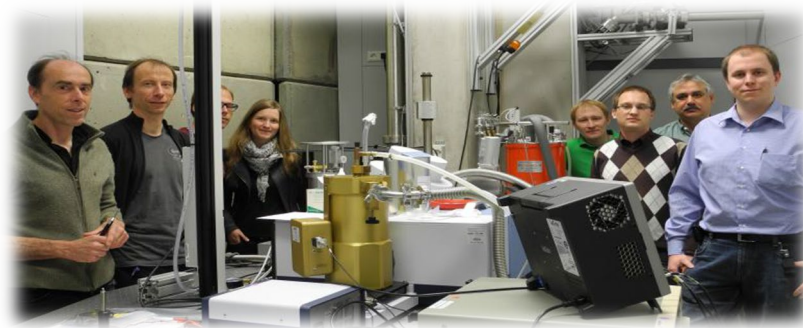
- multi-dip modulation of the electron bunch leads to narrow THz spectrum
- idea: modulate long, chirped laser pulse with Michelson interferometer
- first realized at UVSOR

simple model:

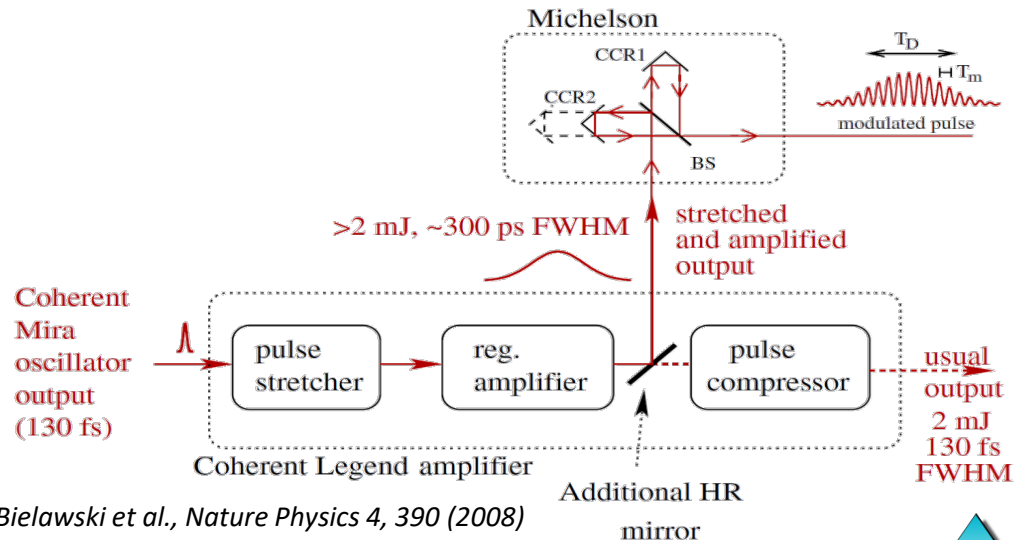
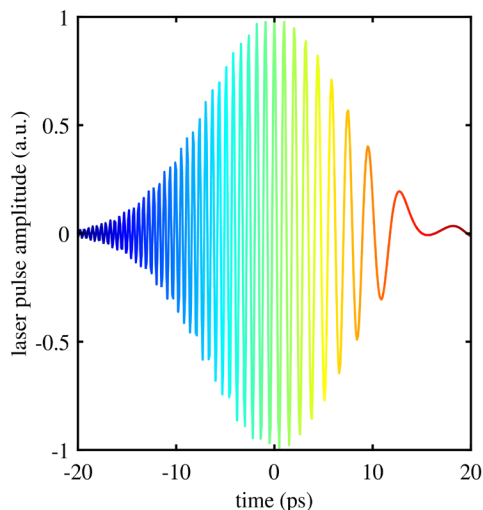


Towards tunable narrowband radiation

cooperation with PhLAM, Lille



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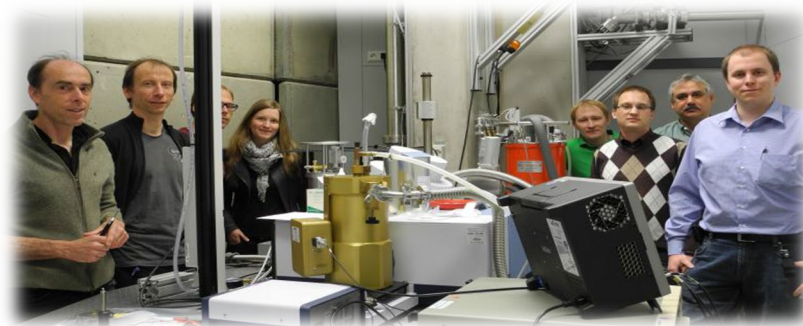


S. Bielawski et al., Nature Physics 4, 390 (2008)

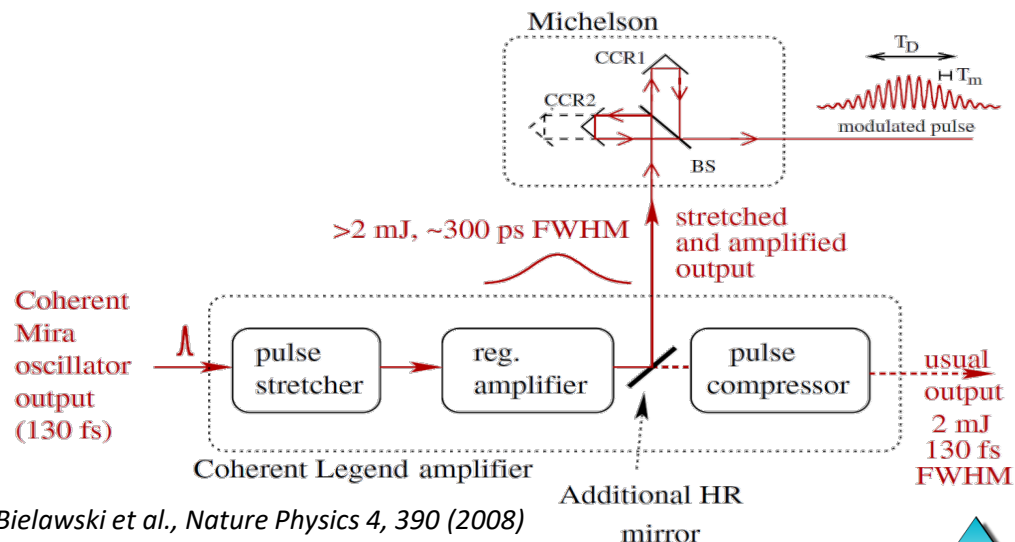
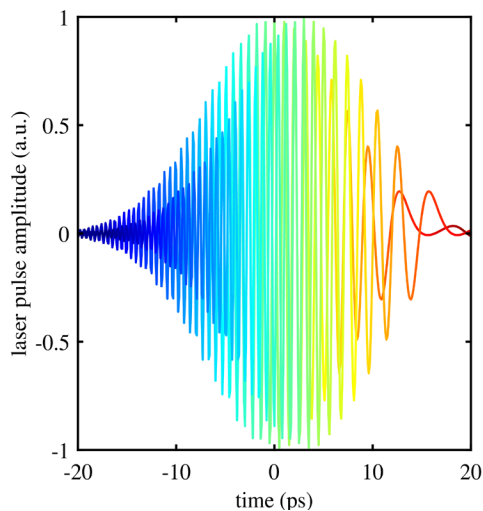
C. Evain et al., PRST-AB 13, 090703 (2010)

