Formation of cylindrical and 3D quasiellipsoidal beams for laser driver



Outline

- Basic idea of pulse shaper
- Mathematical model for formation of cylindrical beams
- Experimental set-up
- Experimental results: formation beams with cylindrical and 3D quasi-ellipsoidal intensity distribution
- Conclusion



Basic idea of pulse shaper



Spectrum control:

$$S_{req}(\Omega) = M(\Omega) \cdot S_{in}(\Omega)$$

Fourier transform

M(Ω) – Spectral mask: in General case - complex function For linearly chirped pulses: $S(Ω) = |S_0(Ω)| \cdot e^{-i\frac{α \cdot Ω^2}{2}}$

At T>>Tf
$$|A(t)|^2 \propto |S(\Omega)|^2 \quad \Omega \propto t$$

Control at frequency domain linearly corresponds to control at time

Disadvantages:

- No axial-symmetry at the profiled beams
- Strong astigmatism from cylindrical telescope



Mathematical model for formation of cylindrical beams

Gaussian spectral intensity distribution:

$$S(\Omega) = S_0 \cdot e^{-2\ln(2) \cdot \frac{\Omega^2}{\Delta\Omega^2} - i\frac{\alpha \cdot \Omega^2}{2}} \implies A(t) = \int_{-\infty}^{\infty} S(\Omega) \cdot e^{i \cdot \Omega \cdot t} d\Omega = \frac{\sqrt{2\pi} \cdot \Delta\Omega \cdot S_0}{\sqrt{4 \cdot \ln(2) + i \cdot \alpha \cdot \Delta\Omega^2}} \cdot e^{-2\ln(2)\frac{t^2}{T^2} + i\frac{\alpha_t \cdot t^2}{2}}$$
$$T = T_F \sqrt{1 + \frac{16 \cdot \ln(2)^2 \cdot \alpha^2}{T_F^4}}, \quad \alpha_t = \frac{\alpha \cdot \Delta\Omega^2}{T_F^2 + \alpha^2 \cdot \Delta\Omega^2}$$

Spectral mask to produce rectangular spectral intensity: $|S_r(\Omega)|^2 = M(\Omega) \cdot |S(\Omega)|^2$

$$M(\Omega) = \begin{cases} \exp\left(4\ln(2)\cdot\left(\frac{\Omega^2 - \Lambda^2}{\Delta\Omega^2}\right)\right) & |\Omega| < \Lambda \\ 0 & |\Omega| > \Lambda \end{cases}$$

Energy efficiency:

$$\eta = \sqrt{\frac{16\ln(2)}{\pi}} \cdot \frac{\Lambda}{\Delta\Omega} \cdot e^{-4\ln(2)\frac{\Lambda^2}{\Delta\Omega^2}} \qquad \eta_{max} = \sqrt{\frac{2}{\pi e}} \sim 0.48 \qquad \Lambda_{max} = \frac{\Delta\Omega}{\sqrt{8\ln(2)}}$$

Optimal pulse chirping

How to find optimal α and T?



The amplitude mask doesn't know about spectral phase

 $2 \cdot \Lambda_{max} = \alpha_t \cdot T_r$

$$T \approx \sqrt{2\ln(2)}T_r \approx 1.17 \cdot T_r$$

Optimal spatial cutting



For cutting Gaussian beams by diaphragm N% intensity deviation from beam center to edge corresponds to N% energy transmission

Efficiency can be increased by implementation of "soft" diaphragm [1].

[1] A. K. Potemkin, et. all., "Compact neodymium phosphate glass laser emitting 100-J, 100-GW pulses for pumping a parametric amplifier of chirped pulses", QUANTUM ELECTRON, 2005, 35 (4), 302–310



Soft-spectral mask: pulse front and PV parameter

$$S(\Omega) = S_0 \cdot e^{-2\ln(2) \cdot \frac{\Omega^2}{\Delta \Omega^2} - i\frac{\alpha \cdot \Omega^2}{2}} \quad M_G(\Omega) = \exp\left[\left(-\frac{\Omega}{\Lambda}\right)^{2N}\right) \cdot \left\{ \exp\left(2\ln(2)\frac{\Omega^2 - \Lambda^2}{\Delta \Omega^2}\right), \quad |\Omega| < \Lambda$$
$$S_r(\Omega) = M_G(\Omega) \cdot S(\Omega)$$



 τ_F – Pulse front duration from 0.1*Max till 0.9*Max



Soft-spectral mask: pulse front and PV parameter

$$M_{G}(\Omega) = \exp\left[\left(-\frac{\Omega}{\Lambda}\right)^{2N}\right) \cdot \begin{cases} \exp\left(2\ln(2)\frac{\Omega^{2} - \Lambda^{2}}{\Delta\Omega^{2}}\right), & |\Omega| < \Lambda \\ 1, & |\Omega| > \Lambda \end{cases}$$





Experimental set up



Output energy: 5.7nJ, pulse duration (FWHM): 41 ps, spectral bandwidth 7.3nm



Used methods for beam diagnostics

Cross-correlator





Y



$$W_3(\tau) \sim \int_{-\infty}^{\infty} I_1(t-\tau) \cdot I_2(t) dt$$
$$I_2(t) \sim \delta(t) \ W_3(\tau) \sim I_1(\tau)$$

In this method we collect transversal cross-sections along beam

Quasi-cylindrical beams: spectral and temporal intensity distribution



Pulse fronts: $\tau_F = 4 ps$ At rectangular pulse duration 41 ps



Quasi-cylindrical beams: 3D intensity distribution

-27ps -14ps 14ps 27ps 0ps -0.5 -0.5 -0.5 -0.5 -0.5 шШ шШ шШ шШ mm 0 0 0 0 0 0.5 0.5 0.5 0.5 0.5 -0.5 0.5 -0.5 0.5 -0.5 0.5 -0.5 -0.5 0 0 0 0 0.5 0 0.5 mm mm mm mm mm 20 100 50 20 40 50 100 100 40 0 0 0 0 0

-6

0.4

0.4 mm

0 و 0.4 الح

Transversal intensity distribution



The measured data was obtained with help of Cross-correlator

-2



4

2

0

mm

3D quasi-ellipsoidal beams



Conclusion

- The mathematical analyze of implementation amplitude mask in pulse shaper has been done
- The pulse shaper scheme was reproduced in IAP RAS with SLM matrix from Hamamatsu
- Cross-correlator and image spectrometer was implemented for beam parameter measurement
- Cylindrical and quasi-ellipsoidal beams were obtained in experiments



Thank you for your attention!