Towards tunable narrowband radiation

cooperation with PhLAM, Lille



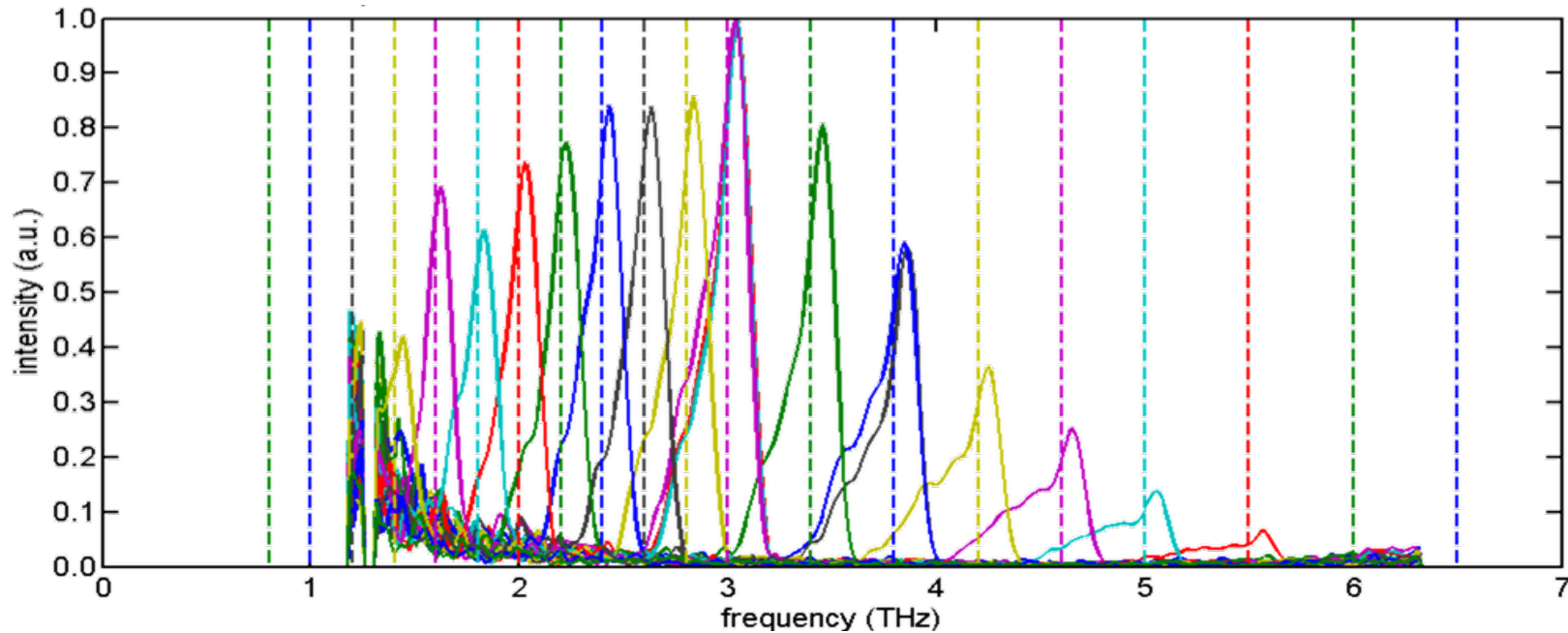
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Tunable THz radiation



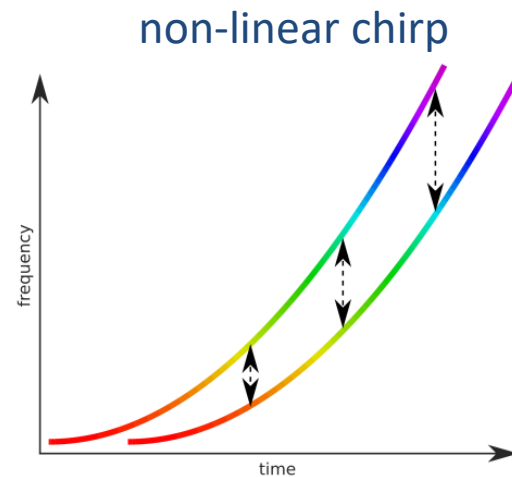
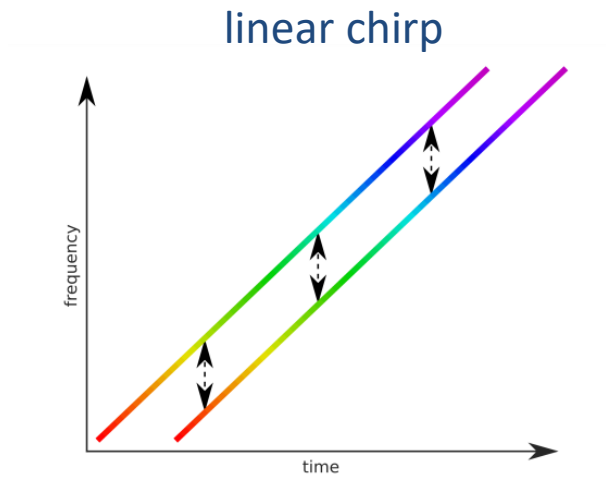
- selectable target frequency (dashed)
- higher order chirp introduces spectral broadening
- asymmetric shape

Higher Order Dispersion

- Taylor expansion of the optical phase:

$$\phi(\omega) = D_0 + D_1 \cdot (\omega - \omega_0) + D_2 \cdot (\omega - \omega_0)^2 + D_3 \cdot (\omega - \omega_0)^3 + \dots$$

D_2 : linear chirp of the pulse



Dispersion measurement

$$\phi(\omega) = D_0 + D_1 \cdot (\omega - \omega_0) + D_2 \cdot (\omega - \omega_0)^2 + D_3 \cdot (\omega - \omega_0)^3 + \dots$$

- D_2 and D_3 must be known

- FROG?



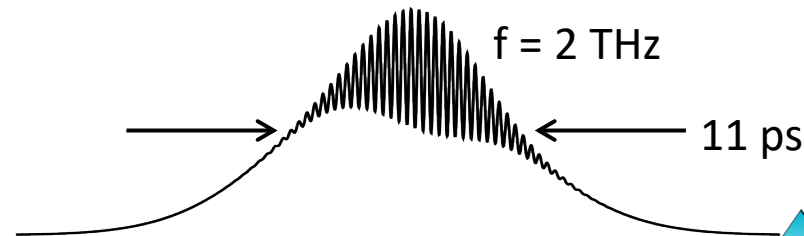
- SPIDER?



autocorrelator!

- pulses are too long for our usual diagnostics

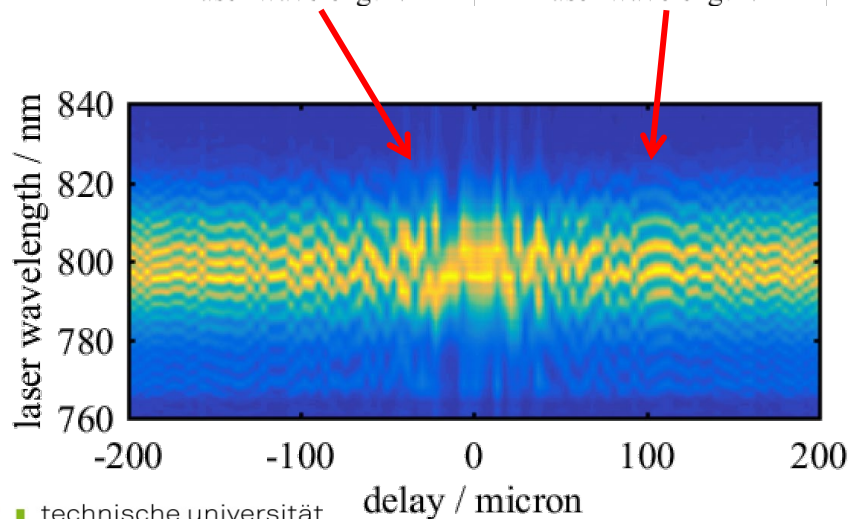
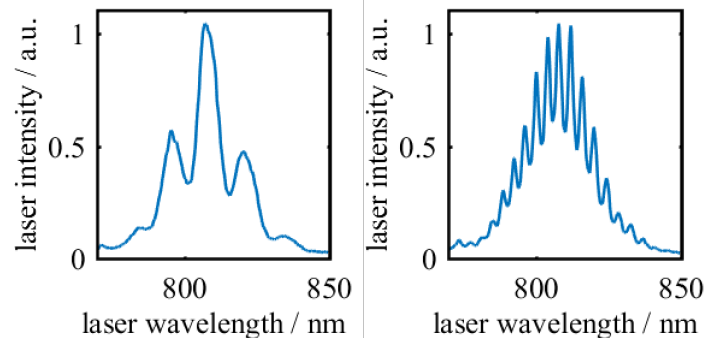
$$I(t) \sim \cos\left(\frac{\tau}{D_2}t + 3D_3\tau t^2\right)$$



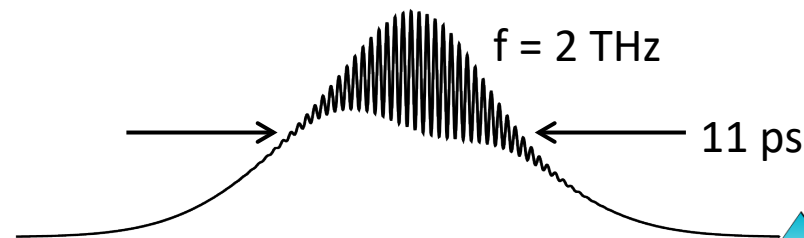
Dispersion measurement

$$\phi(\omega) = D_0 + D_1 \cdot (\omega - \omega_0) + D_2 \cdot (\omega - \omega_0)^2 + D_3 \cdot (\omega - \omega_0)^3 + \dots$$

spectra:

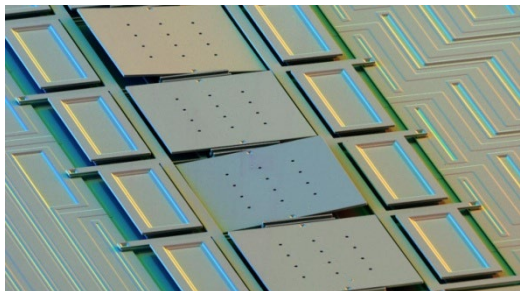


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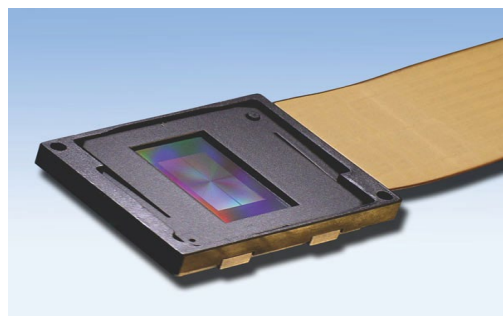


Adaptive optics for spatial light modulation

micro-mirror arrays (MEMS)

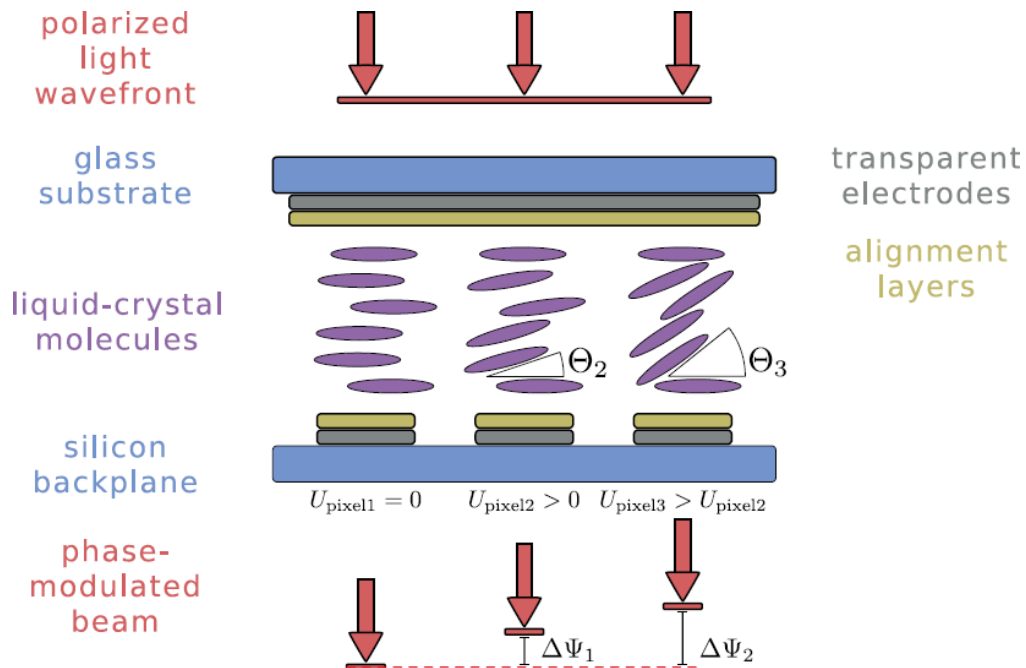


Texas Instruments



Holoeye AG

liquid-crystal modulators

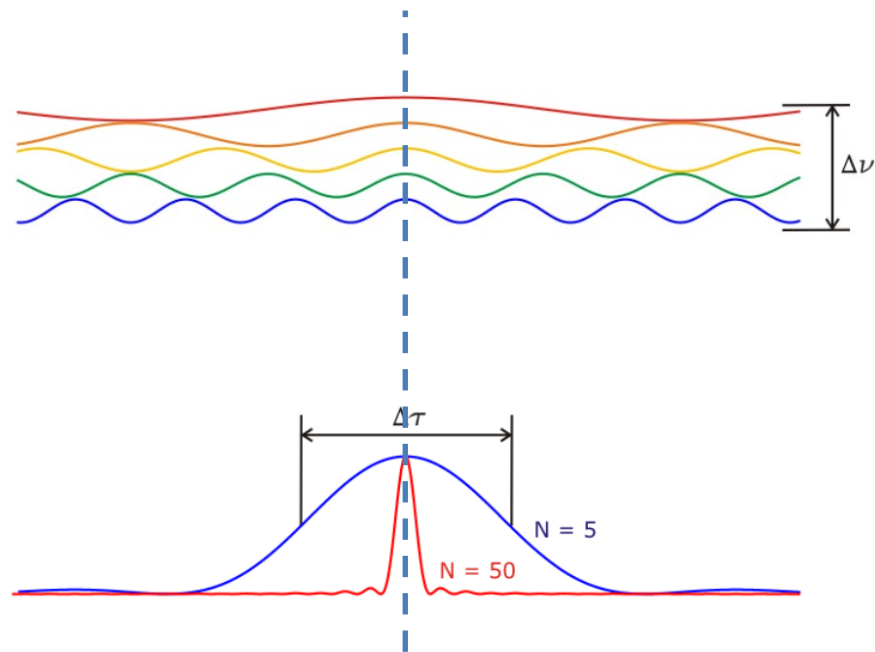


- amplitude and phase modulation
- amplitude- only modulation
- phase-only modulation

Spectrotemporal corrections

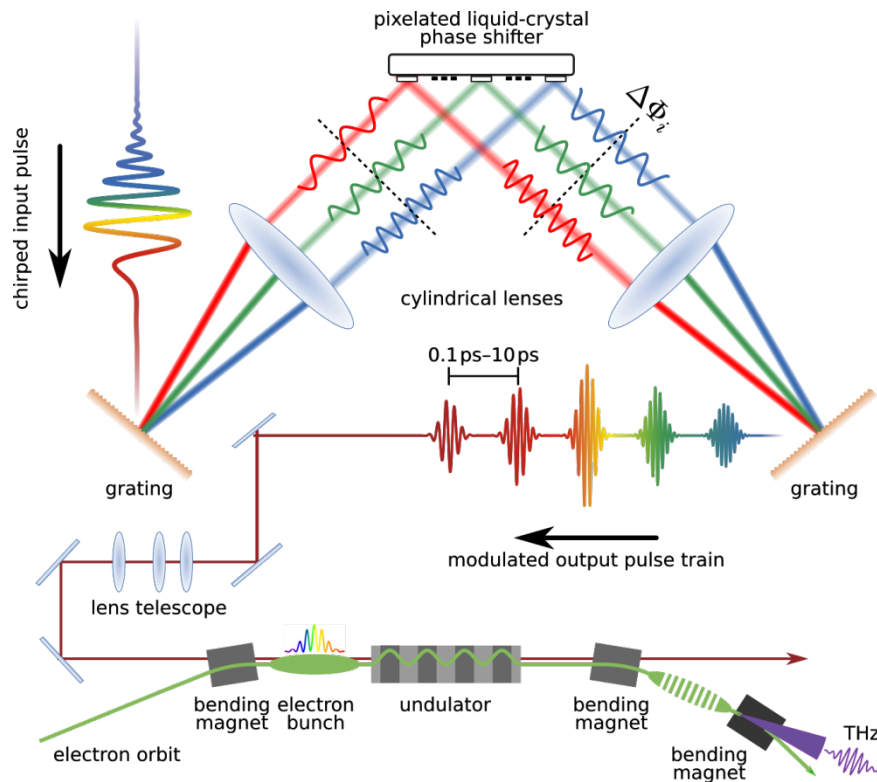
modelocked laser pulse:

contributing laser modes



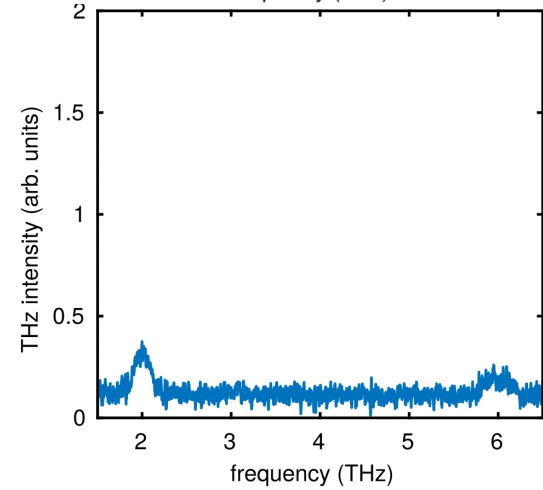
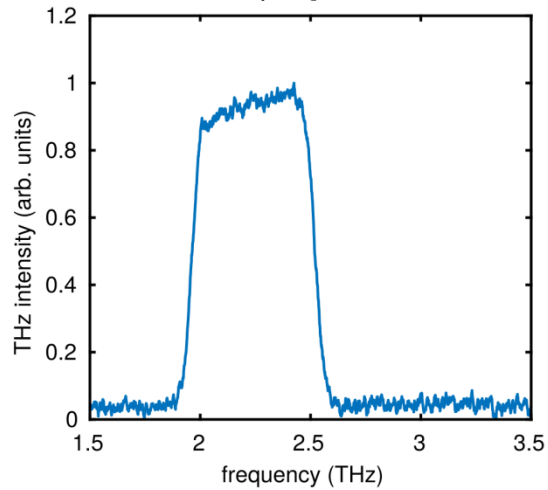
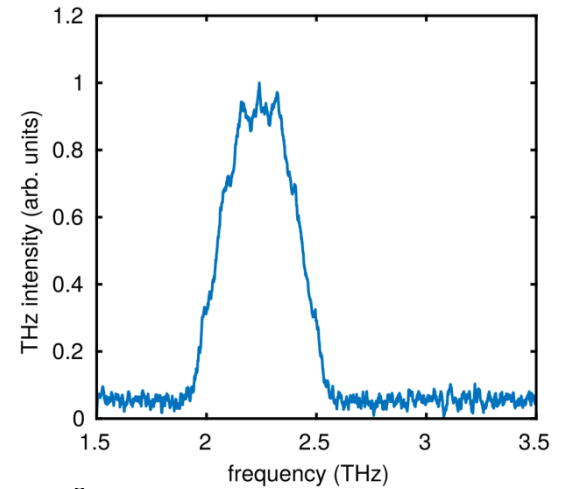
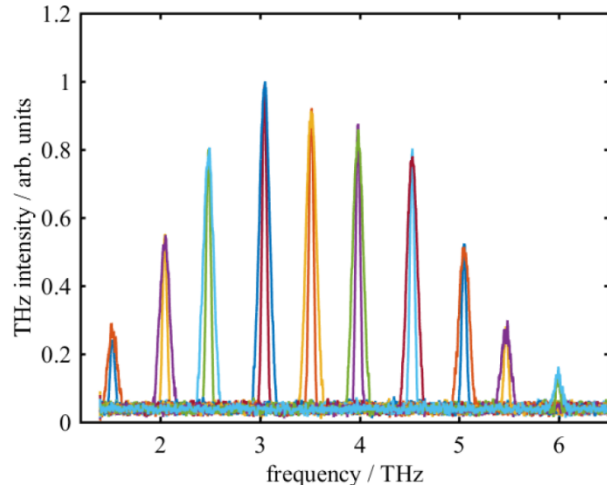
resulting laser pulse

Phase corrected optical setup for laser seeding



zero-dispersion compressor
(4-f shaper)

C. Mai et al., "Pulse shaping methods for laser-induced Generation of THz radiation at the DELTA storage ring", IPAC 2019.

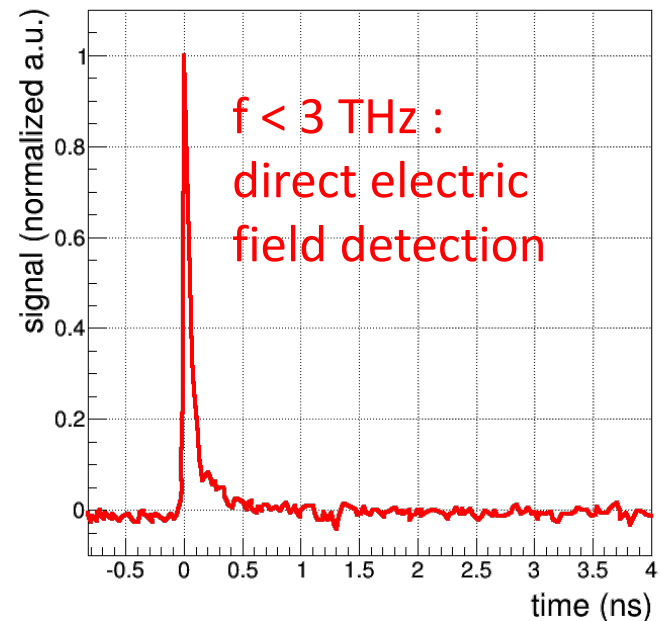
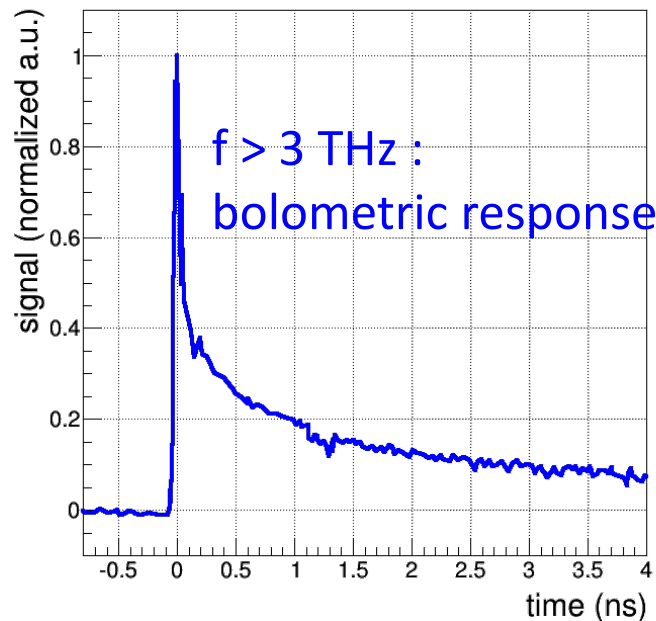


First Applications

- close collaborations to detector development
- frequency dependent behavior of $\text{YBa}_2\text{Cu}_3\text{O}_{7-x}$ -based detectors

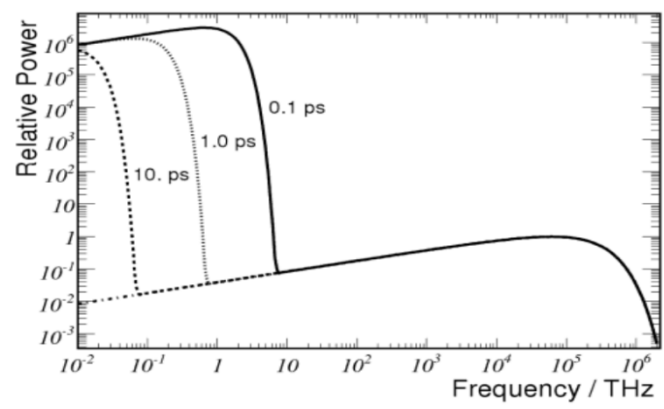


J. Raasch, KIT

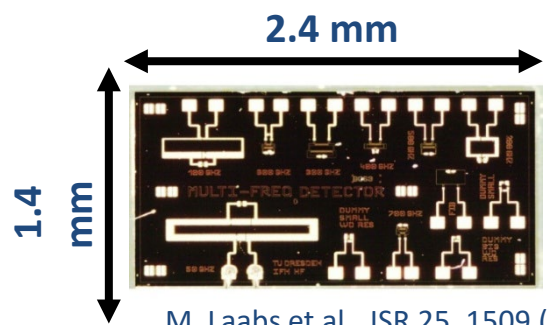
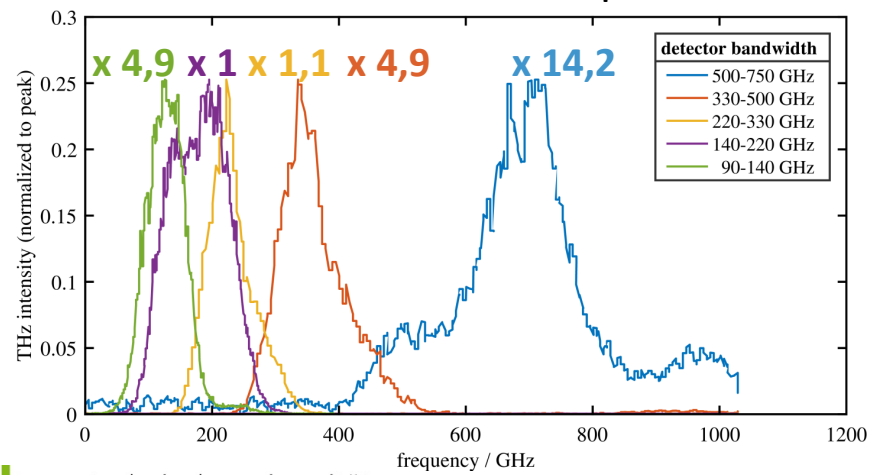


First Applications

- *best* detector for bunch length diagnostics?
- spectral response?
- pulse response?
- comparison under same conditions



direct detector comparison

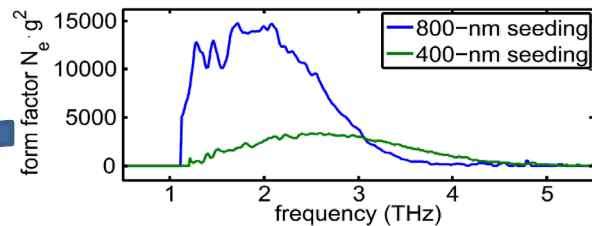


M. Laabs et al., JSR 25, 1509 (2018)

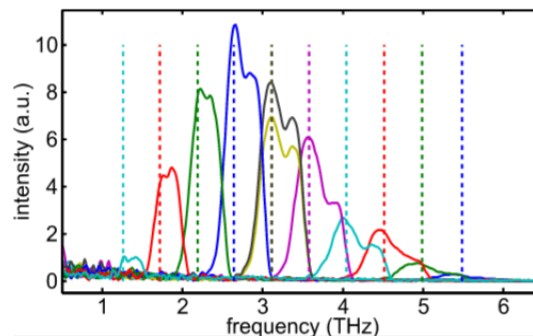


Summary

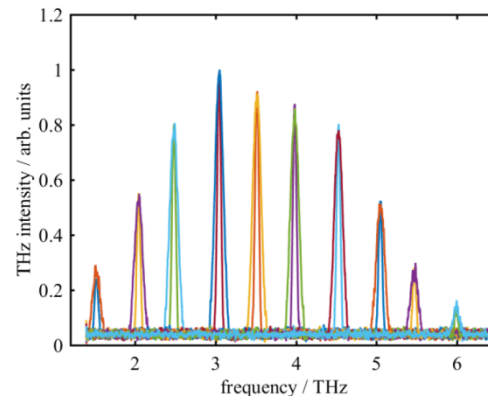
broadband generation



chirped-pulse beating

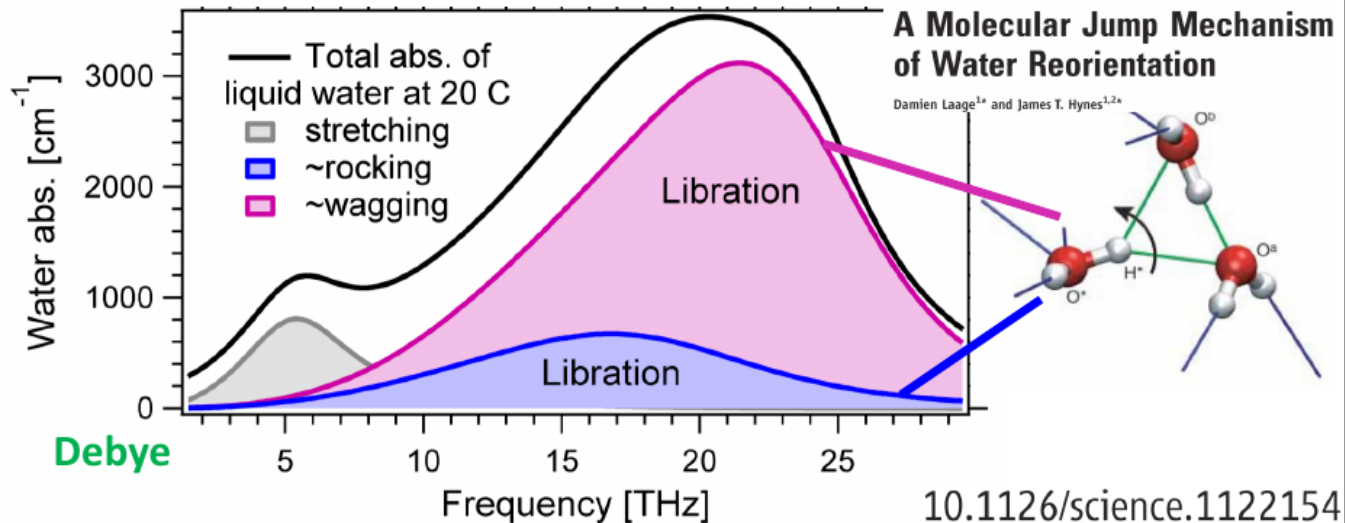


pulse shaping by phase modulation



Outlook

- shaping of laser pulses leads to highly flexible THz source
- full control of optical setups from DELTA control room
- new optics is promising: frequency shift, intensity scaling
- happy users from accelerator R&D groups
- user interest from physical chemistry: water absorption at 6 THz



F. Novelli, Ruhr Universität Bochum

Thank you for your attention!

Thanks for the support of:



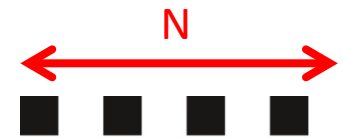
BACKUP

Imaging with spatial light modulators

- image is a diffraction pattern
 - Fourier transform (Fraunhofer approximation)

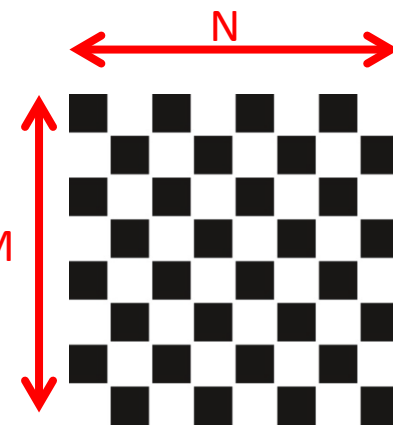
one dimension:

$$F(u) = \sum_m f(n) \cdot \exp(-i \cdot 2\pi / N \cdot n \cdot u)$$

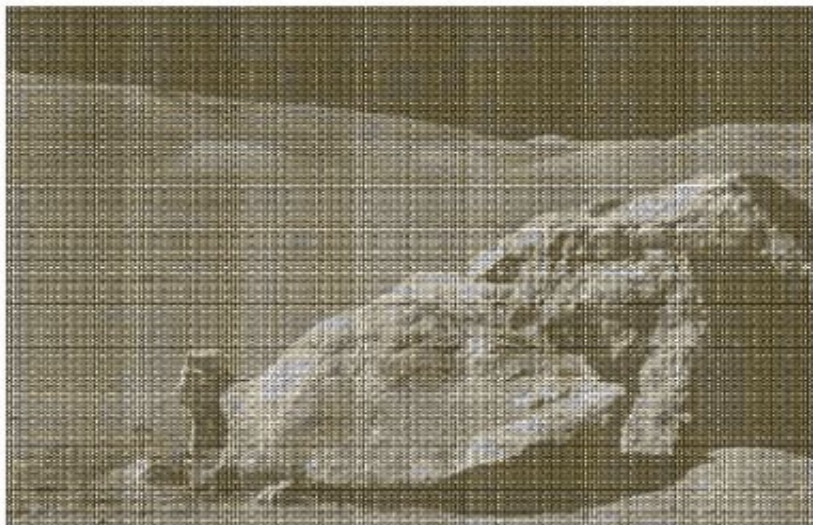


two dimensions:

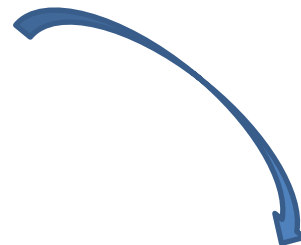
$$F(u, v) = \sum_{m=0}^{M-1} \sum_{n=0}^{N-1} f(m, n) \cdot \exp\left(-i 2\pi \left(\frac{um}{M} + \frac{vn}{N}\right)\right)$$



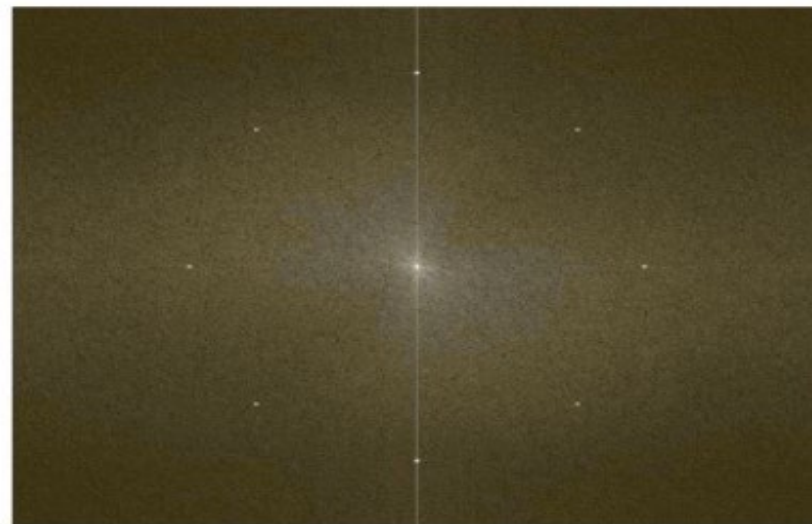
Is optics in Fourier domain of any use?

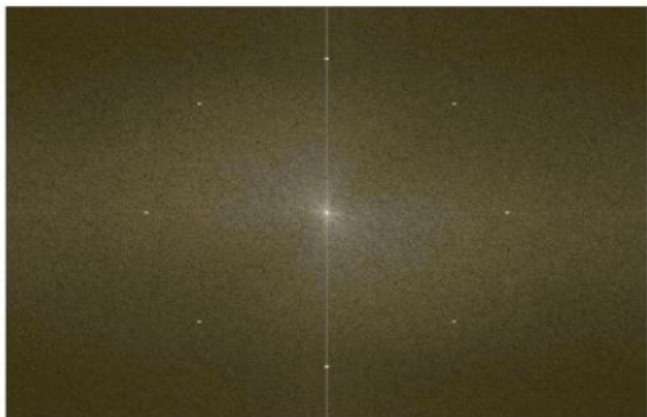


NASA (1969)

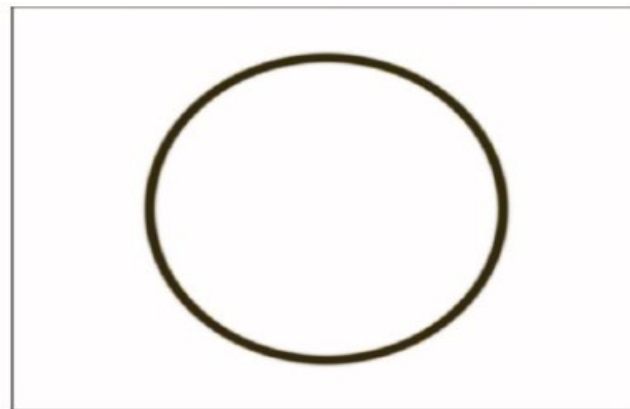


2D-FFT, amplitude spectrum





2D-FFT, amplitude spectrum

 \times 

spectral filter

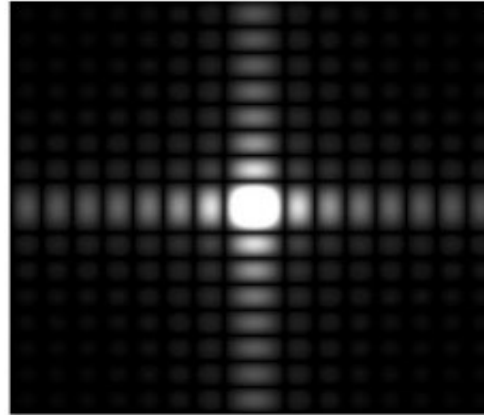
=



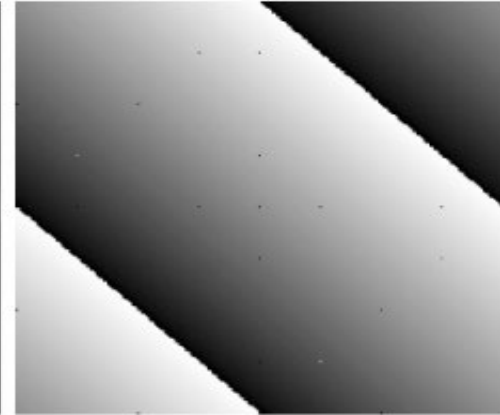
NASA (1969)

Phase and amplitude

Amplitude



Phase



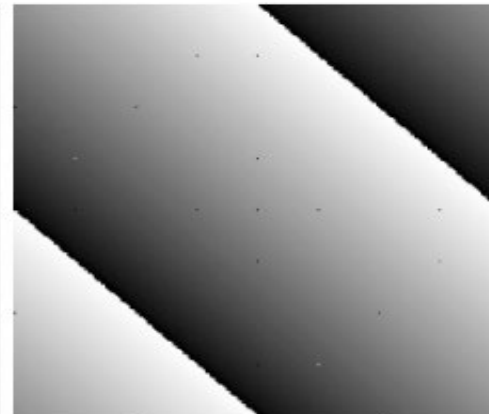
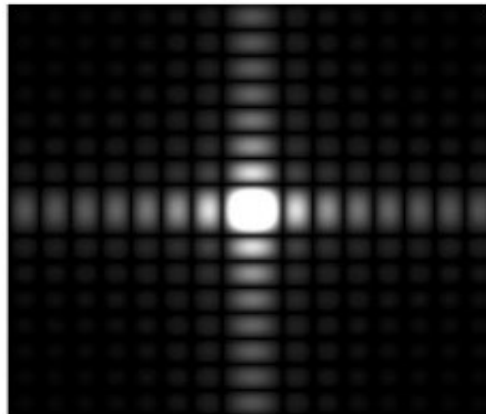
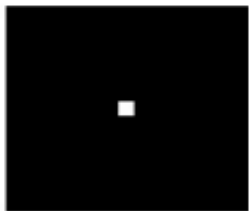
translation:



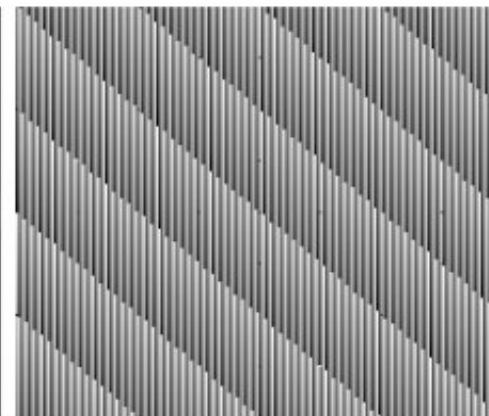
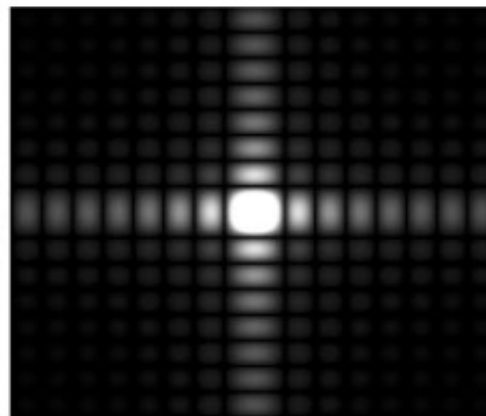
Phase and amplitude

Amplitude

Phase

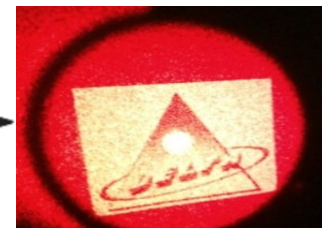
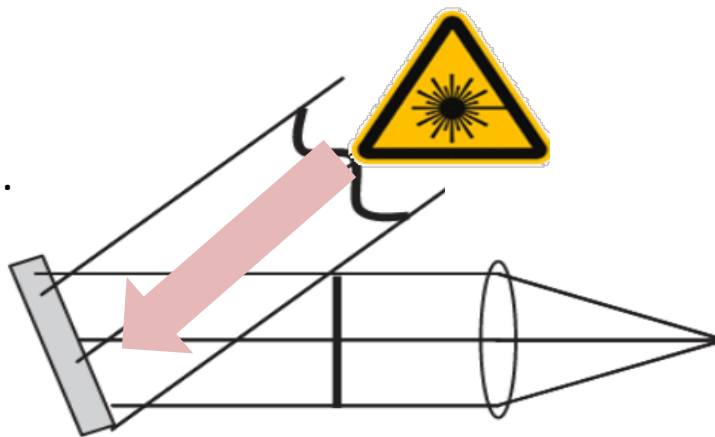
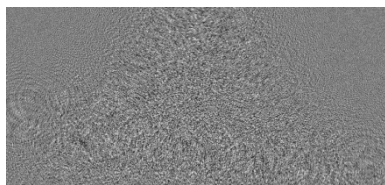


translation:



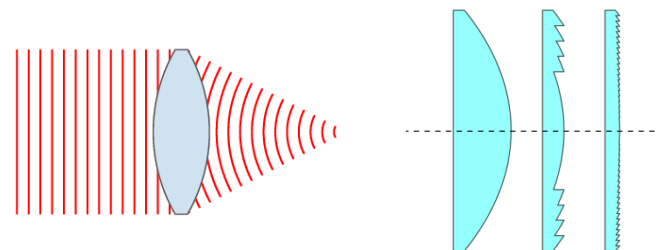
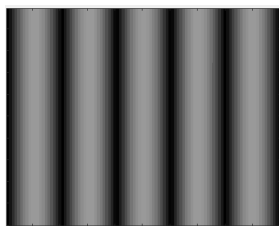
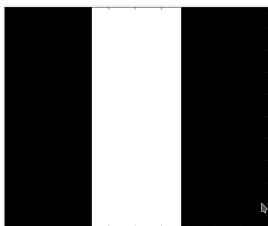
Holographic imaging by phase modulation

a complicated
phase-shifting pattern ...



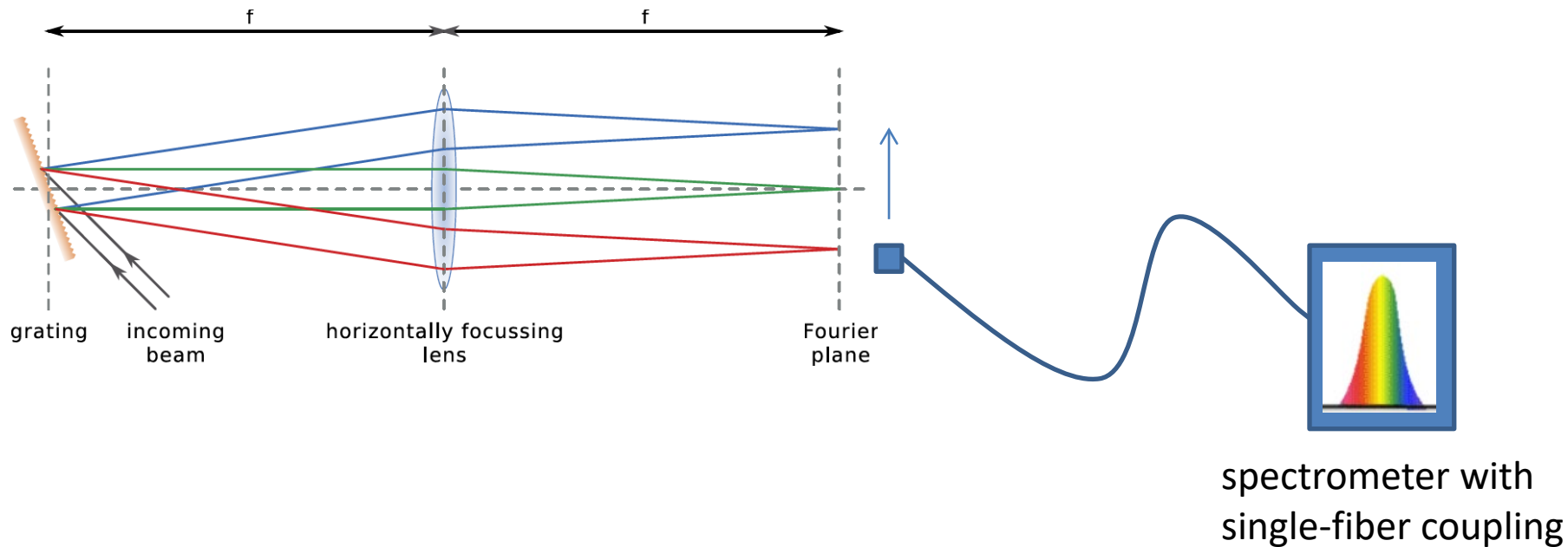
spatial modulator

elementary optics that can be implemented:



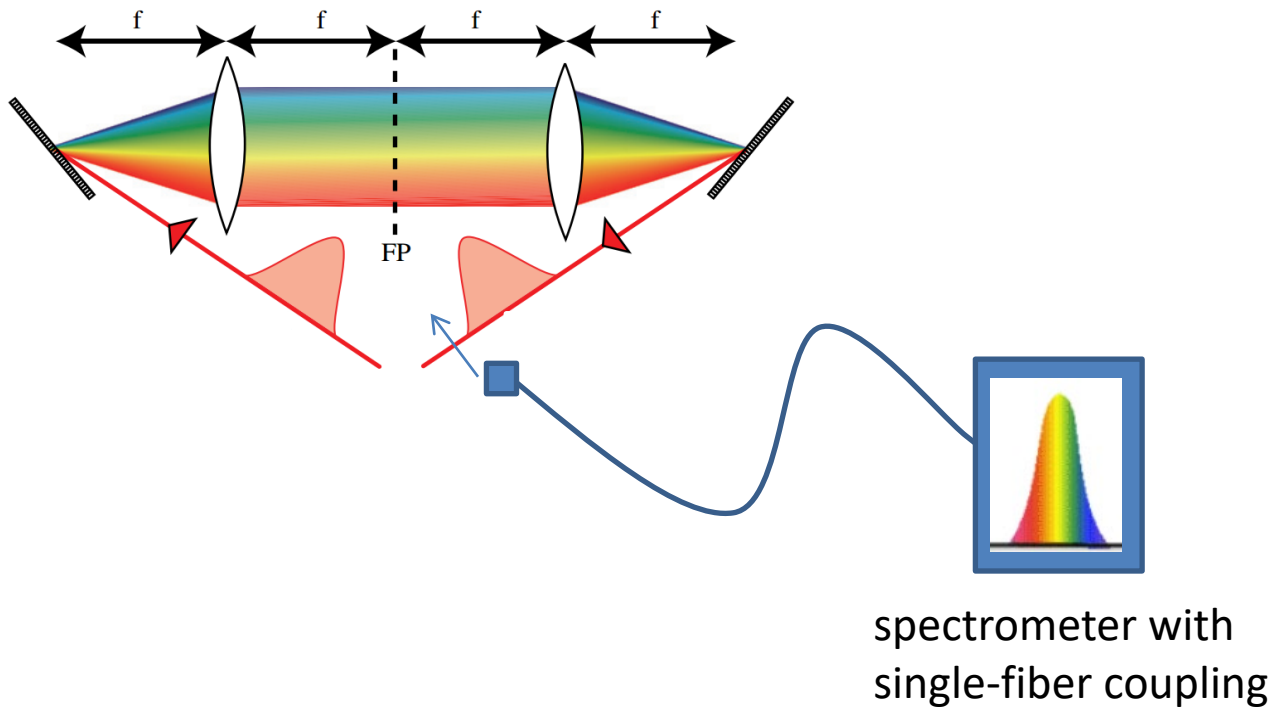
How to calibrate a 4-f pulse shaper

Step 1: tuning the focussing



How to calibrate a 4-f pulse shaper

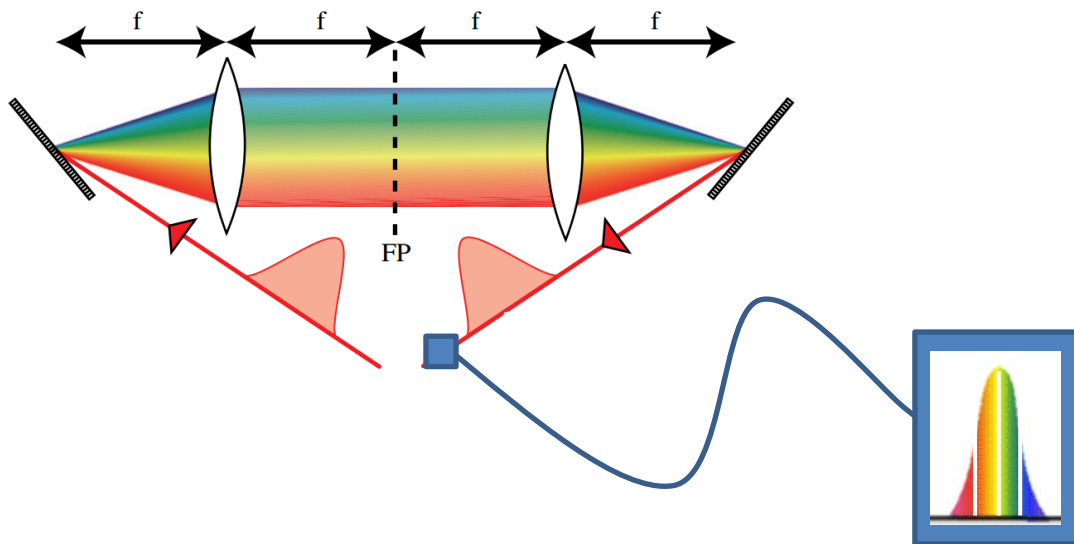
Step 2: match the other half, check for transverse chirp of the beam



How to calibrate a 4-f pulse shaper

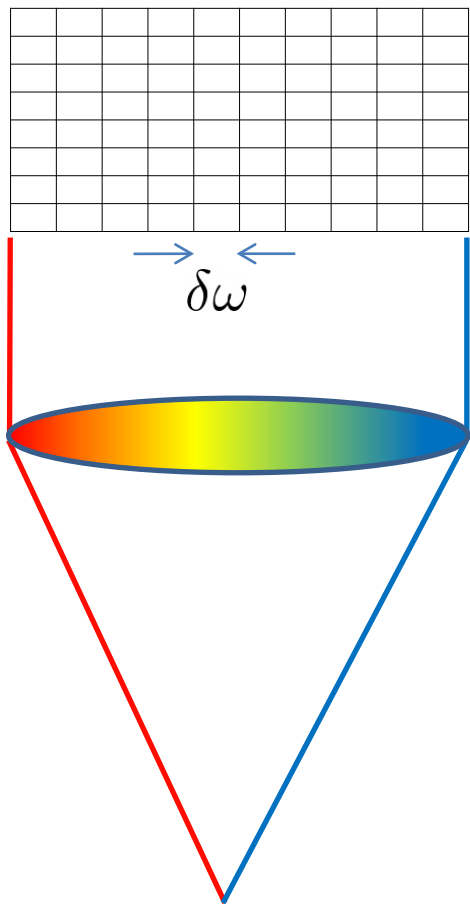
Step 3: calibrate spatial dispersion, apply phase shifts of π for single pixels(/columns)

$$X(\omega_k) - X_0 = f \cdot \tan \left[\arcsin \left(\frac{2\pi c}{\omega_k d} - \sin(\theta_i) \right) - \arcsin \left(\frac{2\pi c}{\omega_0 d} - \sin(\theta_i) \right) \right]$$



spectrometer shows
destructive interference
(intensity drops)

Limitations / Masking effect of the modulator



Spatial modulator: linear complex filter

$$\tilde{E}_{\text{out}}(\omega) = \tilde{M}(\omega) \cdot \tilde{E}_{\text{in}}(\omega)$$

frequency dependent
phase and amplitude

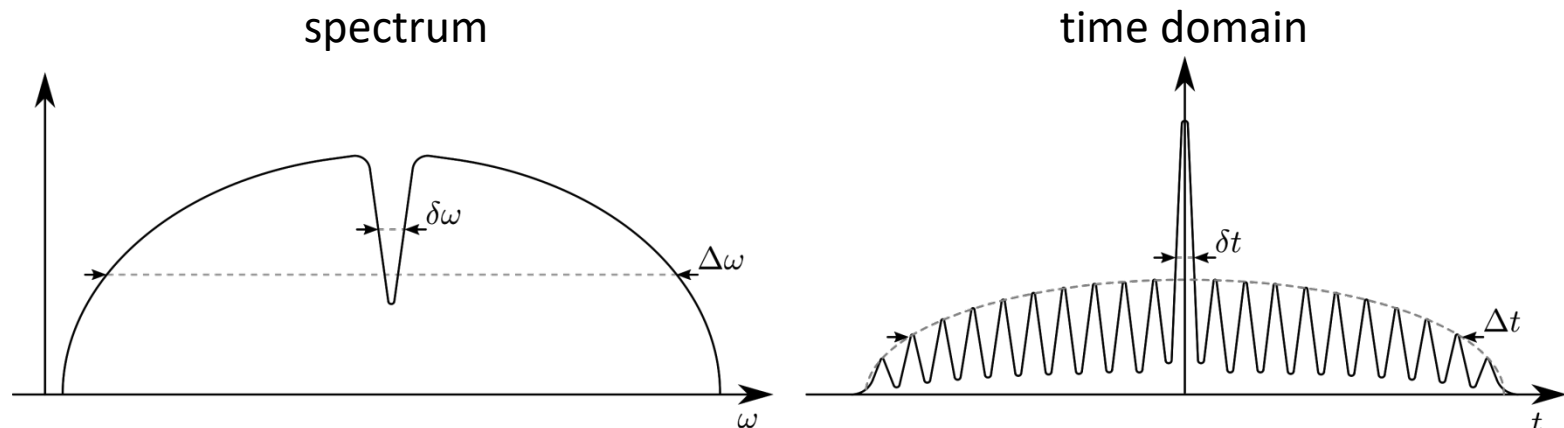
input field

$$\tilde{M}(\omega) = \left(\tilde{F}(\omega) \sum_{n=-N/2}^{n=N/2} \delta(\omega - n \delta\omega) \right) * \text{rect} \left(\frac{\omega}{\delta\omega} \right)$$

$$E_{\text{out}}(t) \sim \left(E_{\text{in}}(t) * \sum f \left(t - n \frac{2\pi}{\delta\omega} \right) \right) \text{sinc} \left(\frac{\delta\omega t}{2} \right)$$

pulse copies

Limitations / What is a (too) complex pulse?



time-bandwidth product (TBP):

$$TBP = \Delta t \cdot \Delta \omega$$

complexity:

$$\begin{aligned} \eta &= \frac{\Delta t}{\delta t} = \frac{\Delta \omega}{\delta \omega} \\ &= \frac{TBP}{4 \ln 2} \quad (\text{gaussian pulse}) \\ &\leq \text{number of modulator pixels} \end{aligned}$$

Spectral limitations of the beamline

gain factor

$$g^2(f) = \left| \int \rho(t) e^{i\omega t} dt \right|^2$$

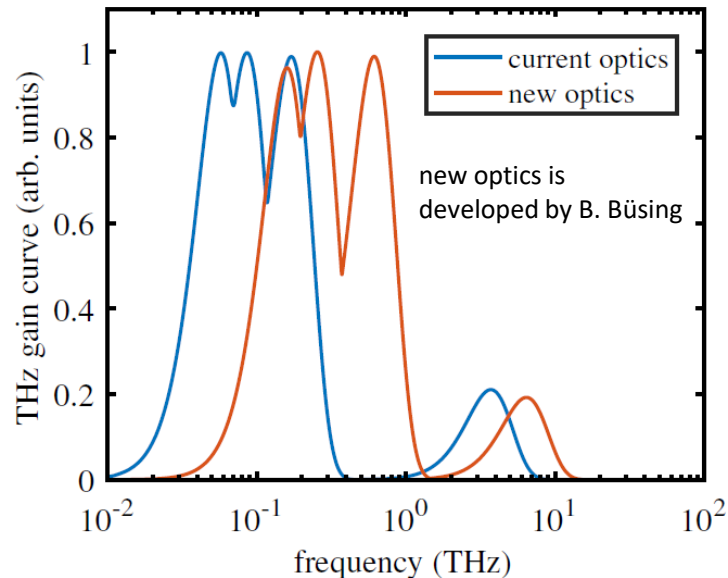
$$\propto \left(r_{56} \frac{\sigma_E}{E} \frac{2\pi f}{c} \right)^4$$

$$\cdot \exp \left[- \left(\frac{2\pi f}{c} \right) \left(r_{51}^2 \sigma_x^2 + r_{52}^2 \sigma_{x'}^2 + r_{56}^2 \frac{\sigma_E^2}{E^2} \right) \right].$$

Phys. Rev. ST Accel. Beams **13**, 090703 (2010)

- storage ring optics: r_{51} , r_{52} , r_{56}

Gain function of DELTA /
accessible frequency region



- new optics leads to modulation of more electrons
- increase of pulse energy